

More stable stems

Daniel C. Isaksen

Guozhen Wang

Zhouli Xu

Author address:

DEPARTMENT OF MATHEMATICS, WAYNE STATE UNIVERSITY, DETROIT, MI
48202, USA

E-mail address: isaksen@wayne.edu

SHANGHAI CENTER FOR MATHEMATICAL SCIENCES, FUDAN UNIVERSITY,
SHANGHAI, CHINA, 200433

E-mail address: wanguozhen@fudan.edu.cn

DEPARTMENT OF MATHEMATICS, MASSACHUSETTS INSTITUTE OF TECHNOLOGY,
CAMBRIDGE, MA 02139

E-mail address: xuzhouli@mit.edu

Contents

List of Tables	ix
Chapter 1. Introduction	1
1.1. New Ingredients	2
1.2. Main results	4
1.3. Remaining uncertainties	9
1.4. Groups of homotopy spheres	9
1.5. Notation	10
1.6. How to use this manuscript	12
1.7. Acknowledgements	12
Chapter 2. Background	13
2.1. Associated graded objects	13
2.2. Motivic modular forms	16
2.3. The cohomology of the \mathbb{C} -motivic Steenrod algebra	17
2.4. Toda brackets	17
Chapter 3. The algebraic Novikov spectral sequence	25
3.1. Naming conventions	25
3.2. Machine computations	26
3.3. h_1 -Bockstein spectral sequence	27
Chapter 4. Massey products	29
4.1. The operator g	29
4.2. The Mahowald operator	30
4.3. Additional computations	30
Chapter 5. Adams differentials	33
5.1. The Adams d_2 differential	33
5.2. The Adams d_3 differential	34
5.3. The Adams d_4 differential	40
5.4. The Adams d_5 differential	45
5.5. Higher differentials	48
Chapter 6. Toda brackets	53
Chapter 7. Hidden extensions	59
7.1. Hidden τ extensions	59
7.2. Hidden 2 extensions	61
7.3. Hidden η extensions	69
7.4. Hidden ν extensions	76

7.5. Miscellaneous hidden extensions	83
7.6. Additional relations	86
Chapter 8. Tables	89
Bibliography	127

Abstract

We compute the stable homotopy groups up to dimension 90, except for some carefully enumerated uncertainties.

2010 *Mathematics Subject Classification*. Primary 14F42, 55Q45, 55S10, 55T15; Secondary 16T05, 55P42, 55Q10, 55S30 .

Key words and phrases. stable homotopy group, stable motivic homotopy theory, May spectral sequence, Adams spectral sequence, cohomology of the Steenrod algebra, Adams-Novikov spectral sequence.

The first author was supported by NSF grant DMS-1606290. The second author was supported by grant NSFC-11801082. The third author was supported by NSF grant DMS-1810638. Many of the associated machine computations were performed on the Wayne State University Grid high performance computing cluster.

List of Tables

1	Stable homotopy groups up to dimension 90	5
1	Notation for $\pi_{*,*}$	91
2	Some hidden values of the unit map of mmf	92
3	Some Massey products in $\text{Ext}_{\mathbb{C}}$	93
4	E_2 -page generators of the \mathbb{C} -motivic Adams spectral sequence	94
5	Some permanent cycles in the \mathbb{C} -motivic Adams spectral sequence	103
6	Adams d_3 differentials in the \mathbb{C} -motivic Adams spectral sequence	104
7	Adams d_4 differentials in the \mathbb{C} -motivic Adams spectral sequence	107
8	Adams d_5 differentials in the \mathbb{C} -motivic Adams spectral sequence	109
9	Higher Adams differentials in the \mathbb{C} -motivic Adams spectral sequence	109
10	Unknown Adams differentials in the \mathbb{C} -motivic Adams spectral sequence	110
11	Some Toda brackets	111
12	Hidden values of inclusion of the bottom cell into $C\tau$	114
13	Hidden values of projection from $C\tau$ to the top cell	114
14	Hidden τ extensions in the \mathbb{C} -motivic Adams spectral sequence	116
15	Hidden 2 extensions	118
16	Possible hidden 2 extensions	119
17	Hidden η extensions	120
18	Possible hidden η extensions	122
19	Hidden ν extensions	123
20	Possible hidden ν extensions	124
21	Miscellaneous hidden extensions	125

CHAPTER 1

Introduction

The computation of stable homotopy groups of spheres is one of the most fundamental and important problems in homotopy theory. It has connections to many topics in topology, such as the cobordism theory of framed manifolds, the classification of smooth structures on spheres, obstruction theory, the theory of topological modular forms, algebraic K-theory, motivic homotopy theory, and equivariant homotopy theory.

Despite their simple definition, which was available eighty years ago, these groups are notoriously hard to compute. All known methods only give a complete answer through a range, and then reach an obstacle until a new method is introduced. The standard approach to computing stable stems is to use Adams type spectral sequences that converge from algebra to homotopy. In turn, to identify the algebraic E_2 -pages, one needs algebraic spectral sequences that converge from simpler algebra to more complicated algebra. For any spectral sequence, difficulties arise in computing differentials and in solving extension problems. Different methods lead to trade-offs. One method may compute some types of differentials and extension problems efficiently, but leave other types unanswered, perhaps even unsolvable by that technique. To obtain complete computations, one must be eclectic, applying and combining different methodologies. Even so, combining all known methods, there are eventually some problems that cannot be solved. Mahowald's uncertainty principle states that no finite collection of methods can completely compute the stable homotopy groups of spheres.

Because stable stems are finite groups (except for the 0-stem), the computation is most easily accomplished by working one prime at a time. At odd primes, the Adams-Novikov spectral sequence and the chromatic spectral sequence, which are based on complex cobordism and formal groups, have yielded a wealth of data [36]. As the prime grows, so does the range of computation. For example, at the primes 3 and 5, we have complete knowledge up to around 100 and 1000 stems respectively [36].

The prime 2, being the smallest prime, remains the most difficult part of the computation. In this case, the Adams spectral sequence is the most effective tool. The manuscript [16] presents a careful analysis of the Adams spectral sequence, in both the classical and \mathbb{C} -motivic contexts, that is essentially complete through the 59-stem. This includes a verification of the details in the classical literature [2] [3] [6] [29]. Subsequently, the second and third authors computed the 60-stem and 61-stem [44].

We also mention [25] [26], which take an entirely different approach to computing stable homotopy groups. However, the computations in [25] [26] are now known to contain several errors. See [44, Section 2] for a more detailed discussion.

The goal of this manuscript is to continue the analysis of the Adams spectral sequence into higher stems at the prime 2. We will present information up to the 90-stem. While we have not been able to resolve all of the possible differentials in this range, we enumerate the handful of uncertainties explicitly. See especially Table 10 for a summary of the possible differentials that remain unresolved.

The charts in [19] and [20] are an essential companion to this manuscript. They present the same information in an easily interpretable graphical format.

Our analysis uses various methods and techniques, including machine-generated homological algebra computations, a deformation of homotopy theories that connects \mathbb{C} -motivic and classical stable homotopy theory, and the theory of motivic modular forms. Here is a quick summary of our approach:

- (1) Compute the cohomology of the \mathbb{C} -motivic Steenrod algebra by machine. These groups serve as the input to the \mathbb{C} -motivic Adams spectral sequence.
- (2) Compute by machine the algebraic Novikov spectral sequence that converges to the cohomology of the Hopf algebroid (BP_*, BP_*BP) . This includes all differentials, and the multiplicative structure of the cohomology of (BP_*, BP_*BP) .
- (3) Identify the \mathbb{C} -motivic Adams spectral sequence for the cofiber of τ with the algebraic Novikov spectral sequence [11]. This includes an identification of the cohomology of (BP_*, BP_*BP) with the homotopy groups of the cofiber of τ .
- (4) Pull back and push forward Adams differentials for the cofiber of τ to Adams differentials for the \mathbb{C} -motivic sphere, along the inclusion of the bottom cell and the projection to the top cell.
- (5) Deduce additional Adams differentials for the \mathbb{C} -motivic sphere with a variety of ad hoc arguments. The most important methods are Toda bracket shuffles and comparison to the motivic modular forms spectrum mmf [10].
- (6) Deduce hidden τ extensions in the \mathbb{C} -motivic Adams spectral sequence for the sphere, using a long exact sequence in homotopy groups.
- (7) Obtain the classical Adams spectral sequence and the classical stable homotopy groups by inverting τ .

The machine-generated data that we obtain in steps (1) and (2) are available at [42]. See also [43] for a discussion of the implementation of the machine computation.

Much of this process is essentially automatic. The exception occurs in step (5) where ad hoc arguments come into play.

This document describes the results of this systematic program through the 90-stem. We anticipate that our approach will allow us to compute into even higher stems, especially towards the last unsolved Kervaire invariant problem in dimension 126. However, we have not yet carried out a careful analysis.

1.1. New Ingredients

We discuss in more detail several new ingredients that allow us to carry out this program.

1.1.1. Machine-generated algebraic data. The Adams-Novikov spectral sequence has been used very successfully to carry out computations at odd primes.

However, at the prime 2, its usage has not been fully exploited in stemwise computations. This is due to the difficulty of computing its E_2 -page. The first author predicted in [16] that “the next major breakthrough in computing stable stems will involve machine computation of the Adams-Novikov E_2 -page.”

The second author achieved this machine computation; the resulting data is available at [42]. The process goes roughly like this. Start with a minimal resolution that computes the cohomology of the Steenrod algebra. Lift this resolution to a resolution of BP_*BP . Finally, use the Curtis algorithm to compute the homology of the resulting complex, and to compute differentials in the associated algebraic spectral sequences, such as the algebraic Novikov spectral sequence and the Bockstein spectral sequence. See [43] for further details.

1.1.2. Motivic homotopy theory. The \mathbb{C} -motivic stable homotopy category gives rise to new methods to compute stable stems. These ideas are used in a critical way in [16] to compute stable stems up to the 59-stem.

The key insight of this article that distinguishes it significantly from [16] is that \mathbb{C} -motivic cellular stable homotopy theory is a deformation of classical stable homotopy theory [11]. From this perspective, the “generic fiber” of \mathbb{C} -motivic stable homotopy theory is classical stable homotopy theory, and the “special fiber” has an entirely algebraic description. The special fiber is the category of BP_*BP -comodules, or equivalently, the category of quasicohherent sheaves on the moduli stack of 1-dimensional formal groups.

In more concrete terms, let $C\tau$ be the cofiber of the \mathbb{C} -motivic stable map τ . The homotopy category of $C\tau$ -modules has an algebraic structure [11]. In particular, the \mathbb{C} -motivic Adams spectral sequence for $C\tau$ is isomorphic to the algebraic Novikov spectral sequence that computes the E_2 -page of the Adams-Novikov spectral sequence. Using naturality of Adams spectral sequences, the differentials in the algebraic Novikov spectral sequence, which are computed by machine, can be lifted to differentials in the \mathbb{C} -motivic Adams spectral sequence for the \mathbb{C} -motivic sphere spectrum. Then the Betti realization functor produces differentials in the classical Adams spectral sequence.

Our use of \mathbb{C} -motivic stable homotopy theory appears to rely on the fundamental computations, due to Voevodsky [40] [41], of the motivic cohomology of a point and of the motivic Steenrod algebra. In fact, recent progress has determined that our results do not depend on this deep and difficult work. There are now purely topological constructions of homotopy categories that have identical computational properties to the cellular stable \mathbb{C} -motivic homotopy category [10] [35]. In these homotopy categories, one can obtain from first principles the fundamental computations of the cohomology of a point and of the Steenrod algebra, using only well-known classical computations. Therefore, the material in this manuscript does not logically depend on Voevodsky’s work, even though the methods were very much inspired by his groundbreaking computations.

1.1.3. Motivic modular forms. In classical chromatic homotopy theory, the theory of topological modular forms, introduced by Hopkins and Mahowald [39], plays a central role in the computations of the $K(2)$ -local sphere.

Using a topological model of the cellular stable \mathbb{C} -motivic homotopy category, one can construct a “motivic modular forms” spectrum mmf [10], whose motivic

cohomology is the quotient of the \mathbb{C} -motivic Steenrod algebra by its subalgebra generated by Sq^1 , Sq^2 , and Sq^4 . Just as tmf plays an essential role in studies of the classical Adams spectral sequence, mmf is an essential tool for motivic computations. The \mathbb{C} -motivic Adams spectral sequence for mmf can be analyzed completely [17], and naturality of Adams spectral sequences along the unit map of mmf provides much information about the behavior of the \mathbb{C} -motivic Adams spectral sequence for the \mathbb{C} -motivic sphere spectrum.

1.2. Main results

We summarize our main results in the following theorem and corollaries.

THEOREM 1.1. *The \mathbb{C} -motivic Adams spectral sequence for the \mathbb{C} -motivic sphere spectrum is displayed in the charts in [19], up to the 90-stem.*

The proof of Theorem 1.1 consists of a series of specific computational facts, which are verified throughout this manuscript.

COROLLARY 1.2. *The classical Adams spectral sequence for the sphere spectrum is displayed in the charts in [19], up to the 90-stem.*

Corollary 1.2 follows immediately from Theorem 1.1. One simply inverts τ , or equivalently ignores τ -torsion.

COROLLARY 1.3. *The Adams-Novikov spectral sequence for the sphere spectrum is displayed in the charts in [20].*

Corollary 1.3 also follows immediately from Theorem 1.1. As described in [16, Chapter 6], the Adams-Novikov spectral sequence can be reverse-engineered from information about \mathbb{C} -motivic stable homotopy groups.

