

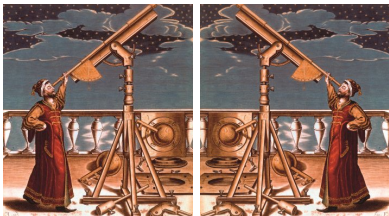


Doug Ravenel

Two equivariant approaches to the telescope conjecture

MIT Topology Seminar

November 20, 2017



Doug Ravenel
University of Rochester

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What is the telescope
conjecture?

The failed approach of MRS

The construction of
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Going equivariant I

Going equivariant II

This talk began in discussions last year with



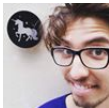
Agnes Beaudry



Mark Behrens



Prasit Bhattacharya



Dominic Culver



Zhouli Xu

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What is the telescope conjecture?

I first made the telescope conjecture in the late '70s and published it in 1984.

LOCALIZATION WITH RESPECT TO CERTAIN PERIODIC HOMOLOGY THEORIES

By DOUGLAS C. RAVENEL*

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It has a version for each prime p and each integer $n \geq 0$.

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It has a version for each prime p and each integer $n \geq 0$.

Let X be a p -local finite spectrum with
 $K(n)_*X \neq 0$ and $K(n-1)_*X = 0$.

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It has a version for each prime p and each integer $n \geq 0$.

Let X be a p -local finite spectrum with $K(n)_*X \neq 0$ and $K(n-1)_*X = 0$. Such complexes are known to exist for all n and p by a theorem of Steve Mitchell.



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Let X be a p -local finite spectrum with $K(n)_*X \neq 0$ and $K(n-1)_*X = 0$. We say that such a complex has **type n** .

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What is the telescope conjecture? (continued)

Let X be a p -local finite spectrum with $K(n)_*X \neq 0$ and $K(n-1)_*X = 0$. We say that such a complex has **type n** .



The Hopkins-Smith periodicity theorem says that any such complex admits a self-map $\Sigma^d X \rightarrow X$ for $d > 0$ that is a $K(n)$ -equivalence.



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Let \widehat{X} be the telescope obtained by iterating this map.



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Let \widehat{X} be the telescope obtained by iterating this map. **The telescope conjecture says it is equivalent to $L_{K(n)}X$.**

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The Hopkins-Smith periodicity theorem says that any such complex admits a self-map $\Sigma^d X \rightarrow X$ for $d > 0$ that is a $K(n)$ -equivalence.



Let \widehat{X} be the telescope obtained by iterating this map. **The telescope conjecture says it is equivalent to $L_{K(n)}X$.**



The $n = 1$ case was proved by Mahowald for $p = 2$ and by Miller for odd primes in 1981.



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In 1989 there was a homotopy theory program at MSRI.



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Something happened there that led me to think I could **disprove** the conjecture for $n \geq 2$.

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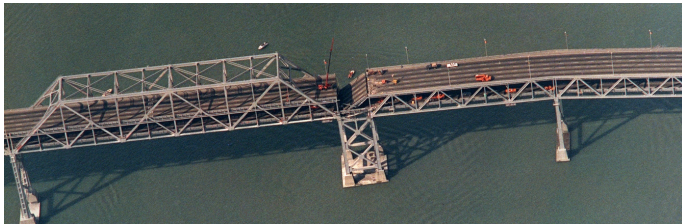
What is the telescope conjecture? (continued)



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Earthquake of October 17, 1989

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The failed approach of Mahowald-R-Shick

The disproof fell through a few years later.

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In 1999 I wrote a paper about it with Mark Mahowald and Paul Shick.

THE TRIPLE LOOP SPACE APPROACH TO THE
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MARK MAHOWALD, DOUGLAS RAVENEL AND PAUL SHICK



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DISCLAIMER: Having bet on
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DISCLAIMER: Having bet on both sides of this question, my credibility now stands at **ZERO**.