COROLLARY 1.4. *Table 1 describes the stable homotopy groups π_k for all values of k up to 90.*

We adopt the following notation in Table 1. An integer n stands for the cyclic abelian group \mathbb{Z}/n ; the symbol \cdot by itself stands for the trivial group; the expression $n \cdot m$ stands for the direct sum $\mathbb{Z}/n \oplus \mathbb{Z}/m$; and n^j stands for the direct sum of j copies of \mathbb{Z}/n . The horizontal line after dimension 61 indicates the range in which our computations are new information.

Table 1 describes each group π_k as the direct sum of three subgroups: the 2-primary v_1 -torsion, the odd primary v_1 -torsion, and the v_1 -periodic subgroups.

The last column of Table 1 describes the groups of homotopy spheres that classify smooth structures on spheres in dimensions at least 5. See Section 1.4 and Theorem 1.7 for more details.

Starting in dimension 82, there remain some uncertainties in the 2-primary v_1 -torsion. In most cases, these uncertainties mean that the order of some stable homotopy groups are known only up to factors of 2. In a few cases, the additive group structures are also undetermined.

These uncertainties have two causes. First, there are a handful of differentials that remain unresolved. Table 10 describes these. Second, there are some possible hidden 2 extensions that remain unresolved.

Table 1: Stable homotopy groups up to dimension 90

k	v_1 -torsion at the prime 2	v_1 -torsion at odd primes	v_1 -periodic	group of smooth structures
1	·	·	2	·
2	·	·	2	·
3	·	·	8·3	·
4	·	·	·	?
5	·	·	·	·
6	2	·	·	·
7	·	·	16·3·5	$\underline{b_2}$
8	2	·	2	$\underline{2}$
9	2	·	2 ²	$\underline{2} \cdot 2^2$
10	·	3	2	2·3
11	·	·	8·9·7	$\underline{b_3}$
12	·	·	·	·
13	·	3	·	3
14	2·2	·	·	2
15	2	·	32·3·5	$\underline{b_4} \cdot 2$
16	2	·	2	$\underline{2}$
17	2 ²	·	2 ²	$\underline{2} \cdot 2^3$
18	8	·	2	2·8
19	2	·	8·3·11	$\underline{b_5} \cdot 2$
20	8	3	·	8·3
21	2 ²	·	·	$\underline{2} \cdot 2^2$
22	2 ²	·	·	2 ²
23	2·8	3	16·9·5·7·13	$\underline{b_6} \cdot 2 \cdot 8 \cdot 3$
24	2	·	2	$\underline{2}$
25	·	·	2 ²	$\underline{2} \cdot 2$
26	2	3	2	2 ² ·3
27	·	·	8·3	$\underline{b_7}$
28	2	·	·	$\underline{2}$
29	·	3	·	3
30	2	3	·	3
31	2 ²	·	64·3·5·17	$\underline{b_8} \cdot 2^2$
32	2 ³	·	2	$\underline{2^3}$
33	2 ³	·	2 ²	$\underline{2} \cdot 2^4$
34	2 ² ·4	·	2	2 ³ ·4
35	2 ²	·	8·27·7·19	$\underline{b_9} \cdot 2^2$
36	2	3	·	2·3
37	2 ²	3	·	$\underline{2} \cdot 2^2 \cdot 3$
38	2·4	3·5	·	2·4·3·5
39	2 ⁵	3	16·3·25·11	$\underline{b_{10}} \cdot 2^5 \cdot 3$
40	2 ⁴ ·4	3	2	2 ⁴ ·4·3
41	2 ³	·	2 ²	$\underline{2} \cdot 2^4$
42	2·8	3	2	2 ² ·8·3

Table 1: Stable homotopy groups up to dimension 90

k	v_1 -torsion at the prime 2	v_1 -torsion at odd primes	v_1 -periodic	group of smooth structures
43	\cdot	\cdot	$8 \cdot 3 \cdot 23$	$\underline{b_{11}}$
44	8	\cdot	\cdot	8
45	$2^3 \cdot 16$	9·5	\cdot	$\underline{2} \cdot 2^3 \cdot 16 \cdot 9 \cdot 5$
46	2^4	3	\cdot	$2^4 \cdot 3$
47	$2^3 \cdot 4$	3	$32 \cdot 9 \cdot 5 \cdot 7 \cdot 13$	$\underline{b_{12}} \cdot 2^3 \cdot 4 \cdot 3$
48	$2^3 \cdot 4$	\cdot	2	$\underline{2^3} \cdot 4$
49	\cdot	3	2^2	$\underline{2} \cdot 2 \cdot 3$
50	2^2	3	2	$2^3 \cdot 3$
51	$2 \cdot 8$	\cdot	$8 \cdot 3$	$\underline{b_{13}} \cdot 2 \cdot 8$
52	2^3	3	\cdot	$\underline{2^3} \cdot 3$
53	2^4	\cdot	\cdot	$\underline{2} \cdot 2^4$
54	$2 \cdot 4$	\cdot	\cdot	$2 \cdot 4$
55	\cdot	3	$16 \cdot 3 \cdot 5 \cdot 29$	$\underline{b_{14}} \cdot 3$
56	\cdot	\cdot	2	\cdot
57	2	\cdot	2^2	$\underline{2} \cdot 2^2$
58	2	\cdot	2	2^2
59	2^2	\cdot	$8 \cdot 9 \cdot 7 \cdot 11 \cdot 31$	$\underline{b_{15}} \cdot 2^2$
60	4	\cdot	\cdot	4
61	\cdot	\cdot	\cdot	\cdot
62	2^4	3	\cdot	$2^3 \cdot 3$
63	$2^2 \cdot 4$	\cdot	$128 \cdot 3 \cdot 5 \cdot 17$	$\underline{b_{16}} \cdot 2^2 \cdot 4$
64	$2^5 \cdot 4$	\cdot	2	$\underline{2^5} \cdot 4$
65	$2^7 \cdot 4$	3	2^2	$\underline{2} \cdot 2^8 \cdot 4 \cdot 3$
66	$2^5 \cdot 8$	\cdot	2	$2^6 \cdot 8$
67	$2^3 \cdot 4$	\cdot	$8 \cdot 3$	$\underline{b_{17}} \cdot 2^3 \cdot 4$
68	2^3	3	\cdot	$\underline{2^3} \cdot 3$
69	2^4	\cdot	\cdot	$\underline{2} \cdot 2^4$
70	$2^5 \cdot 4^2$	\cdot	\cdot	$2^5 \cdot 4^2$
71	$2^6 \cdot 4 \cdot 8$	\cdot	$16 \cdot 27 \cdot 5 \cdot 7 \cdot 13 \cdot 19 \cdot 37$	$\underline{b_{18}} \cdot 2^6 \cdot 4 \cdot 8$
72	2^7	3	2	$\underline{2^7} \cdot 3$
73	2^5	\cdot	2^2	$\underline{2} \cdot 2^6$
74	4^3	3	2	$2 \cdot 4^3 \cdot 3$
75	2	9	$8 \cdot 3$	$\underline{b_{19}} \cdot 2 \cdot 9$
76	$2^2 \cdot 4$	5	\cdot	$\underline{2^2} \cdot 4 \cdot 5$
77	$2^5 \cdot 4$	\cdot	\cdot	$\underline{2} \cdot 2^5 \cdot 4$
78	$2^3 \cdot 4^2$	3	\cdot	$2^3 \cdot 4^2 \cdot 3$
79	$2^6 \cdot 4$	\cdot	$32 \cdot 3 \cdot 25 \cdot 11 \cdot 41$	$\underline{b_{20}} \cdot 2^6 \cdot 4$
80	2^8	\cdot	2	$\underline{2^8}$
81	$2^3 \cdot 4 \cdot 8$	3^2	2^2	$\underline{2} \cdot 2^4 \cdot 4 \cdot 8 \cdot 3^2$
82	$2^5 \cdot 8$ or $2^4 \cdot 8$ or $2^3 \cdot 4 \cdot 8$	3·7	2	$2^6 \cdot 8 \cdot 3 \cdot 7$ or $2^5 \cdot 8 \cdot 3 \cdot 7$ or $2^4 \cdot 4 \cdot 8 \cdot 3 \cdot 7$
83	$2^3 \cdot 8$ or $2^3 \cdot 4$	5	$8 \cdot 9 \cdot 49 \cdot 43$	$\underline{b_{21}} \cdot 2^3 \cdot 8 \cdot 5$ or $\underline{b_{21}} \cdot 2^3 \cdot 4 \cdot 5$
84	2^6 or 2^5	3^2	\cdot	$\underline{2^6} \cdot 3^2$ or $2^5 \cdot 3^2$

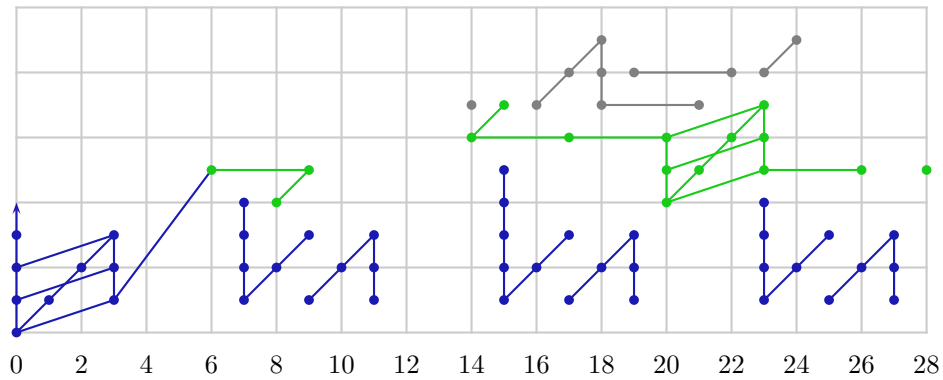
Table 1: Stable homotopy groups up to dimension 90

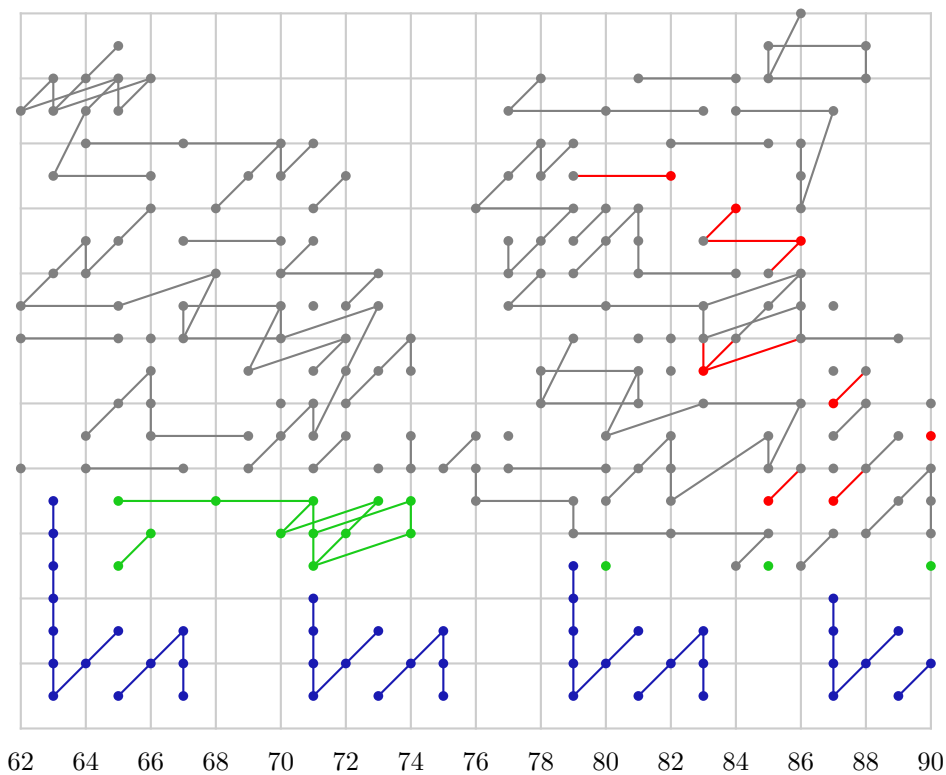
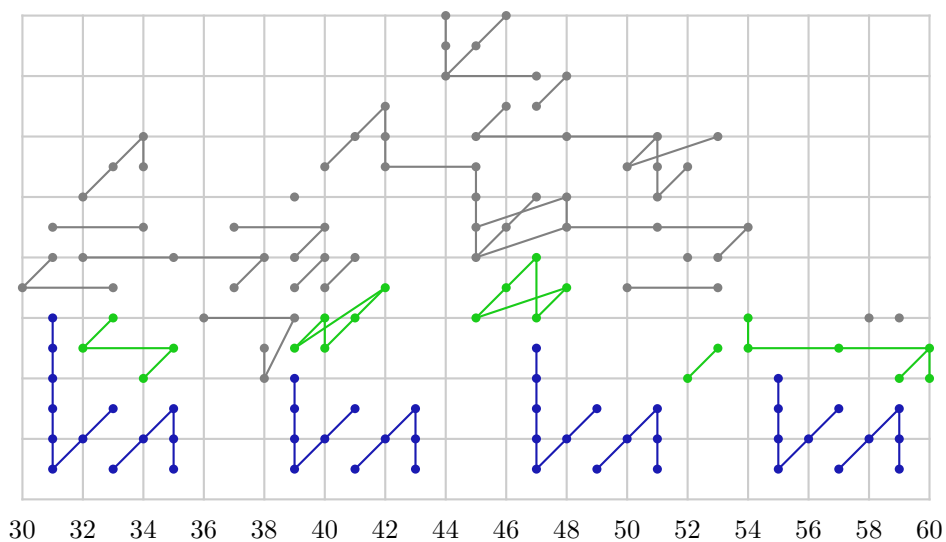
k	v_1 -torsion at the prime 2	v_1 -torsion at odd primes	v_1 -periodic	group of smooth structures
85	$2^6 \cdot 4^2$ or $2^5 \cdot 4^2$ or $2^4 \cdot 4^3$	3^2	.	$2^6 \cdot 4^2 \cdot 3^2$ or $2^5 \cdot 4^2 \cdot 3^2$ or $2^4 \cdot 4^3 \cdot 3^2$
86	$2^5 \cdot 8^2$ or $2^4 \cdot 8^2$ or $2^3 \cdot 4 \cdot 8^2$ or $2^2 \cdot 4 \cdot 8^2$	$3 \cdot 5$.	$2^5 \cdot 8^2 \cdot 3 \cdot 5$ or $2^4 \cdot 8^2 \cdot 3 \cdot 5$ or $2^3 \cdot 4 \cdot 8^2 \cdot 3 \cdot 5$ or $2^2 \cdot 4 \cdot 8^2 \cdot 3 \cdot 5$
87	2^8 or 2^7 or $2^6 \cdot 4$ or $2^5 \cdot 4$.	$16 \cdot 3 \cdot 5 \cdot 23$	$\underline{b_{22}} \cdot 2^8$ or $\underline{b_{22}} \cdot 2^7$ or $\underline{b_{22}} \cdot 2^6 \cdot 4$ or $\underline{b_{22}} \cdot 2^5 \cdot 4$
88	$2^4 \cdot 4$.	2	$2^4 \cdot 4$
89	2^3	.	2^2	$2 \cdot 2^4$
90	$2^3 \cdot 8$ or $2^2 \cdot 8$	3	2	$2^4 \cdot 8 \cdot 3$ or $2^3 \cdot 8 \cdot 3$

Figure 1 displays the 2-primary stable homotopy groups in a graphical format originally promoted by Allen Hatcher. Vertical chains of n dots indicate $\mathbb{Z}/2^n$. The non-vertical lines indicate multiplications by η and ν . The blue dots represent the v_1 -periodic subgroups. The green dots are associated to the topological modular forms spectrum tmf . These elements are detected by the unit map from the sphere spectrum to tmf , either in homotopy or in the algebraic Ext groups that serve as Adams E_2 -pages.

Finally, the red dots indicate uncertainties. In addition, in higher stems, there are possible extensions by 2, η , and ν that are not indicated in Figure 1. See Tables 16, 18, and 20 for more details about these possible extensions.

FIGURE 1. 2-primary stable homotopy groups





The orders of individual 2-primary stable homotopy groups do not follow a clear pattern, with large increases and decreases seemingly at random. However, an empirically observed pattern emerges if we consider the cumulative size of the groups, i.e., the product of the orders of all 2-primary stable homotopy groups from dimension 1 to dimension k .

Our data strongly suggest that asymptotically, there is a linear relationship between k^2 and the logarithm of this product of orders. In other words, the number of dots in Figure 1 in stems 1 through k is linearly proportional to k^2 . Correspondingly, the number of dots in the classical Adams E_∞ -page in stems 1 through k is linearly proportional to k^2 . Thus, in extending from dimension 60 to dimension 90, the overall size of the computation more than doubles. Specifically, through dimension 60, the cumulative rank of the Adams E_∞ -page is 199, and is 435 through dimension 90. Similarly, through dimension 60, the cumulative rank of the Adams E_2 -page is 488, and is 1,461 through dimension 90.

CONJECTURE 1.5. *Let $f(k)$ be the product of the orders of the 2-primary stable homotopy groups in dimensions 1 through k . There exists a non-zero constant C such that*

$$\lim_{k \rightarrow \infty} \frac{\log_2 f(k)}{k^2} = C.$$

One interpretation of this conjecture is that the expected value of the logarithm of the order of the 2-primary component of π_k grows linearly in k . We have only data to support the conjecture, and we have not formulated a mathematical rationale. It is possible that in higher stems, new phenomena occur that alter the growth rate of the stable homotopy groups.

By comparison, data indicates that the growth rate of the Adams E_2 -page is qualitatively greater than the growth rate of the Adams E_∞ -page. This apparent mismatch has implications for the frequency of Adams differentials.

1.3. Remaining uncertainties

Some uncertainties remain in the analysis of the first 90 stable stems. Table 10 lists all possible differentials in this range that are undetermined. Some of these unknown differentials are inconsequential because possible error terms are killed by later differentials. Others are inconsequential because they only affect the names (and Adams filtrations) of the elements in the Adams E_∞ -page that detect particular stable homotopy elements. A few of the unknown differentials affect the structure of the stable homotopy groups more significantly. This means that the orders of some of the stable homotopy groups are known only up to factors of 2.

In addition, there are some possible hidden extensions by 2, η , and ν that remain unresolved. Tables 16, 18, and 20 summarize these possibilities. The presence of unknown hidden extensions by 2 means that the group structures of some stable homotopy groups are not known, even though their orders are known.

1.4. Groups of homotopy spheres

An important application of stable homotopy group computations is to the work of Kervaire and Milnor [23] on the classification of smooth structures on spheres in dimensions at least 5. Let Θ_n be the group of h -cobordism classes of homotopy n -spheres. This group classifies the differential structures on S^n for $n \geq 5$. It has a subgroup Θ_n^{bp} , which consists of homotopy spheres that bound parallelizable manifolds. The relation between Θ_n and the stable homotopy group π_n is summarized in Theorem 1.6. See also [32] for a survey on this subject.

THEOREM 1.6. *(Kervaire-Milnor [23]) Suppose that $n \geq 5$.*

(1) The subgroup Θ_n^{bp} is cyclic, and has the following order:

$$|\Theta_n^{bp}| = \begin{cases} 1, & \text{if } n \text{ is even,} \\ 1 \text{ or } 2, & \text{if } n = 4k + 1, \\ b_k, & \text{if } n = 4k - 1. \end{cases}$$

Here b_k is $2^{2k-2}(2^{2k-1} - 1)$ times the numerator of $4B_{2k}/k$, where B_{2k} is the $2k$ th Bernoulli number.

(2) For $n \not\equiv 2 \pmod{4}$, there is an exact sequence

$$0 \longrightarrow \Theta_n^{bp} \longrightarrow \Theta_n \longrightarrow \pi_n/J \longrightarrow 0.$$

Here π_n/J is the cokernel of the J -homomorphism.

(3) For $n \equiv 2 \pmod{4}$, there is an exact sequence

$$0 \longrightarrow \Theta_n^{bp} \longrightarrow \Theta_n \longrightarrow \pi_n/J \xrightarrow{\Phi} \mathbb{Z}/2 \longrightarrow \Theta_{n-1}^{bp} \longrightarrow 0.$$

Here the map Φ is the Kervaire invariant.

The first few values, and then estimates, of the numbers b_k are given by the sequence

$$28, 992, 8128, 261632, 1.45 \times 10^9, 6.71 \times 10^7, 1.94 \times 10^{12}, 7.54 \times 10^{14}, \dots$$

THEOREM 1.7. *The last column of Table 1 describes the groups Θ_n for $n \leq 90$, with the exception of $n = 4$. The underlined symbols denote the contributions from Θ_n^{bp} .*

The cokernel of the J -homomorphism is slightly different than the v_1 -torsion part of π_n at the prime 2. In dimensions $8m + 1$ and $8m + 2$, there are classes detected by $P^m h_1$ and $P^m h_1^2$ in the Adams spectral sequence. These classes are v_1 -periodic, in the sense that they are detected by the $K(1)$ -local sphere. However, they are also in the cokernel of the J -homomorphism.

We restate the following conjecture from [44], which is based on the current knowledge of stable stems and a problem proposed by Milnor [32].

CONJECTURE 1.8. *In dimensions greater than 4, the only spheres with unique smooth structures are S^5 , S^6 , S^{12} , S^{56} , and S^{61} .*

Uniqueness in dimensions 5, 6 and 12 was known to Kervaire and Milnor [23]. Uniqueness in dimension 56 is due to the first author [16], and uniqueness in dimension 61 is due to the second and the third authors [44].

Conjecture 1.8 is equivalent to the claim that the group Θ_n is not of order 1 for dimensions greater than 61. This conjecture has been confirmed in all odd dimensions by the second and the third authors [44] based on the work of Hill, Hopkins, and Ravenel [13], and in even dimensions up to 140 by Behrens, Hill, Hopkins, and Mahowald [4].

1.5. Notation

The cohomology of the Steenrod algebra is highly irregular, so consistent naming systems for elements presents a challenge. A list of multiplicative generators

appears in Table 4. To a large extent, we rely on the traditional names for elements, as used in [7], [16], [37], and elsewhere. However, we have adopted some new conventions in order to partially systematize the names of elements.

First, we use the symbol Δx to indicate an element that is represented by $v_2^4 x$ in the May spectral sequence. This use of Δ is consistent with the role that v_2^4 plays in the homotopy of tmf , where it detects the discriminant element Δ . For example, instead of the traditional symbol r , we use the name Δh_2^2 .

Second, the symbol M indicates the Massey product operator $\langle -, h_0^3, g_2 \rangle$. For example, instead of the traditional symbol B_1 , we use the name Mh_1 .

Similarly, the symbol g indicates the Massey product operator $\langle -, h_1^4, h_4 \rangle$. For example, we write $h_2 g$ for the indecomposable element $\langle h_2, h_1^4, h_4 \rangle$.

Eventually, we encounter elements that neither have traditional names, nor can be named using symbols such as P , Δ , M , and g . In these cases, we use arbitrary names of the form $x_{s,f}$, where s and f are the stem and Adams filtration of the element.

The last column of Table 4 gives alternative names, if any, for each multiplicative generator. These alternative names appear in at least one of [7] [16] [37].

REMARK 1.9. One specific element deserves further discussion. We define τQ_3 to be the unique element such that $h_3 \cdot \tau Q_3 = 0$. This choice is not compatible with the notation of [16]. The element τQ_3 from [16] equals the element $\tau Q_3 + \tau n_1$ in this manuscript.

We shall also extensively study the Adams spectral sequence for the cofiber of τ . See Section 3.1 for more discussion of the names of elements in this spectral sequence, and how they relate to the Adams spectral sequence for the sphere.

Table 1 gives some notation for elements in $\pi_{*,*}$. Many of these names follow standard usage, but we have introduced additional non-standard elements such as κ_1 and $\bar{\kappa}_2$. These elements are defined by the classes in the Adams E_∞ -page that detect them. In some cases, this style of definition leaves indeterminacy because of the presence of elements in the E_∞ -page in higher filtration. In some of these cases, Table 1 provides additional defining information. Beware that this additional defining information does not completely specify a unique element in $\pi_{*,*}$ in all cases. For the purposes of our computations, these remaining indeterminacies are not consequential.

- (1) $C\tau$ represents the cofiber of $\tau : S^{0,-1} \rightarrow S^{0,0}$. We can also write S/τ for this \mathbb{C} -motivic spectrum, but the latter notation is more cumbersome.
- (2) $\text{Ext} = \text{Ext}_{\mathbb{C}}$ is the cohomology of the \mathbb{C} -motivic Steenrod algebra. It is graded in the form (s, f, w) , where s is the stem (i.e., the total degree minus the Adams filtration), f is the Adams filtration (i.e., the homological degree), and w is the motivic weight.
- (3) Ext_{cl} is the cohomology of the classical Steenrod algebra. It is graded in the form (s, f) , where s is the stem (i.e., the total degree minus the Adams filtration), and f is the Adams filtration (i.e., the homological degree).
- (4) $\pi_{*,*}$ is the 2-completed \mathbb{C} -motivic stable homotopy groups.
- (5) $H^*(S; BP)$ is the Adams-Novikov E_2 -page for the classical sphere spectrum, i.e., $\text{Ext}_{BP_* BP}(BP_*, BP_*)$.
- (6) $H^*(S/2; BP)$ is the Adams-Novikov E_2 -page for the classical mod 2 Moore spectrum, i.e., $\text{Ext}_{BP_* BP}(BP_*, BP_*/2)$.

1.6. How to use this manuscript

The manuscript is oriented around a series of tables to be found in Chapter 8. In a sense, the rest of the manuscript consists of detailed arguments for establishing each of the computations listed in the tables. We have attempted to give references and cross-references within these tables, so that the reader can more easily find the specific arguments pertaining to each computation.

We have attempted to make the arguments accessible to users who do not intend to read the manuscript in its entirety. To some extent, with an understanding of how the manuscript is structured, it is possible to extract information about a particular homotopy class in isolation.

We assume that the reader is also referring to the Adams charts in [19] and [17]. These charts describe the same information as the tables, except in graphical form.

This manuscript is very much a sequel to [16]. We will frequently refer to discussions in [16], rather than repeat that same material here in an essentially redundant way. This is especially true for the first parts of Chapters 2, 3, and 4 of [16], which discuss respectively the general properties of Ext, the May spectral sequence, and Massey products; the Adams spectral sequence and Toda brackets; and hidden extensions.