The failed MRS approach (continued)

The central character in our paper is a spectrum we call $y(n)$,

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The central character in our paper is a spectrum we call $y(n)$, which is defined for each prime p

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The central character in our paper is a spectrum we call $y(n)$, which is defined for each prime p and each integer $n > 0$. In this talk p will always be 2.

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The central character in our paper is a spectrum we call $y(n)$, which is defined for each prime p and each integer $n > 0$. In this talk p will always be 2.

I will outline the construction of $y(n)$ later in the talk.

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Our spectrum $y(n)$ has the following properties.

$$\textcircled{1} H_*(y(n); \mathbf{Z}/2) = \mathbf{Z}/2[\xi_1, \xi_2, \dots, \xi_n]$$

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- 1 $H_*(y(n); \mathbf{Z}/2) = \mathbf{Z}/2[\xi_1, \xi_2, \dots, \xi_n]$ where ξ_i is the Milnor generator of the dual Steenrod algebra A_* .

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- 2 It is an associative ring spectrum with a v_n self-map

$$v_n : \Sigma^{2(2^n-1)}y(n) \rightarrow y(n)$$

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- 3 There is a localized Adams spectral sequence converging to $\pi_* Y(n)$



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- 4 There is an Adams-Novikov spectral sequence converging to $\pi_* L_{K(n)}y(n)$,



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- 4 There is an Adams-Novikov spectral sequence converging to $\pi_* L_{K(n)}y(n)$, also with a known E_2 -term.



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- 5 There is a conjectured pattern of Adams differentials



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- 5 There is a conjectured pattern of Adams differentials that shows $Y(n)$ and $L_{K(n)}y(n)$ are **very different**.



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- 5 There is a conjectured pattern of Adams differentials that shows $Y(n)$ and $L_{K(n)}y(n)$ are **very different**. **If correct, it would disprove the telescope conjecture.**



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Our program failed because

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Our program failed because we could not rule out spurious Adams differentials that could mess up the calculation.

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OUR HOPE NOW:

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OUR HOPE NOW: By making $y(n)$ either the fixed point set or the underlying spectrum of a C_2 -equivariant spectrum,

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Experience has shown that an equivariant perspective

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Experience has shown that an equivariant perspective can lead to new insights into nonequivariant problems.

I will describe two different ways we might do this.

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OUR HOPE NOW: By making $y(n)$ either the fixed point set or the underlying spectrum of a C_2 -equivariant spectrum, we would have some additional structure that would give us more control over the Adams differentials.

Experience has shown that an equivariant perspective can lead to new insights into nonequivariant problems.

I will describe two different ways we might do this. **It is too early to tell if either approach will work.**

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Consider the diagram

$$\begin{array}{ccc} S^1 & \xrightarrow{f} & BO \\ & \searrow i & \nearrow g \\ & & \Omega^2 S^3 \end{array}$$

where

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Consider the diagram

$$\begin{array}{ccc} S^1 & \xrightarrow{f} & BO \\ & \searrow i & \nearrow g \\ & \Omega^2 S^3 & \end{array}$$

where

- f represents the nontrivial element of $\pi_1 BO = \mathbf{Z}/2$,

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- g is the extension of f given by the infinite loop space structure on BO .

We know that

$$H_* \Omega^2 S^3 = \mathbf{Z}/2[u_1, u_2, \dots] \quad \text{with } |u_n| = 2^n - 1 = |\xi_n|.$$



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$$\begin{array}{ccc} S^1 & \xrightarrow{f} & BO \\ & \searrow i & \nearrow g \\ & \Omega^2 S^3 & \end{array}$$

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Let $y(\infty)$ denote the Thom spectrum induced by g .

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We will construct subspaces Y_n of $\Omega^2 S^3$ with

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We will construct subspaces Y_n of $\Omega^2 S^3$ with

$$H_* Y_n = \mathbf{Z}/2[u_1, u_2, \dots, u_n],$$

and $y(n)$ will be the corresponding Thom spectrum.



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In the early 50s Ian James defined the reduced product $J_k X$ (for any space X) as a certain quotient of $X^{\times k}$

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In the early 50s Ian James defined the reduced product $J_k X$ (for any space X) as a certain quotient of $X^{\times k}$ and showed that $J_\infty X$ is equivalent to $\Omega\Sigma X$.