Chapter 2 provides some additional miscellaneous background material not already covered in [16]. Chapter 3 discusses the nature of the machine-generated data that we rely on. In particular, it describes our data on the algebraic Novikov spectral sequence, which is equal to the Adams spectral sequence for the cofiber of τ . Chapter 4 provides some tools for computing Massey products in Ext, and gives some specific computations. Chapter 5 carries out a detailed analysis of Adams differentials. Chapter 6 computes some miscellaneous Toda brackets that are needed for various specific arguments elsewhere. Chapter 7 methodically studies hidden extensions by τ , 2, η , and ν in the E_∞ -page of the \mathbb{C} -motivic Adams spectral sequence. This chapter also gives some information about other miscellaneous hidden extensions. Finally, Chapter 8 includes the tables that summarize the multitude of specific computations that contribute to our study of stable homotopy groups.

1.7. Acknowledgements

We thank Agnes Beaudry, Mark Behrens, Robert Bruner, Robert Burklund, Dexter Chua, Paul Goerss, Jesper Grodal, Lars Hesselholt, Mike Hopkins, Peter May, Haynes Miller, Christian Nassau, Doug Ravenel, and John Rognes for their support and encouragement throughout this project.

CHAPTER 2

Background

2.1. Associated graded objects

DEFINITION 2.1. A filtered object A consists of a finite chain

$$A = F_0A \supseteq F_1A \supseteq F_2A \supseteq \cdots \supseteq F_{p-1}A \supseteq F_pA = 0$$

of inclusions descending from A to 0.

We will only consider finite chains because these are the examples that arise in our Adams spectral sequences. Thus we do not need to refer to “exhaustive” and “Hausdorff” conditions on filtrations, and we avoid subtle convergence issues associated with infinite filtrations.

EXAMPLE 2.2. The \mathbb{C} -motivic stable homotopy group $\pi_{14,8} = \mathbb{Z}/2 \oplus \mathbb{Z}/2$ is a filtered object under the Adams filtration. The generators of this group are σ^2 and κ . The subgroup F_5 is zero, the subgroup $F_3 = F_4$ is generated by κ , and the subgroup $F_0 = F_1 = F_2$ is generated by σ^2 and κ .

DEFINITION 2.3. Let A be a filtered object. The associated graded object $\text{Gr } A$ is

$$\bigoplus_{i=0}^p F_iA/F_{i+1}A.$$

If a is an element of $\text{Gr } A$, then we write $\{a\}$ for the set of elements of A that are detected by a . In general, $\{a\}$ consists of more than one element of A , unless a happens to have maximal filtration. More precisely, the element a is a coset $\alpha + F_{i+1}A$ for some α in A , and $\{a\}$ is another name for this coset. In this situation, we say that a detects α .

In this manuscript, the main example of a filtered object is a \mathbb{C} -motivic homotopy group $\pi_{p,q}$, equipped with its Adams filtration.

EXAMPLE 2.4. Consider the \mathbb{C} -motivic stable homotopy group $\pi_{14,8}$ with its Adams filtration, as described in Example 2.2. The associated graded object is non-trivial only in degrees 2 and 4, and it is generated by h_3^2 and d_0 respectively.

DEFINITION 2.5. Let A and B be filtered objects, perhaps with filtrations of different lengths. A map $f : A \rightarrow B$ is filtration preserving if $f(F_iA)$ is contained in F_iB for all i .

Let $f : A \rightarrow B$ be a filtration preserving map of filtered objects. We write $\text{Gr } f : \text{Gr } A \rightarrow \text{Gr } B$ for the induced map on associated graded objects.

DEFINITION 2.6. Let a and b be elements of $\text{Gr } A$ and $\text{Gr } B$ respectively. We say that b is the (not hidden) value of a under f if $\text{Gr } f(a) = b$.

We say that b is the hidden value of a under f if:

- (1) $\text{Gr } f(a) = 0$.
- (2) there exists an element α of $\{a\}$ in A such that $f(\alpha)$ is contained in $\{b\}$ in B .
- (3) there is no element γ in filtration strictly higher than α such that $f(\gamma)$ is contained in $\{b\}$.

The motivation for condition (3) may not be obvious. The point is to avoid situations in which condition (2) is satisfied trivially. Suppose that there is an element γ such that $f(\gamma)$ is contained in $\{b\}$. Let a be any element of $\text{Gr } A$ whose filtration is strictly less than the filtration of γ . Now let α be any element of $\{a\}$ such that $f(\alpha) = 0$. (It may not be possible to choose such an α in general, but sometimes it is possible.) Then $\alpha + \gamma$ is another element of $\{a\}$ such that $f(\alpha + \gamma)$ is contained in $\{b\}$. Thus f takes some element of $\{a\}$ into $\{b\}$, but only because of the presence of γ . Condition (3) is designed to exclude this situation.

EXAMPLE 2.7. We illustrate the role of condition (3) in Definition 2.6 with a specific example. Consider the map $\eta : \pi_{14,8} \rightarrow \pi_{15,9}$. The associated graded map $\text{Gr}(\eta)$ takes h_3^2 to 0 and takes d_0 to $h_1 d_0$.

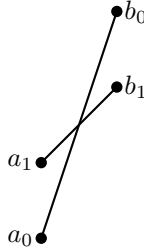
The coset $\{h_3^2\}$ in $\pi_{14,8}$ consists of two elements σ^2 and $\sigma^2 + \kappa$. One of these elements is non-zero after multiplying by η . (In fact, $\eta\sigma^2$ equals zero, and $\eta(\sigma^2 + \kappa) = \eta\kappa$ is non-zero, but that is not relevant here.) Conditions (1) and (2) of Definition 2.6 are satisfied, but condition (3) fails because of the presence of κ in higher filtration.

Suppose that b is the hidden value of a under f . It is typically the case that $f(\alpha)$ is contained in $\{b\}$ for every α in A . However, an even more complicated situation can occur in which this is not true.

Suppose that b_0 is the hidden value of a_0 under f , and suppose that b_1 is the (hidden or not hidden) value of a_1 under f . Moreover, suppose that the filtration of a_0 is strictly lower than the filtration of a_1 , and the filtration of b_0 is strictly greater than the filtration of b_1 . In this situation, we say that the value of a_0 under f crosses the value of a_1 under f .

The terminology arises from the usual graphical calculus, in which elements of higher filtration are drawn above elements of lower filtration, and values of maps are indicated by line segments, as in Figure 1.

FIGURE 1. Crossing values



EXAMPLE 2.8. For any map $X \rightarrow Y$ of \mathbb{C} -motivic spectra, naturality of the Adams spectral sequence induces a filtration preserving map $\pi_{p,q}X \rightarrow \pi_{p,q}Y$. We

are often interested in inclusion $S^{0,0} \rightarrow C\tau$ of the bottom cell into $C\tau$, and in projection $C\tau \rightarrow S^{1,-1}$ from $C\tau$ to the top cell. We also consider the unit map $S^{0,0} \rightarrow mmf$.

2.1.1. Indeterminacy in hidden values. Definition 2.6 allows for the possibility of some essentially redundant cases. In order to avoid this redundancy, we introduce indeterminacy into our definition.

Suppose, as in Definition 2.6, that b is the hidden value of a under f , so there exists some α in $\{a\}$ such that $f(\alpha)$ is contained in $\{b\}$. Suppose also that there is another element a' in $\text{Gr } A$ in degree strictly greater than the degree of a , such that $f(\alpha')$ is contained in $\{b'\}$, where α' is in $\{a'\}$ and b' has the same degree as b . Then $b + b'$ is also a hidden value of a under f , since $\alpha + \alpha'$ is contained in $\{a\}$ and $f(\alpha + \alpha')$ is contained in $\{b + b'\}$. In this case, we say that b' belongs to the target indeterminacy of the hidden value.

EXAMPLE 2.9. Consider the map $\eta : \pi_{63,33} \rightarrow \pi_{64,34}$. The element h_3Q_2 is a hidden value of τh_1H_1 under this map. This hidden value has target indeterminacy generated by $\tau h_1X_2 = h_1 \cdot (\tau X_2 + \tau C')$.

2.1.2. Hidden extensions. Let α be an element of $\pi_{a,b}$. Then multiplication by α induces a filtration preserving map $\pi_{p,q} \rightarrow \pi_{p+a,q+b}$. A hidden value of this map is precisely the same as a hidden extension by α in the sense of [16, Definition 4.2]. For clarity, we repeat the definition here.

DEFINITION 2.10. Let α be an element of $\pi_{*,*}$ that is detected by an element a of the E_∞ -page of the \mathbb{C} -motivic Adams spectral sequence. A hidden extension by α is a pair of elements b and c of E_∞ such that:

- (1) $ab = 0$ in the E_∞ -page.
- (2) There exists an element β of $\{b\}$ such that $\alpha\beta$ is contained in $\{c\}$.
- (3) If there exists an element β' of $\{b'\}$ such that $\alpha\beta'$ is contained in $\{c\}$, then the Adams filtration of b' is less than or equal to the Adams filtration of b .

A crossing value for the map $\alpha : \pi_{p,q} \rightarrow \pi_{p+a,q+b}$ is precisely the same as a crossing extension in the sense of [16, Examples 4.6 and 4.7].

The discussion target indeterminacy applies to the case of hidden extensions. For example, the hidden η extension from h_3Q_2 to τh_1H_1 has target indeterminacy generated by τh_1X_2 .

Typically, there is a symmetry in the presence of hidden extensions, in the following sense. Let α and β be detected by a and b respectively. If there is a hidden α extension from b to c , then usually there is also a hidden β extension from a to c as well. However, this symmetry does not always occur, as the following example demonstrates.

EXAMPLE 2.11. We prove in Lemma 7.156 that $(\sigma^2 + \kappa)\theta_5$ is zero. Therefore, there is no $(\sigma^2 + \kappa)$ extension on h_5^2 . On the other hand, Table 21 shows that $\sigma^2\theta_5$ is non-zero and detected by h_0h_4A . Therefore, there is a hidden θ_5 extension from h_3^2 to h_0h_4A .

In later chapters, we will thoroughly explore hidden extensions by 2, η , and ν . We warn the reader that a complete understanding of such hidden extensions does not necessarily lead to a complete understanding of multiplication by 2, η , and ν in the \mathbb{C} -motivic stable homotopy groups.

For example, in the 45-stem, there exists an element $\theta_{4,5}$ that is detected by $h_2^2 h_5$ such that $4\theta_{4,5}$ is detected by $h_0 h_5 d_0$. This is an example of a hidden 4 extension. However, there is no hidden 2 extension from $h_0 h_2^2 h_5$ to $h_0 h_5 d_0$; condition (3) of Definition 2.6 is not satisfied.

In fact, a complete understanding of *all* hidden extensions leads to a complete understanding of the multiplicative structure of the \mathbb{C} -motivic stable homotopy groups, but the process is perhaps more complicated than expected.

For example, we mentioned in Example 2.7 that either $\eta(\sigma^2 + \kappa)$ or $\eta\sigma^2$ is non-zero, but these cases cannot be distinguished by a study of hidden η extensions. However, we can express that $\eta\sigma^2$ is zero by observing that there is no hidden σ extension from $h_1 h_3$ to $h_1 d_0$.

There are even further complications. For example, the equation $h_2^3 + h_1^2 h_3 = 0$ does not prove that $\nu^3 + \eta^2 \sigma$ equals zero because it could be detected in higher filtration. In fact, this does occur. Toda's relation [38] says that

$$\eta^2 \sigma + \nu^3 = \eta \epsilon,$$

where $\eta \epsilon$ is detected by $h_1 c_0$.

We can express Toda's relation in terms of a "matric hidden extension". We have a map $[\nu \quad \eta] : \pi_{6,4} \oplus \pi_{8,5} \rightarrow \pi_{9,6}$. The associated graded map takes $(h_2^2, h_1 h_3)$ to zero, but $h_1 c_0$ is the hidden value of $(h_2^2, h_1 h_3)$ under this map, in the sense of Definition 2.6.

2.2. Motivic modular forms

Over \mathbb{C} , a "motivic modular forms" spectrum mmf has recently been constructed [10]. From our computational perspective, mmf is a ring spectrum whose cohomology is $A//A(2)$, i.e., the quotient of the \mathbb{C} -motivic Steenrod algebra by the subalgebra generated by Sq^1 , Sq^2 , and Sq^4 . By the usual change-of-rings isomorphism, this implies that the homotopy groups of mmf are computed by an Adams spectral sequence whose E_2 -page is the cohomology of \mathbb{C} -motivic $A(2)$ [14]. The Adams spectral sequence for mmf has been completely computed [17].

By naturality, the unit map $S^{0,0} \rightarrow mmf$ yields a map of Adams spectral sequences. This map allows us to transport information from the thoroughly understood spectral sequence for mmf to the less well understood spectral sequence for $S^{0,0}$. This comparison technique is essential at many points throughout our computations.

We rely on notation from [14] and [17] for the Adams spectral sequence for mmf , except that we use a and n instead of α and ν respectively.

For the most part, the map $\pi_{*,*} \rightarrow \pi_{*,*} mmf$ is detected on Adams E_∞ -pages. However, this map does have some hidden values.

THEOREM 2.12. *Through dimension 90, Table 2 lists all hidden values of the map $\pi_{*,*} \rightarrow \pi_{*,*} mmf$.*

PROOF. Most of these hidden values follow from hidden τ extensions in the Adams spectral sequences for $S^{0,0}$ and for mmf . For example, for $S^{0,0}$, there is a hidden τ extension from $h_1 h_3 g$ to d_0^2 . For mmf , there is a hidden τ extension from cg to d^2 . This implies that cg is the hidden value of $h_1 h_3 g$.

A few cases are slightly more difficult. The hidden values of $\Delta h_1 h_3$ and $h_0 h_5 i$ follow from the Adams-Novikov spectral sequences for $S^{0,0}$ and for mmf . These two values are detected on Adams-Novikov E_∞ -pages in filtration 2.

Next, the hidden value on Ph_2h_5j follows from multiplying the hidden value on h_0h_5i by d_0 . Finally, the hidden values on $\Delta h_1^2h_3$, $h_0h_2h_5i$, and Ph_5j follow from already established hidden values, relying on h_1 extensions and h_2 extensions. \square

REMARK 2.13. Through the 90-stem, there are no crossing values for the map $\pi_{*,*} \rightarrow \pi_{*,*}mmf$. Moreover, in this range, there is only one hidden value that has target indeterminacy. Namely, Δ^2h_2d is the hidden value of Ph_5j , with target indeterminacy generated by $\tau^3\Delta h_1g^2$.

2.3. The cohomology of the \mathbb{C} -motivic Steenrod algebra

We have implemented machine computations of Ext , i.e., the cohomology of the \mathbb{C} -motivic Steenrod algebra, in detail through the 110-stem. We take this computational data for granted. It is depicted graphically in the chart of the E_2 -page shown in [19]. The data is available at [42]. See [43] for a discussion of the implementation.

In addition to the additive structure of Ext , we also have complete information about multiplications by h_0 , h_1 , h_2 , and h_3 . We do not have complete multiplicative information. Occasionally we must deduce some multiplicative information on an ad hoc basis.

Similarly, we do not have systematic machine-generated Massey product information about Ext . We deduce some of the necessary information about Massey products in Chapter 4.

In the classical situation, Bruner has carried out extensive machine computations of the cohomology of the classical Steenrod algebra [7]. This data includes complete primary multiplicative information, but no higher Massey product structure. We rely heavily on this information. Our reliance on this data is so ubiquitous that we will not give repeated citations.

The May spectral sequence is the key tool for a conceptual computation of Ext . See [16] for full details. In this manuscript, we use the May spectral sequence to compute some Massey products that we need for various specific purposes; see Remark 2.26 for more details.

For convenience, we restate the following structural theorem about a portion of $\text{Ext}_{\mathbb{C}}$ [16, Theorem 2.19].

THEOREM 2.14. *There is a highly structured isomorphism from Ext_{cl} to the subalgebra of Ext consisting of elements in degrees (s, f, w) with $s + f - 2w = 0$. This isomorphism takes classical elements of degree (s, f) to motivic elements of degree $(2s + f, f, s + f)$.*

2.4. Toda brackets

Toda brackets are an essential computational tool for understanding stable homotopy groups. Brackets appear throughout the various stages of the computations, including in the analysis of Adams differentials and in the resolution of hidden extensions.

It is well-known that the stable homotopy groups form a ring under the composition product. The higher Toda bracket structure is an extension of this ring structure that is much deeper and more intricate. Our philosophy is that the stable homotopy groups are not really understood until the Toda bracket structure is revealed.

A complete analysis of all Toda brackets (even in a range) is not a practical goal. There are simply too many possibilities to take into account methodically, especially when including matrix Toda brackets (and possibly other more exotic non-linear types of brackets). In practice, we compute only the Toda brackets that we need for our specific computational purposes.

2.4.1. The Moss Convergence Theorem. We next discuss the Moss Convergence Theorem [33], which is the essential tool for computing Toda brackets in stable homotopy groups. In order to make this precise, we must clarify the various types of bracket operations that arise.

First, the Adams E_2 -page has Massey products arising from the fact that it is the homology of the cobar complex, which is a differential graded algebra. We typically refer to these simply as “Massey products”, although we write “Massey products in the E_2 -page” for clarification when necessary.

Next, each higher E_r -page also has Massey products, since it is the homology of the E_{r-1} -page, which is a differential graded algebra. We always refer to these as “Massey products in the E_r -page” to avoid confusion with the more familiar Massey products in the E_2 -page. This type of bracket appears only occasionally throughout the manuscript.

Beware that the higher E_r -pages do not inherit Massey products from the preceding pages. For example, τh_1^2 equals the Massey product $\langle h_0, h_1, h_0 \rangle$ in the E_2 -page. However, in the E_3 -page, the bracket $\langle h_0, h_1, h_0 \rangle$ equals zero, since the product $h_0 h_1$ is already equal to zero in the E_2 -page before taking homology to obtain the E_3 -page.

On the other hand, the Massey product $\langle h_1, h_0, h_3^2 \rangle$ is not a well-defined Massey product in the E_2 -page since $h_0 h_3^2$ is non-zero, while $\langle h_1, h_0, h_3^2 \rangle$ in the E_3 -page equals $h_1 h_4$ because of the differential $d_2(h_4) = h_0 h_3^2$.

Finally, we have Toda brackets in the stable homotopy groups $\pi_{*,*}$. The point of the Moss Convergence Theorem is to relate these various kinds of brackets.

DEFINITION 2.15. Let a and b be elements in the E_r -page of the \mathbb{C} -motivic Adams spectral sequence such that $ab = 0$, and let $n \geq 1$. We call a nonzero differential

$$d_{r+n}x = y$$

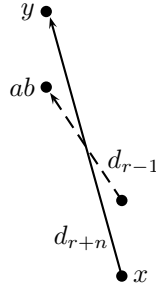
a **crossing differential** for the product ab in the E_r page, if the element y has the same stem and motivic weight as the product $ab = 0$, and the difference between the Adams filtration of y and of ab is strictly greater than 0 but strictly less than $n + 1$.

Figure 2 depicts the situation of a crossing differential in a chart for the E_r -page. Typically, the product ab is zero in the E_r -page because it was hit by a d_{r-1} differential, as shown by the dashed arrow in the figure. However, it may very well be the case that the product ab is already zero in the E_{r-1} -page (or even in an earlier page), in which case the dashed d_{r-1} differential is actually $d_{r-1}(0) = 0$.

THEOREM 2.16 (Moss Convergence Theorem). *Suppose that a , b , and c are permanent cycles in the E_r -page of the \mathbb{C} -motivic Adams spectral sequence that detect homotopy classes α , β , and γ in $\pi_{*,*}$ respectively. Suppose further that*

- (1) *the Massey product $\langle a, b, c \rangle$ is defined in the E_r -page, i.e., $ab = 0$ and $bc = 0$ in the E_r -page.*

FIGURE 2. Crossing differentials



- (2) the Toda bracket $\langle \alpha, \beta, \gamma \rangle$ is defined in $\pi_{*,*}$, i.e., $\alpha\beta = 0$ and $\beta\gamma = 0$.
 (3) there are no crossing differentials for the products ab and bc in the E_r -page.

Then there exists an element e contained in the Massey product $\langle a, b, c \rangle$ in the E_r -page, such that

- (1) the element e is a permanent cycle.
 (2) the element e detects a homotopy class in the Toda bracket $\langle \alpha, \beta, \gamma \rangle$.

REMARK 2.17. The homotopy classes α , β , and γ are usually not unique. The presence of elements in higher Adams filtration implies that a , b , and c detect more than one homotopy class. Moreover, it may be the case that $\langle \alpha, \beta, \gamma \rangle$ is defined for only some choices of α , β , and γ , while the Toda bracket is not defined for other choices.

REMARK 2.18. The Moss Convergence Theorem 2.16 says that a certain Massey product $\langle a, b, c \rangle$ in the E_r -page contains an element with certain properties. The theorem does not claim that every element of $\langle a, b, c \rangle$ has these properties. In the presence of indeterminacies, there can be elements in $\langle a, b, c \rangle$ that do not satisfy the given properties.

REMARK 2.19. Beware that the Toda bracket $\langle \alpha, \beta, \gamma \rangle$ may have non-zero indeterminacy. In this case, we only know that e detects one element of the Toda bracket. Other elements of the Toda bracket could possibly be detected by other elements of the Adams E_∞ -page; these occurrences must be determined by inspection.

REMARK 2.20. In practice, one computes a Toda bracket $\langle \alpha, \beta, \gamma \rangle$ by first studying its corresponding Massey product $\langle a, b, c \rangle$ in a certain page of the Adams spectral sequence. In the case that the Massey product $\langle a, b, c \rangle$ equals zero in the E_r -page in Adams filtration f , the Moss Convergence Theorem 2.16 does not imply that the Toda bracket $\langle \alpha, \beta, \gamma \rangle$ contains zero. Rather, the Toda bracket contains an element (possibly zero) that is detected in Adams filtration at least $f + 1$.

EXAMPLE 2.21. Consider the Toda bracket $\langle \nu, \eta, \nu \rangle$. The elements h_1 and h_2 are permanent cycles that detect η and ν , and the product $\eta\nu$ is zero. We have that $\langle h_2, h_1, h_2 \rangle$ equals h_1h_3 , with no indeterminacy, in the E_2 -page. There are no crossing differentials for the product $h_1h_2 = 0$ in the E_2 -page, so the Moss Convergence Theorem 2.16 implies that h_1h_3 detects a homotopy class in $\langle \nu, \eta, \nu \rangle$.

Note that h_1h_3 detects the homotopy class $\eta\sigma$ because h_3 is a permanent cycle that detects σ . However, we cannot conclude that $\eta\sigma$ is contained in $\langle\nu, \eta, \nu\rangle$. The presence of the permanent cycle c_0 in higher filtration means that h_1h_3 detects both $\eta\sigma$ and $\eta\sigma + \epsilon$, where ϵ is the unique homotopy class that is detected by c_0 . The Moss Convergence Theorem 2.16 implies that either $\eta\sigma$ or $\eta\sigma + \epsilon$ is contained in the Toda bracket $\langle\nu, \eta, \nu\rangle$. In fact, $\eta\sigma + \epsilon$ is contained in the Toda bracket, but determining this requires further analysis.

EXAMPLE 2.22. Consider the Toda bracket $\langle\sigma^2, 2, \eta\rangle$. The elements h_3^2 , h_0 , and h_1 are permanent cycles that detect σ^2 , 2, and η respectively, and the products $2\sigma^2$ and 2η are both zero. Due to the Adams differential $d_2(h_4) = h_0h_3^2$, the Massey product $\langle h_3^2, h_0, h_1 \rangle$ equals h_1h_4 in the E_3 -page, with zero indeterminacy. There are no crossing differentials for the products $h_0h_3^2 = 0$ and $h_0h_1 = 0$ in the E_3 -page. The Moss Convergence Theorem 2.16 implies that h_1h_4 detects a homotopy class in the Toda bracket $\langle\sigma^2, 2, \eta\rangle$.

The element h_3^2 also detects $\sigma^2 + \kappa$, where κ is the unique homotopy class that is detected by d_0 , and the product $2(\sigma^2 + \kappa)$ is zero. The Moss Convergence Theorem 2.16 also implies that h_1h_4 detects a homotopy class in the Toda bracket $\langle\sigma^2 + \kappa, 2, \eta\rangle$.

EXAMPLE 2.23. Consider the Toda bracket $\langle\kappa, 2, \eta\rangle$. The elements d_0 , h_0 , and h_1 are permanent cycles that detect κ , 2, and η respectively, and the products 2κ and 2η are both zero. Due to the Adams differential $d_3(h_0h_4) = h_0d_0$, the Massey product $\langle d_0, h_0, h_1 \rangle$ equals $h_0h_4 \cdot h_1 = 0$ in Adams filtration 3 in the E_4 -page, with zero indeterminacy. There are no crossing differentials for the products $h_0d_0 = 0$ and $h_0h_1 = 0$ in the E_4 -page. The Moss Convergence Theorem 2.16 implies that the Toda bracket $\langle\kappa, 2, \eta\rangle$ either contains zero, or it contains a non-zero element detected in Adams filtration greater than 3.

The only possible detecting element is Pc_0 . There is a hidden η extension from $h_0^3h_4$ to Pc_0 , so Pc_0 detects an element in the indeterminacy of $\langle\kappa, 2, \eta\rangle$. Consequently, the Toda bracket is $\{0, \eta\rho_{15}\}$, where ρ_{15} is detected by $h_0^3h_4$.

EXAMPLE 2.24. The Massey product $\langle h_2, h_3^2, h_0^2 \rangle$ equals $\{f_0, f_0 + h_0^2h_2h_4\}$ in the E_2 -page. The elements h_2 , h_3^2 , and h_0^2 are permanent cycles that detect ν , σ^2 , and 4 respectively, and the products $\nu\sigma^2$ and $4\sigma^2$ are both zero. However, the product $h_0^2h_3^2$ has a crossing differential $d_3(h_0h_4) = h_0d_0$. The Moss Convergence Theorem 2.16 does not apply, and we cannot conclude anything about the Toda bracket $\langle\nu, \sigma^2, 4\rangle$. In particular, we cannot conclude that $\{f_0, f_0 + h_0^2h_2h_4\}$ contains a permanent cycle. In fact, both elements support Adams d_2 differentials.

REMARK 2.25. There is a version of the Moss Convergence Theorem 2.16 for computing fourfold Toda brackets $\langle\alpha, \beta, \gamma, \delta\rangle$ in terms of fourfold Massey products $\langle a, b, c, d \rangle$ in the E_r -page. In this case, the crossing differential condition applies not only to the products ab , bc , and cd , but also to the subbrackets $\langle a, b, c \rangle$ and $\langle b, c, d \rangle$.