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In the early 50s Ian James defined the reduced product $J_k X$ (for any space X) as a certain quotient of $X^{\times k}$ and showed that $J_\infty X$ is equivalent to $\Omega \Sigma X$.

He showed there is a 2-local fiber sequence

$$\Omega^2 S^{2^{n+1}+1} \rightarrow J_{2^n-1} S^2 \rightarrow \Omega S^3 \rightarrow \Omega S^{2^{n+1}+1}.$$

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Note that ΩS^3 is equivalent to a CW complex with a single cell in each even dimension.

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Our space Y_n is $\Omega J_{2^n-1} S^2$,

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Our space Y_n is $\Omega J_{2^n-1} S^2$, so it maps to $\Omega^2 S^3$ as desired.

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Note that ΩS^3 is equivalent to a CW complex with a single cell in each even dimension. $J_{2^n-1} S^2$ is its $(2^{n+1} - 1)$ -skeleton.

Our space Y_n is $\Omega J_{2^n-1} S^2$, so it maps to $\Omega^2 S^3$ as desired. The MRS spectrum $y(n)$ is the Thomification of

$$\Omega J_{2^n-1} S^2 \longrightarrow \Omega^2 S^3 \xrightarrow{g} BO.$$



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The MRS spectrum $y(n)$ is the Thomification of

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From James' 2-local fiber sequence

$$\Omega^3 S^{2^{n+1}+1} \rightarrow \Omega J_{2^n-1} S^2 \rightarrow \Omega^2 S^3$$

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The MRS spectrum $y(n)$ is the Thomification of

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From James' 2-local fiber sequence

$$\Omega^3 S^{2^{n+1}+1} \rightarrow \Omega J_{2^n-1} S^2 \rightarrow \Omega^2 S^3$$

we get maps of spectra

$$\Sigma^\infty S^{|v_n|} \rightarrow \Sigma^\infty \Omega^3 S^{2^{n+1}+1} \rightarrow y(n) \rightarrow H\mathbb{Z}/2.$$

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The MRS spectrum $y(n)$ is the Thomification of

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$$\Sigma^\infty S^{|\nu_n|} \rightarrow \Sigma^\infty \Omega^3 S^{2^{n+1}+1} \rightarrow y(n) \rightarrow H\mathbb{Z}/2.$$

where the map $S^{|\nu_n|} \rightarrow \Omega^3 S^{2^{n+1}+1}$ is the inclusion of the bottom cell.

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The MRS spectrum $y(n)$ is the Thomification of

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From James' 2-local fiber sequence

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$$\Sigma^\infty S^{|v_n|} \rightarrow \Sigma^\infty \Omega^3 S^{2^{n+1}+1} \rightarrow y(n) \rightarrow H\mathbb{Z}/2.$$

where the map $S^{|v_n|} \rightarrow \Omega^3 S^{2^{n+1}+1}$ is the inclusion of the bottom cell. Since $y(n)$ is the Thom spectrum for a loop map, it is an associative ring spectrum. The composite map above leads to the desired v_n -self map of $y(n)$.

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If the MRS approach is to succeed, we need some more structure in the localized Adams spectral sequence for $Y(n)$.

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If the MRS approach is to succeed, we need some more structure in the localized Adams spectral sequence for $Y(n)$. Here I will outline the first of two ways to get $y(n)$ and $Y(n)$ into a C_2 -equivariant setting. Each of them will be a retract of the fixed point set of a C_2 -spectrum.