REMARK 2.26. Just as the Moss Convergence Theorem 2.16 is the key tool for computing Toda brackets with the Adams spectral sequence, the May Convergence Theorem is the key tool for computing Massey products with the May spectral sequence. The statement of the May Convergence Theorem is entirely analogous to the statement of the Moss Convergence Theorem, with Adams differentials replaced

by May differentials; Adams E_r -pages replaced by May E_r -pages; $\pi_{*,*}$ replaced by Ext; and Toda brackets replaced by Massey products. An analogous crossing differential condition applies. See [16, Section 2.2] [30] for more details. We will use the May Convergence Theorem to compute various Massey products that we need for specific purposes.

2.4.2. Moss's higher Leibniz rule. Occasionally, we will use Moss's higher Leibniz rule [33], which describes how Massey products in the E_r -page interact with the Adams d_r differential. This theorem is a direct generalization of the usual Leibniz rule $d_r(ab) = d_r(a)b + ad_r(b)$ for twofold products.

THEOREM 2.27. [33] *Suppose that a , b , and c are elements in the E_r -page of the \mathbb{C} -motivic Adams spectral sequence such that $ab = 0$, $bc = 0$, $d_r(b)a = 0$, and $d_r(b)c = 0$. Then*

$$d_r\langle a, b, c \rangle \subseteq \langle d_r(a), b, c \rangle + \langle a, d_r(b), c \rangle + \langle a, b, d_r(c) \rangle,$$

where all brackets are computed in the E_r -page.

REMARK 2.28. By the Leibniz rule, the conditions $d_r(b)a = 0$ and $d_r(b)c = 0$ imply that $d_r(a)b = 0$ and $d_r(c)b = 0$. Therefore, all of the Massey products in Theorem 2.27 are well-defined.

REMARK 2.29. The Massey products in Moss's higher Leibniz rule 2.27 may have indeterminacies, so the statement involves an inclusion of sets, rather than an equality.

REMARK 2.30. Beware that Moss's higher Leibniz rule 2.27 cannot be applied to Massey products in the E_r -page to study differentials in higher pages. For example, we cannot use it to compute the d_3 differential on a Massey product in the E_2 -page. In fact, there are versions of the higher Leibniz rule that apply to higher differentials [24, Theorem 8.2] [30, Theorem 4.3], but these results have strong vanishing conditions that often do not hold in practice.

EXAMPLE 2.31. Consider the element $\tau\Delta_1 h_1^2$, which was called G in [37]. Table 4 shows that there is an Adams differential $d_2(\tau\Delta_1 h_1^2) = Mh_1 h_3$, which follows by comparison to $C\tau$. To illustrate Moss's higher Leibniz rule 2.27, we shall give an independent derivation of this differential.

Table 3 shows that $\tau\Delta_1 h_1^2$ equals the Massey product $\langle h_1, h_0, D_1 \rangle$, with no indeterminacy. By Moss's higher Leibniz rule 2.27, the element $d_2(\tau\Delta_1 h_1^2)$ is contained in

$$\langle 0, h_0, D_1 \rangle + \langle h_1, 0, D_1 \rangle + \langle h_1, h_0, d_2(D_1) \rangle.$$

By inspection, the first two terms vanish. Also, Table 4 shows that $d_2(D_1)$ equals $h_0^2 h_3 g_2$.

Therefore, $d_2(\tau\Delta_1 h_1^2)$ is contained in the bracket $\langle h_1, h_0, h_0^2 h_3 g_2 \rangle$, which equals $\langle h_1, h_0, h_0^2 g_2 \rangle h_3$. Finally, Table 3 shows that $\langle h_1, h_0, h_0^2 g_2 \rangle$ equals Mh_1 . This shows that $d_2(\tau\Delta_1 h_1^2)$ equals $Mh_1 h_3$.

EXAMPLE 2.32. Consider the element $\tau e_0 g$ in the Adams E_3 -page. Because of the Adams differential $d_2(e_0) = h_1^2 d_0$, we have that $\tau e_0 g$ equals $\langle d_0, h_1^2, \tau g \rangle$ in the Adams E_3 -page. The higher Leibniz rule 2.27 implies that $d_3(\tau e_0 g)$ is contained in

$$\langle 0, h_1^2, \tau g \rangle + \langle d_0, 0, \tau g \rangle + \langle d_0, h_2^2, 0 \rangle,$$

which equals $\{0, c_0 d_0^2\}$. In this case, the higher Leibniz rule 2.27 does not help to determine the value of $d_3(\tau e_0 g)$ because the indeterminacy is too large. (In fact, $d_3(\tau e_0 g)$ does equal $c_0 d_0^2$, but we need a different argument.)

EXAMPLE 2.33. Lemma 5.29 shows that $d_3(\Delta h_2^2 h_6)$ equals $h_1 h_6 d_0^2$. This argument uses that $\Delta h_2^2 h_6$ equals $\langle \Delta h_2^2, h_5^2, h_0 \rangle$ in the E_3 -page, because of the Adams differential $d_2(h_6) = h_0 h_5^2$.

2.4.3. Shuffling formulas for Toda brackets. Toda brackets satisfy various types of formal relations that we will use extensively. The most important example of such a relation is the shuffle formula

$$\alpha \langle \beta, \gamma, \delta \rangle = \langle \alpha, \beta, \gamma \rangle \delta,$$

which holds whenever both Toda brackets are defined. Note the equality of sets here; the indeterminacies of both expressions are the same.

The following theorem states some formal properties of threefold Toda brackets that we will use later. We apply these results so frequently that we typically use them without further mention.

THEOREM 2.34. *Let $\alpha, \alpha', \beta, \gamma$, and δ be homotopy classes in $\pi_{*,*}$. Each of the following relations involving threefold Toda brackets holds up to a sign, whenever the Toda brackets are defined:*

- (1) $\langle \alpha + \alpha', \beta, \gamma \rangle \subseteq \langle \alpha, \beta, \gamma \rangle + \langle \alpha', \beta, \gamma \rangle$.
- (2) $\langle \alpha, \beta, \gamma \rangle = \langle \gamma, \beta, \alpha \rangle$.
- (3) $\alpha \langle \beta, \gamma, \delta \rangle \subseteq \langle \alpha \beta, \gamma, \delta \rangle$.
- (4) $\langle \alpha \beta, \gamma, \delta \rangle \subseteq \langle \alpha, \beta \gamma, \delta \rangle$.
- (5) $\alpha \langle \beta, \gamma, \delta \rangle = \langle \alpha, \beta, \gamma \rangle \delta$.
- (6) $0 \in \langle \alpha, \beta, \gamma \rangle + \langle \beta, \gamma, \alpha \rangle + \langle \gamma, \alpha, \beta \rangle$.

We next turn our attention to fourfold Toda brackets. New complications arise in this context. If $\alpha\beta = 0$, $\beta\gamma = 0$, $\gamma\delta = 0$, $\langle \alpha, \beta, \gamma \rangle$ contains zero, and $\langle \beta, \gamma, \delta \rangle$ contains zero, then the fourfold bracket $\langle \alpha, \beta, \gamma, \delta \rangle$ is not necessarily defined. Problems can arise when both threefold subbrackets have indeterminacy. See [15] for a careful analysis of this problem in the analogous context of Massey products.

However, when at least one of the threefold subbrackets is strictly zero, then these difficulties vanish. Every fourfold bracket that we use has at least one threefold subbracket that is strictly zero.

Another complication with fourfold Toda brackets lies in the description of the indeterminacy. If at least one threefold subbracket is strictly zero, then the indeterminacy of $\langle \alpha, \beta, \gamma, \delta \rangle$ is the linear span of the sets $\langle \alpha, \beta, \epsilon \rangle$, $\langle \alpha, \epsilon, \delta \rangle$, and $\langle \epsilon, \gamma, \delta \rangle$, where ϵ ranges over all possible values in the appropriate degree for which the Toda bracket is defined.

The following theorem states some formal properties of fourfold Toda brackets that we will use later. We apply these results so frequently that we typically use them without further mention.

THEOREM 2.35. *Let $\alpha, \alpha', \beta, \gamma, \delta$, and ϵ be homotopy classes in $\pi_{*,*}$. Each of the following relations involving fourfold Toda brackets holds up to a sign, whenever the Toda brackets are defined:*

- (1) $\langle \alpha + \alpha', \beta, \gamma, \delta \rangle \subseteq \langle \alpha, \beta, \gamma, \delta \rangle + \langle \alpha', \beta, \gamma, \delta \rangle$.
- (2) $\langle \alpha, \beta, \gamma, \delta \rangle = \langle \delta, \gamma, \beta, \alpha \rangle$.

- (3) $\alpha\langle\beta, \gamma, \delta, \epsilon\rangle \subseteq \langle\alpha\beta, \gamma, \delta, \epsilon\rangle$.
- (4) $\langle\alpha\beta, \gamma, \delta, \epsilon\rangle \subseteq \langle\alpha, \beta\gamma, \delta, \epsilon\rangle$.
- (5) $\alpha\langle\beta, \gamma, \delta, \epsilon\rangle = \langle\alpha, \beta, \gamma, \delta\rangle\epsilon$.
- (6) $\alpha\langle\beta, \gamma, \delta, \epsilon\rangle \subseteq \langle\langle\alpha, \beta, \gamma\rangle, \delta, \epsilon\rangle$.
- (7) $0 \in \langle\langle\alpha, \beta, \gamma\rangle, \delta, \epsilon\rangle + \langle\alpha, \langle\beta, \gamma, \delta\rangle, \epsilon\rangle + \langle\alpha, \beta, \langle\gamma, \delta, \epsilon\rangle\rangle$.

Part (6) of Theorem 2.35 requires some further explanation. In the expression $\langle\langle\alpha, \beta, \gamma\rangle, \delta, \epsilon\rangle$, we have a set $\langle\alpha, \beta, \gamma\rangle$ as the first input to a threefold Toda bracket. The expression $\langle\langle\alpha, \beta, \gamma\rangle, \delta, \epsilon\rangle$ is defined to be the union of all sets of the form $\langle\zeta, \delta, \epsilon\rangle$, where ζ ranges over all elements of $\langle\alpha, \beta, \gamma\rangle$. Part (7) uses the same notational convention.

We will make occasional use of matrix Toda brackets. We will not describe their shuffling properties in detail, except to observe that they obey analogous matrix versions of the properties in Theorems 2.34 and 2.35.

The algebraic Novikov spectral sequence

Consider the cofiber sequence

$$S^{0,-1} \xrightarrow{\tau} S^{0,0} \xrightarrow{i} C\tau \xrightarrow{p} S^{1,-1},$$

where $C\tau$ is the cofiber of τ . The inclusion i of the bottom cell and projection p to the top cell are tools for comparing the \mathbb{C} -motivic Adams spectral sequence for $S^{0,0}$ to the \mathbb{C} -motivic Adams spectral sequence for $C\tau$. In [16], we analyzed both spectral sequences simultaneously, playing the structure of each against the other in order to obtain more detailed information about both. Then we used the structure of the homotopy of $C\tau$ to reverse-engineer the structure of the classical Adams-Novikov spectral sequence.

In this manuscript, we use $C\tau$ in a different, more powerful way, because we have a deeper understanding of the connection between the homotopy of $C\tau$ and the structure of the classical Adams-Novikov spectral sequence. Namely, the \mathbb{C} -motivic spectrum $C\tau$ is an E_∞ -ring spectrum [9] (here E_∞ is used in the naive sense, not in some enriched motivic sense). Moreover, with appropriate finiteness conditions, the homotopy category of $C\tau$ -modules is equivalent to the category of BP_*BP -comodules [11] [27]. By considering endomorphisms of unit objects, this comparison of homotopy categories gives a structured explanation for the identification of the homotopy of $C\tau$ and the classical Adams-Novikov E_2 -page.

From a computational perspective, there is an even better connection. Namely, the algebraic Novikov spectral sequence for computing the Adams-Novikov E_2 -page [34] [31] is identical to the \mathbb{C} -motivic Adams spectral sequence for computing the homotopy of $C\tau$. This rather shocking, and incredibly powerful, identification of spectral sequences allows us to transform purely algebraic computations directly into information about Adams differentials for $C\tau$. Finally, naturality along the inclusion i of the bottom cell and along the projection p to the top cell allows us to deduce information about Adams differentials for $S^{0,0}$.

Due to the large quantity of data, we do not explicitly describe the structure of the Adams spectral sequence for $C\tau$ in this manuscript. We refer the interested reader to the charts in [21], which provide details in a graphical form.

3.1. Naming conventions

Our naming convention for elements of the algebraic Novikov spectral sequence (and for elements of the Adams-Novikov spectral sequence) differs from previous approaches. Our names are chosen to respect the inclusion i of the bottom cell and the projection p to the top cell. Specifically, if x is an element of the \mathbb{C} -motivic Adams E_2 -page for $S^{0,0}$, then we use the same letter x to indicate its image $i_*(x)$ in the Adams E_2 -page for $C\tau$. It is certainly possible that $i_*(x)$ is zero, but we

will only use this convention in cases where $i_*(x)$ is non-zero, i.e., when x is not a multiple of τ .

On the other hand, if x is an element of the \mathbb{C} -motivic Adams E_2 -page for $S^{0,0}$ such that τx is zero, then we use the symbol \bar{x} to indicate an element of $p_*^{-1}(x)$ in the Adams E_2 -page for $C\tau$. There is often more than one possible choice for \bar{x} , and the indeterminacy in this choice equals the image of i_* in the appropriate degree. We will not usually be explicit about these choices. However, potential confusion can arise in this context. For example, it may be the case that one choice of \bar{x} supports an h_1 extension, while another choice of \bar{x} supports an h_2 extension, but there is no possible choice of \bar{x} that simultaneously supports both extensions. (The authors dwell on this point because this precise issue has generated confusion about specific computations.)

3.2. Machine computations

We have analyzed the algebraic Novikov spectral sequence by computer in a large range. Roughly speaking, our algorithm computes a Curtis table for a minimal resolution. Significant effort went into optimizing the linear algebra algorithms to complete the computation in a reasonable amount of time. The data is available at [42]. See [43] for a discussion of the implementation.

Our machine computations give us a full description of the additive structure of the algebraic Novikov E_2 -page, together with all d_r differentials for $r \geq 2$. It thus yields a full description of the additive structure of the algebraic Novikov E_∞ -page.

Moreover, the data also gives full information about multiplication by 2, h_1 , and h_2 in the Adams-Novikov E_2 -page for the classical sphere spectrum, which we denote by $H^*(S; BP)$.

We have also conducted machine computations of the Adams-Novikov E_2 -page for the classical cofiber of 2, which we denote by $H^*(S/2; BP)$. Note that $H^*(S; BP)$ is the homology of a differential graded algebra (i.e., the cobar complex) that is free as a \mathbb{Z}_2 -module. Therefore, $H^*(S/2; BP)$ is the homology of this differential graded algebra modulo 2. We have computed this homology by machine, including full information about multiplication by h_1 , h_2 , and h_3 . These computations are related by a long exact sequence

$$\cdots \longrightarrow H^*(S; BP) \xrightarrow{j} H^*(S/2; BP) \xrightarrow{q} H^*(S; BP) \longrightarrow \cdots$$

Because h_2^2, h_3^2, h_4^2 , and h_5^2 are annihilated by 2 in $H^*(S; BP)$, there are classes $\widetilde{h}_2^2, \widetilde{h}_3^2, \widetilde{h}_4^2$, and \widetilde{h}_5^2 in $H^*(S/2; BP)$ such that $q(\widetilde{h}_i^2)$ equals h_i^2 for $2 \leq i \leq 5$. We also have full information about multiplication by $\widetilde{h}_2^2, \widetilde{h}_3^2, \widetilde{h}_4^2$, and \widetilde{h}_5^2 in $H^*(S/2; BP)$.

This multiplicative information allows us to determine some of the Massey product structure in the Adams-Novikov E_2 -page for the sphere spectrum. There are several cases to consider.

First, let x and y be elements of $H^*(S; BP)$. If the product $j(x)j(y)$ is non-zero in $H^*(S/2; BP)$, then xy must also be non-zero in $H^*(S; BP)$.

In the second case, let x be an element of $H^*(S; BP)$, and let \tilde{y} be an element of $H^*(S/2; BP)$ such that $q(\tilde{y}) = y$. If the product $x \cdot \tilde{y}$ is non-zero in $H^*(S/2; BP)$ and equals $j(z)$ for some z in $H^*(S; BP)$, then z belongs to the Massey product $\langle 2, y, x \rangle$. This follows immediately from the relationship between Massey products and the multiplicative structure of a cofiber, as discussed in [16, Section 3.1.1].

Third, let \tilde{x} and \tilde{y} be elements of $H^*(S/2; BP)$ such that $q(\tilde{x}) = x$, $q(\tilde{y}) = y$, and $q(\tilde{x} \cdot \tilde{y}) = z$. Then z belongs to the Massey product $\langle x, 2, y \rangle$ in $H^*(S; BP)$. This follows immediately from the multiplicative snake lemma 3.3.

EXAMPLE 3.1. Computer data shows that the product $\widetilde{h}_4^2 \cdot \widetilde{h}_5^2$ does not equal zero in $H^*(S/2; BP)$. This implies that the Massey product $\langle h_4^2, 2, h_5^2 \rangle$ does not contain zero in $H^*(S; BP)$, which in turn implies that the Toda bracket $\langle \theta_4, 2, \theta_5 \rangle$ does not contain zero in $\pi_{93,48}$.

REMARK 3.2. let \tilde{x} and \tilde{y} be elements of $H^*(S/2; BP)$ such that $q(\tilde{x})$ and $q(\tilde{y})$ equal x and y , and such that $\tilde{x} \cdot \tilde{y}$ equals $j(z)$ for some z in $H^*(S; BP)$. It appears that z has some relationship to the fourfold Massey product $\langle 2, x, 2, y \rangle$, but we have not made this precise.

LEMMA 3.3 (Multiplicative Snake Lemma). *Let A be a differential graded algebra that has no 2-torsion, and let $H(A)$ be its homology. Also let $H(A/2)$ be the homology of $A/2$, and let $\delta : H(A/2) \rightarrow H(A)$ be the boundary map associated to the short exact sequence*

$$0 \longrightarrow A \xrightarrow{2} A \longrightarrow A/2 \longrightarrow 0.$$

Suppose that a and b are elements of $H(A/2)$ such that $2\delta(a) = 0$ and $2\delta(b) = 0$ in $H(A)$. Then the Massey product $\langle \delta(a), 2, \delta(b) \rangle$ in $H(A)$ contains $\delta(ab)$.

PROOF. We carry out a diagram chase in the spirit of the snake lemma. Write ∂ for the boundary operators in A and $A/2$.

Let x and y be cycles in $A/2$ that represent a and b respectively. Let x' and y' be elements in A that reduce to x and y . Then $\partial x'$ and $\partial y'$ reduce to zero in $A/2$ because x and y are cycles. Therefore, $\partial x' = 2\tilde{x}$ and $\partial y' = 2\tilde{y}$ for some \tilde{x} and \tilde{y} in A .

By definition of the boundary map, $\delta(a)$ and $\delta(b)$ are represented by \tilde{x} and \tilde{y} . By the definition of Massey products, the cycle $\tilde{x}y' + x'\tilde{y}$ is contained in $\langle \delta(a), 2, \delta(b) \rangle$.

Now we compute $\delta(ab)$. Note that $x'y'$ is an element of A that reduces to ab . Then

$$\partial(x'y') = \partial(x')y' + x'\partial(y') = 2(\tilde{x}y' + x'\tilde{y}).$$

This shows that $\delta(ab)$ is represented by $\tilde{x}y' + x'\tilde{y}$. \square

3.3. h_1 -Bockstein spectral sequence

The charts in [21] show graphically the algebraic Novikov spectral sequence, i.e., the Adams spectral sequence for $C\tau$. Essentially all of the information in the charts can be read off from machine-generated data. This includes hidden extensions in the E_∞ -page.

One aspect of these charts requires further explanation. The \mathbb{C} -motivic Adams E_2 -page for $C\tau$ contains a large number of h_1 -periodic elements, i.e., elements that support infinitely many h_1 multiplications. The behavior of these elements is entirely understood [12], at least up to many multiplications by h_1 , i.e., in an h_1 -periodic sense.

On the other hand, it takes some work to “desperiodicize” this information. For example, we can immediately deduce from [12] that $d_2(h_1^k e_0) = h_1^{k+2} d_0$ for large values of k , but that does not necessarily determine the behavior of Adams differentials for small values of k .

The behavior of these elements is a bit subtle in another sense, as illustrated by Example 3.4.

EXAMPLE 3.4. Consider the h_1 -periodic element $\overline{c_0 e_0}$. Machine computations tell us that this element supports a d_2 differential, but there is more than one possible value for $d_2(\overline{c_0 e_0})$ because of the presence of both $h_1^2 \overline{c_0 d_0}$ and $P e_0$.

In fact, $d_2(\overline{c_0 d_0})$ equals $P d_0$, and $d_2(P e_0)$ equals $P h_1^2 d_0$. Therefore, $P e_0 + h_1^2 \overline{c_0 d_0}$ is the only non-zero d_2 cycle, and it follows that $d_2(\overline{c_0 e_0})$ must equal $P e_0 + h_1^2 \overline{c_0 d_0}$.

In higher stems, it becomes more and more difficult to determine the exact values of the Adams d_2 differentials on h_1 -periodic classes. Eventually, these complications become unmanageable because they involve sums of many monomials.

Fortunately, we only need concern ourselves with the Adams d_2 differential in this context. The h_1 -periodic E_3 -page equals the h_1 -periodic E_∞ -page, and the only non-zero classes are well-understood v_1 -periodic families running along the top of the Adams chart.

Our solution to this problem, as usual, is to introduce a filtration that hides the filtration. In this case, we filter by powers of h_1 . The effect is that terms involving higher powers of h_1 are ignored, and the formulas become much more manageable.

This h_1 -Bockstein spectral sequence starts with an E_0 -page, because there are some differentials that do not increase h_1 divisibility. For example, we have Bockstein differentials $d_0(\overline{h_1^2 e_0}) = \overline{h_1^4 d_0}$ and $d_0(\overline{c_0 d_0}) = P d_0$, reflecting the Adams differentials $d_2(\overline{h_1^2 e_0}) = \overline{h_1^4 d_0}$ and $d_2(\overline{c_0 d_0}) = P d_0$.

There are also plenty of higher h_1 -Bockstein differentials, such as $d_2(e_0) = h_1^2 d_0$, and $d_7(e_0^2 g) = M h_1^8$.

REMARK 3.5. Beware that filtering by powers of h_1 changes the multiplicative structure in perhaps unexpected ways. For example, $P h_1$ and d_0 are not h_1 -multiples, so their h_1 -Bockstein filtration is zero. One might expect their product to be $P h_1 d_0$, but the h_1 -Bockstein filtration of this element is 1. Therefore, $P h_1 \cdot d_0$ equals 0 in the h_1 -Bockstein spectral sequence.

But not all $P h_1$ multiplications are trivial in the h_1 -Bockstein spectral sequence. For example, we have $P h_1 \cdot \overline{c_0 d_0} = \overline{P h_1 c_0 d_0}$ because the h_1 -Bockstein filtrations of all three elements are zero.

In Example 3.4, we explained that there is an Adams differential $d_2(\overline{c_0 e_0}) = P e_0 + h_1^2 \overline{c_0 d_0}$. When we throw out higher powers of h_1 , we obtain the h_1 -Bockstein differential $d_0(\overline{c_0 e_0}) = P e_0$. We also have an h_1 -Bockstein differential $d_0(\overline{c_0 d_0}) = P d_0$.

The first four charts in [21] show graphically how this h_1 -Bockstein spectral sequence plays out in practice. The main point is that the h_1 -Bockstein E_∞ -page reveals which (formerly) h_1 -periodic classes contribute to the Adams E_3 -page for $C\tau$.

Massey products

The purpose of this chapter is to provide some general tools, and to give some specific computations, of Massey products in Ext . This material contributes to Table 3, which lists a number of Massey products in Ext that we need for various specific purposes. Most commonly, these Massey products yield information about Toda brackets via the Moss Convergence Theorem 2.16.

We begin with a \mathbb{C} -motivic version of a classical theorem of Adams about symmetric Massey products [1, Lemma 2.5.4].

THEOREM 4.1.

- (1) If h_0x is zero, then $\langle h_0, x, h_0 \rangle$ contains τh_1x .
- (2) If $n \geq 1$ and h_nx is zero, then $\langle h_n, x, h_n \rangle$ contains $h_{n+1}x$.

4.1. The operator g

The projection map $p : A_* \rightarrow A(2)_*$ induces a map $p_* : \text{Ext}_{\mathbb{C}} \rightarrow \text{Ext}_{A(2)}$. Because $\text{Ext}_{A(2)}$ is completely known [14], this map is useful for detecting structure in $\text{Ext}_{\mathbb{C}}$. Proposition 4.2 provides a tool for using p_* to compute certain types of Massey products.

PROPOSITION 4.2. *Let x be an element of $\text{Ext}_{\mathbb{C}}$ such that $h_1^4x = 0$. Then $p_*(\langle h_4, h_1^4, x \rangle)$ equals the element $gp_*(x)$ in $\text{Ext}_{A(2)}$.*

PROOF. The idea of the proof is essentially the same as in [18, Proposition 3.1]. The $\text{Ext}_{\mathbb{C}}$ -module $\text{Ext}_{A(2)}$ is a ‘‘Toda module’’, in the sense that Massey products $\langle x, a, b \rangle$ are defined for all x in $\text{Ext}_{A(2)}$ and all a and b in $\text{Ext}_{\mathbb{C}}$ such that $x \cdot a = 0$ and $ab = 0$. In particular, the bracket $\langle 1, h_4, h_1^4 \rangle$ is defined in $\text{Ext}_{A(2)}$. We wish to compute this bracket.

We use the May Convergence Theorem in order to compute the bracket. The crossing differentials condition on the theorem is satisfied because there are no possible differentials that could interfere.

The key point is the May differential $d_4(b_{21}^2) = h_1^4h_4$. This shows that g is contained in $\langle 1, h_4, h_1^4 \rangle$. Also, the bracket has no indeterminacy by inspection.

Now suppose that x is an element of $\text{Ext}_{\mathbb{C}}$ such that $h_1^4x = 0$. Then

$$p_*(\langle h_4, h_1^4, x \rangle) = 1 \cdot \langle h_4, h_1^4, x \rangle = \langle 1, h_4, h_1^4 \rangle \cdot x = gp_*(x).$$

□

EXAMPLE 4.3. We illustrate the practical usefulness of Proposition 4.2 with a specific example. Consider the Massey product $\langle h_1^3h_4, h_1, h_2 \rangle$. The proposition says that

$$p_*(\langle h_1^3h_4, h_1, h_2 \rangle) = h_2g$$

in $\text{Ext}_{A(2)}$. This implies that $\langle h_1^3h_4, h_1, h_2 \rangle$ equals h_2g in $\text{Ext}_{\mathbb{C}}$.