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Recall that the construction of $y(n)$ involved the diagram

$$\begin{array}{ccccc} S^1 & \xrightarrow{i} & \Omega^2 S^3 & \xrightarrow{g} & BO \\ & & \uparrow & & \\ & & \Omega J_{2^n-1} S^2 & & \end{array}$$



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We can add another space and get

$$\begin{array}{ccccc} S^1 & \xrightarrow{i} & \Omega^2 S^3 & \xrightarrow{g} & BO \\ & & \uparrow & & \uparrow \\ \Omega J_{2^n-1} S^2 & \Rightarrow & \Omega(SU(k+1)/SO(k+1)) & & \end{array} \quad \text{for } k \gg 0.$$



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$$\begin{array}{ccccc} S^1 & \xrightarrow{i} & \Omega^2 S^3 & \xrightarrow{g} & BO \\ & & \uparrow & & \uparrow a_k \\ & & \Omega J_{2^n-1} S^2 & \xrightarrow{g_n} & \Omega(SU(k+1)/SO(k+1)). \end{array}$$

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The map a_k for $k \gg 0$ is related to Bott's proof of his Periodicity Theorem.

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The map a_k for $k \gg 0$ is related to Bott's proof of his Periodicity Theorem. In mod 2 homology we have

$$H_* BO = \mathbf{Z}/2[b_1, b_2, \dots] \quad \text{where } |b_i| = i,$$

$$H_* \Omega(SU(k+1)/SO(k+1)) = \mathbf{Z}/2[b_1, \dots, b_k]$$

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and the loop map g_n exists for $k \geq 2^n - 1$.

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and the loop map g_n exists for $k \geq 2^n - 1$. Thomifying the square on the right gives

$$\begin{array}{ccc}
 H\mathbf{Z}/2 & \longrightarrow & MO \\
 \uparrow & & \uparrow \\
 y(n) & \longrightarrow & w(k),
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where $w(k)$ is the Thom spectrum induced by the map a_k . We can show that $w(k)$ splits as a wedge of suspensions of $y(n)$.



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One can show that

$$\begin{array}{ccccc} S^1 & \xrightarrow{i} & \Omega^2 S^3 & \xrightarrow{g} & BO \\ & & \uparrow & & \uparrow a_k \\ & & \Omega J_{2^n-1} S^2 & \xrightarrow{g_n} & \Omega(SU(k+1)/SO(k+1)). \end{array}$$

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 & & \uparrow & & \uparrow a_k \\
 & & \Omega J_{2^n-1} S^2 & \xrightarrow{g_n} & \Omega(SU(k+1)/SO(k+1)).
 \end{array}$$

is the fixed point set of the following diagram of C_2 -spaces:

$$\begin{array}{ccccc}
 S^\rho & \xrightarrow{i} & \Omega^{1+\rho} S^{1+2\rho} & \xrightarrow{g} & BU_{\mathbf{R}} \\
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 \end{array}$$

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One can show that

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where

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where

- $BU_{\mathbf{R}}$ and $SU_{\mathbf{R}}$ denote the spaces BU and SU equipped with a C_2 -action related to complex conjugation,
- σ denotes the sign representation of C_2 and
- $\rho = 1 + \sigma$ denotes its regular representation.



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with Thom spectra indicated on the right. Taking 2-local fibers of the vertical maps in the square gives

$$\begin{array}{ccc}
 \Omega^{1+\rho} S^{1+2\rho} & \xrightarrow{g} & BU_{\mathbf{R}} \\
 \uparrow & & \uparrow a_k \\
 \Omega^\rho J_{2^{n-1}} S^{2\rho} & \xrightarrow{g_n} & \Omega^\sigma SU(k+1)_{\mathbf{R}} \\
 \uparrow & & \uparrow \\
 \Omega^{2+\rho} S^{1+2^{n+1}\rho} & \longrightarrow & \Omega^\rho(SU/SU(k+1))_{\mathbf{R}}
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 \end{array}$$

The two fibers have the same connectivity when $k = 2^{n+1} - 2 = |v_n|$.



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$$S^{|v_n|} \rightarrow \Omega^3 S^{2^{n+1}+1} \rightarrow y(n) \rightarrow w(|v_n|),$$



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$$\begin{array}{ccc}
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$$\Sigma^{(1+|v_n|)\rho-1} X(|v_n|)_{\mathbf{R}} \rightarrow X(|v_n|)_{\mathbf{R}}.$$



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The underlying spectrum of this telescope is contractible



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The underlying spectrum of this telescope is contractible because the underlying map is known to be nilpotent.



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In this approach,

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In this approach, $y(n)$ and $Y(n)$ will be given nontrivial C_2 -actions.

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In this approach, $y(n)$ and $Y(n)$ will be given nontrivial C_2 -actions. They will be the **underlying spectra** of a pair of C_2 -spectra.

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Recall that the starting point of the construction of $y(n)$ was the diagram

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Here ρ denotes the (2-dimensional) regular representation of the group C_2 , S^V denote the one point compactification of V , and the twisted loop space $\Omega^\rho X$ is space of pointed continuous (but not necessarily equivariant) maps $S^\rho \rightarrow X$ for a pointed C_2 -space X .



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$$\begin{array}{ccc} S^1 & \xrightarrow{f} & BO \\ & \searrow i & \nearrow g \\ & \Omega^\rho S^{1+\rho} & \end{array}$$

It is known that BO is the 0th space in a C_2 Ω -spectrum,

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It is known that BO is the 0th space in a C_2 Ω -spectrum, so we can deloop it ρ times. This means f deloops to a map $S^{1+\rho} \rightarrow B^\rho BO$,

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The resulting equivariant Thom spectrum is the subject of a recent paper by Mark Behrens and Dylan Wilson.



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They show that it is the C_2 -spectrum $\underline{HZ}/2$,

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The resulting equivariant Thom spectrum is the subject of a recent paper by Mark Behrens and Dylan Wilson.



They show that it is the C_2 -spectrum $\underline{HZ}/2$, where $\underline{\mathbf{Z}}/2$ denotes the constant $\mathbf{Z}/2$ -valued Mackey functor,

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$$\begin{array}{ccc} S^1 & \xrightarrow{f} & BO \\ & \searrow i & \nearrow g \\ & \Omega^\rho S^{1+\rho} & \end{array}$$

It is known that BO is the 0th space in a C_2 Ω -spectrum, so we can deloop it ρ times. This means f deloops to a map $S^{1+\rho} \rightarrow B^\rho BO$, which leads to the map g above.



The resulting equivariant Thom spectrum is the subject of a recent paper by Mark Behrens and Dylan Wilson.



They show that it is the C_2 -spectrum $H\underline{\mathbf{Z}}/2$, where $\underline{\mathbf{Z}}/2$ denotes the constant $\mathbf{Z}/2$ -valued Mackey functor, [as expected](#).

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$$\begin{array}{ccc} S^1 & \xrightarrow{f} & BO \\ & \searrow i & \nearrow g \\ & \Omega^\rho S^{1+\rho} = \Omega^{1+\sigma} S^{2+\sigma} & \end{array}$$

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Behrens-Wilson show that the Thom spectrum of g is the C_2 -equivariant mod 2 Eilenberg-Mac Lane spectrum $H\mathbb{Z}/2$.

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There are two ways we might do this equivariantly.

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1



Use a twisted version of the James construction, due to Slawomir Rybicki, to filter the twisted loop space $\Omega^\sigma S^{1+\rho} = \Omega^\sigma \Sigma^\sigma S^2$.



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Use a twisted version of the James construction, due to Slawomir Rybicki, to filter the twisted loop space $\Omega^\sigma S^{1+\rho} = \Omega^\sigma \Sigma^\sigma S^2$.

2 Use the usual James construction to filter $\Omega S^{1+\rho}$,



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Use a twisted version of the James construction, due to Slawomir Rybicki, to filter the twisted loop space $\Omega^\sigma S^{1+\rho} = \Omega^\sigma \Sigma^\sigma S^2$.

- 2 Use the usual James construction to filter $\Omega S^{1+\rho}$, and then look at twisted loop spaces of certain of its skeleta.



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There are two ways we might study $\Omega^\rho S^{1+\rho} = \Omega^{1+\sigma} S^{2+\sigma}$.

- 1 Use a twisted version of the James construction, due to Rybicki, to filter the twisted loop space $\Omega^\sigma \Sigma^\sigma S^2$, and then look at the loop spaces of certain of its skeleta.
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There are two ways we might study $\Omega^\rho S^{1+\rho} = \Omega^{1+\sigma} S^{2+\sigma}$.

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The first approach does not appear to give us the right Thom spectra.

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There are two ways we might study $\Omega^\rho S^{1+\rho} = \Omega^{1+\sigma} S^{2+\sigma}$.

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The first approach does not appear to give us the right Thom spectra.

The second approach may give the right Thom spectra,

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There are two ways we might study $\Omega^\rho S^{1+\rho} = \Omega^{1+\sigma} S^{2+\sigma}$.

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The first approach does not appear to give us the right Thom spectra.

The second approach may give the right Thom spectra, but it suffers from a technical problem.

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There are two ways we might study $\Omega^\rho S^{1+\rho} = \Omega^{1+\sigma} S^{2+\sigma}$.

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The first approach does not appear to give us the right Thom spectra.

The second approach may give the right Thom spectra, but it suffers from a technical problem. The twisted loop space $\Omega^\sigma X$ of a C_2 -space X is **NOT** a C_2 -H-space,

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There are two ways we might study $\Omega^\rho S^{1+\rho} = \Omega^{1+\sigma} S^{2+\sigma}$.

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- 2 Use the James construction to filter $\Omega \Sigma S^{1+\sigma}$, and then look at the twisted loop spaces of certain of its skeleta.

The first approach does not appear to give us the right Thom spectra.

The second approach may give the right Thom spectra, but it suffers from a technical problem. The twisted loop space $\Omega^\sigma X$ of a C_2 -space X is **NOT** a C_2 -H-space, even though it is underlain by an ordinary H-space.

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- 2 Use the James construction to filter $\Omega S^{1+\rho}$, and then look at the twisted loop spaces of certain of its skeleta.

The twisted loop space $\Omega^\sigma X$ of a C_2 -space X is **NOT** an H-space.



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- 2 Use the James construction to filter $\Omega S^{1+\rho}$, and then look at the twisted loop spaces of certain of its skeleta.

The twisted loop space $\Omega^\sigma X$ of a C_2 -space X is **NOT** an H-space. The reason for this is that there is **no equivariant pinch map**



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- 2 Use the James construction to filter $\Omega S^{1+\rho}$, and then look at the twisted loop spaces of certain of its skeleta.

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The twisted loop space $\Omega^\sigma X$ of a C_2 -space X is **NOT** an H-space. The reason for this is that there is **no equivariant pinch map** $S^\sigma \rightarrow S^\sigma \vee S^\sigma$. We would need one to get the multiplication map

$$\Omega^\sigma X = \text{Map}_*(S^\sigma, X) \leftarrow \text{Map}_*(S^\sigma \vee S^\sigma, X) = \Omega^\sigma X \times \Omega^\sigma X.$$



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Instead there is cofiber sequence

$$S^0 \rightarrow S^\sigma \rightarrow C_{2+} \wedge S^1.$$



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This leads to a **twisted multiplication**

$$\Omega^\sigma X \leftarrow \text{Map}_*(C_{2+} \wedge S^1, X) = N^{C_2} \Omega X,$$



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This leads to a **twisted multiplication**

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where the space on the right is the **C_2 -norm** of X .



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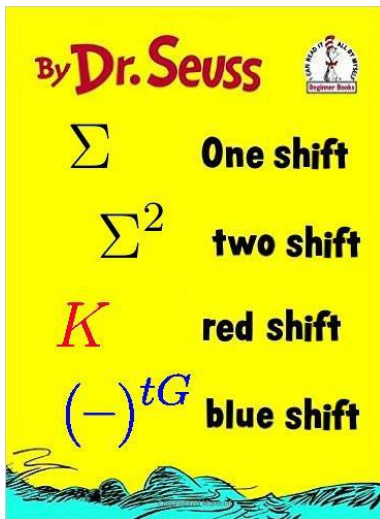
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THANK YOU!

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