REMARK 4.4. The Massey product computation in Example 4.3 is in relatively low dimension, and it can be computed using other more direct methods. Table 3 lists additional examples, including some that cannot be determined by more elementary methods.

4.2. The Mahowald operator

We recall some results from [18] about the Mahowald operator. The Mahowald operator is defined to be $Mx = \langle x, h_0^3, g_2 \rangle$ for all x such that h_0^3x equals zero. As always, one must be cautious about indeterminacy in Mx .

There exists a subalgebra B of the \mathbb{C} -motivic Steenrod algebra whose cohomology $\text{Ext}_B(\mathbb{M}_2, \mathbb{M}_2)$ equals $\mathbb{M}_2[v_3] \otimes_{\mathbb{M}_2} \text{Ext}_{A(2)}$. The inclusion of B into the \mathbb{C} -motivic Steenrod algebra induces a map $p_* : \text{Ext}_{\mathbb{C}} \rightarrow \text{Ext}_B$.

PROPOSITION 4.5. [18, Theorem 1.1] *The map $p_* : \text{Ext}_{\mathbb{C}} \rightarrow \text{Ext}_B$ takes Mx to the product $(e_0v_3^2 + h_1^3v_3^3)p_*(x)$, whenever Mx is defined.*

Proposition 4.5 is useful in practice for detecting certain Massey products of the form $\langle x, h_0^3, g_2 \rangle$. For example, if x is an element of $\text{Ext}_{\mathbb{C}}$ such that h_0^3x equals zero and $e_0p_*(x)$ is non-zero in $\text{Ext}_{A(2)}$, then $\langle x, h_0^3, g_2 \rangle$ is non-zero.

EXAMPLE 4.6. Proposition 4.5 shows that $\langle h_1, h_0, h_0^2g_2 \rangle$ is non-zero. There is only one non-zero element in the appropriate degree, so we have identified the Massey product. We give this element the name Mh_1 .

EXAMPLE 4.7. Expanding on Example 4.6, Proposition 4.5 also shows that $\langle Mh_1, h_0, h_0^2g_2 \rangle$ is non-zero. Again, there is only one non-zero element in the appropriate degree, so we have identified the Massey product. We give this element the name M^2h_1 .

4.3. Additional computations

LEMMA 4.8. *The Massey product $\langle h_1^3h_4, h_1, \tau gn \rangle$ equals τg^2n , with indeterminacy generated by $Mh_0h_2^2g$.*

PROOF. We start by analyzing the indeterminacy. The product $Mc_0 \cdot h_1^3h_4$ equals

$$\langle g_2, h_0^3, c_0 \rangle h_1^3h_4 = \langle g_2, h_0^3, h_1^3h_4c_0 \rangle = \langle g_2, h_0^3, h_0h_2 \cdot h_2g \rangle = \langle g_2, h_0^3, h_2g \rangle h_0h_2,$$

which equals $Mh_0h_2^2g$. The equalities hold because the indeterminacies are zero, and the first and last brackets in this computation are given by Table 3. This shows that $Mh_0h_2^2g$ belongs to the indeterminacy.

Table 3 shows that

$$\langle h_2, h_1^3h_4, h_1 \rangle = \langle h_2, h_1, h_1^3h_4 \rangle$$

equals h_2g . Then

$$h_2 \langle h_1^3h_4, h_1, \tau gn \rangle = \langle h_2, h_1^3h_4, h_1 \rangle \tau gn = \tau h_2g^2n.$$

This implies that $\langle h_1^3h_4, h_1, \tau gn \rangle$ contains either τg^2n or $\tau g^2n + \Delta h_3g^2$. However, the shuffle

$$h_1 \langle h_1^3h_4, h_1, \tau gn \rangle = \langle h_1, h_1^3h_4, h_1 \rangle \tau gn = 0$$

eliminates $\tau g^2n + \Delta h_3g^2$. \square

REMARK 4.9. The Massey product of Lemma 4.8 cannot be established with Proposition 4.2 because $p_*(\tau gn) = 0$ in $\text{Ext}_{A(2)}$.

LEMMA 4.10. *The Massey product $\langle \Delta e_1 + C_0, h_1^3, h_1 h_4 \rangle$ equals $(\Delta e_1 + C_0)g$, with no indeterminacy.*

PROOF. Consider the Massey product $\langle \tau(\Delta e_1 + C_0), h_1^4, h_4 \rangle$. By inspection, this Massey product has no indeterminacy. Therefore,

$$\langle \tau(\Delta e_1 + C_0), h_1^4, h_4 \rangle = (\Delta e_1 + C_0) \langle \tau, h_1^4, h_4 \rangle.$$

Table 3 shows that the latter bracket equals τg , so the expression equals $\tau(\Delta e_1 + C_0)g$.

On the other hand, it also equals $\tau \langle \Delta e_1 + C_0, h_1^3, h_1 h_4 \rangle$. Therefore, the bracket $\langle \Delta e_1 + C_0, h_1^3, h_1 h_4 \rangle$ must contain $(\Delta e_1 + C_0)g$. Finally, the indeterminacy can be computed by inspection. \square

LEMMA 4.11. *The Massey product $\langle h_3, p', h_2 \rangle$ equals $h_0 e_2$, with no indeterminacy.*

PROOF. We have

$$\langle h_3, p', h_2 \rangle h_4^2 = h_3 \langle p', h_2, h_4^2 \rangle = p' \langle h_2, h_4^2, h_3 \rangle.$$

Table 3 shows that the last Massey product equals c_2 . Observe that $p'c_2$ equals $h_0 h_4^2 e_2$.

This shows that $\langle h_3, p', h_2 \rangle$ equals either $h_0 e_2$ or $h_0 e_2 + h_6 e_0$. However, shuffle to obtain

$$\langle h_3, p', h_2 \rangle h_1 = h_3 \langle p', h_2, h_1 \rangle,$$

which must equal zero. Since $h_1(h_0 e_2 + h_6 e_0)$ is non-zero, it cannot equal $\langle h_3, p', h_2 \rangle$.

The indeterminacy is zero by inspection. \square

LEMMA 4.12. *The Massey product $\langle \tau^3 g G_0, h_0 h_2, h_2 \rangle$ has indeterminacy generated by $\tau M^2 h_2$, and it either contains zero or $\tau e_0 x_{76,9}$. In particular, it does not contain any linear combination of $\Delta^2 h_1 g_2$ with other elements.*

PROOF. The indeterminacy can be computed by inspection.

The only possible elements in the Massey product $\langle \tau^2 g G_0, h_0 h_2, h_2 \rangle$ are linear combinations of $e_0 x_{76,9}$ and $M^2 h_2$. The inclusion

$$\tau \langle \tau^2 g G_0, h_0 h_2, h_2 \rangle \subseteq \langle \tau^3 g G_0, h_0 h_2, h_2 \rangle$$

gives the desired result. \square

LEMMA 4.13. *The Massey product $\langle h_1^2, h_4^2, h_1^2, h_4^2 \rangle$ equals $\Delta_1 h_3^2$.*

PROOF. Table 3 shows that Δh_2^2 equals the Massey product $\langle h_0^2, h_3^2, h_0^2, h_3^2 \rangle$. Recall the isomorphism between classical Ext groups and \mathbb{C} -motivic Ext groups in degrees satisfying $s + f - 2w = 0$, as described in Theorem 2.14. This shows that $\Delta_1 h_3^2$ equals $\langle h_1^2, h_4^2, h_1^2, h_4^2 \rangle$. \square

Adams differentials

The goal of this chapter is to describe the values of the Adams differentials in the motivic Adams spectral sequence. These values are given in Tables 4, 6, 7, 8, and 9. See also the Adams charts in [19] for a graphical representation of the computations.

5.1. The Adams d_2 differential

THEOREM 5.1. *Table 4 lists the values of the Adams d_2 differential on all multiplicative generators, through the 95-stem, except that:*

- (1) $d_2(D_3^2)$ equals either h_1X_3 or $h_1X_3 + \tau h_1C''$.
- (2) $d_2(x_{94,8})$ might equal $\tau h_1^2x_{91,8}$.
- (3) $d_2(x_{95,7})$ equals either $x_{94,9}$ or $x_{94,9} + \tau d_1H_1$.

PROOF. The fourth column of Table 4 gives information on the proof of each differential. Most follow immediately by comparison to the Adams spectral sequence for $C\tau$. A few additional differentials follow by comparison to the classical Adams spectral sequence for tmf .

If an element is listed in the fourth column of Table 4, then the corresponding differential can be deduced from a straightforward argument using a multiplicative relation. For example, it is possible that $d_2(\Delta h_1h_3)$ equals τd_0e_0 . However, $h_0 \cdot \Delta h_1h_3$ is zero, while $h_0 \cdot \tau d_0e_0$ is non-zero. Therefore, $d_2(\Delta h_1h_3)$ must equal zero.

In some cases, it is necessary to combine these different techniques to establish the differential.

The remaining more difficult computations are carried out in the following lemmas. \square

Table 4 lists all of the multiplicative generators of the Adams E_2 -page through the 95-stem. The third column indicates the value of the d_2 differential, if it is non-zero. A blank entry in the third column indicates that the d_2 differential is zero. The fourth column indicates the proof. A blank entry in the fourth column indicates that there are no possible values for the differential. The fifth column gives alternative names for the element, as used in [37], [7], and [16].

REMARK 5.2. Note that $d_3(x_{94,8})$ is non-zero in the Adams spectral sequence for $C\tau$. Therefore, either $d_2(x_{94,8})$ equals $\tau h_1^2x_{91,8}$, or $d_3(x_{94,8})$ equals $h_1x_{92,10}$. In either case, $x_{94,8}$ does not survive to the E_∞ -page, and either $\tau h_1^2x_{91,8}$ or $h_1x_{92,10}$ is hit by a differential.

LEMMA 5.3. $d_2(\Delta x) = h_0^2B_4 + \tau Mh_1d_0$.

PROOF. We have a differential $d_2(\Delta x) = h_0^2B_4$ in the Adams spectral sequence for $C\tau$. Therefore, $d_2(\Delta x)$ equals either $h_0^2B_4$ or $h_0^2B_4 + \tau Mh_1d_0$.

We have the relation $h_1^2 \cdot \Delta x = Ph_1 \cdot \tau \Delta_1 h_1^2$, so $h_1^2 d_2(\Delta x) = Ph_1 d_2(\tau \Delta_1 h_1^2) = Ph_1 h_3 \cdot Mh_1^3 d_0$. Therefore, $d_2(\Delta x)$ must equal $h_0^2 B_4 + \tau Mh_1 d_0$. \square

REMARK 5.4. The proof of [16, Lemma 3.50] is incorrect. We claimed that $h_1^2 \cdot \Delta x$ equals $h_3 \cdot \Delta^2 h_1 h_3$, when in fact $h_1^2 \cdot \Delta x$ equals $\tau h_3 \cdot \Delta^2 h_1 h_3$.

LEMMA 5.5. $d_2(x_{77,7}) = \tau Mh_1 h_4^2$.

The following proof was suggested to us by Dexter Chua.

PROOF. This follows from the interaction between algebraic squaring operations and classical Adams differentials [6, Theorem 2.2], applied to the element x in the 37-stem. The theorem says that

$$d_* \text{Sq}^2 x = \text{Sq}^3 d_2 x \dot{+} h_0 \text{Sq}^3 x.$$

The notation means that there is an Adams differential on $\text{Sq}^2 x$ hitting either $\text{Sq}^3 d_2 x = 0$ or $h_0 \text{Sq}^3 x$, depending on which element has lower Adams filtration. Therefore $d_2 \text{Sq}^2 x = h_0 \text{Sq}^3 x$.

Next, observe from [8] that $\text{Sq}^3 x = h_0^2 x_{76,6} + \tau^2 d_1 g_2$, so

$$h_0 \text{Sq}^3 x = h_0^3 x_{76,6} + \tau Mh_1 h_4^2.$$

Therefore, there is a d_2 differential whose value is $\tau Mh_1 h_4^2$, and the possibility is that $d_2(x_{77,7})$ equals $\tau Mh_1 h_4^2$. \square

LEMMA 5.6. $d_2(\tau B_5 g) = \tau Mh_0^2 g^2$.

PROOF. We use the Mahowald operator methods of Section 4.2. The map $p_* : \text{Ext}_C \rightarrow \text{Ext}_B$ takes $d_0 \cdot \tau B_5 g$ to $\tau h_0 a g^3 v_3^2$, which is non-zero. We deduce that the product $d_0 \cdot \tau B_5 g$ is non-zero in Ext , and the only possibility is that it equals $\tau M g \cdot h_0 m$.

Now $d_2(\tau M g \cdot h_0 m)$ equals $\tau M g \cdot h_0^2 e_0^2$, which we also know is non-zero since it maps to the non-zero element $\tau h_0^2 \text{deg}^2 v_3^2$ of Ext_B . It follows that $d_2(\tau B_5 g)$ must equal $\tau Mh_0^2 g^2$. \square

5.2. The Adams d_3 differential

THEOREM 5.7. *Table 6 lists some values of the Adams d_3 differential on multiplicative generators. The Adams d_3 differential is zero on all multiplicative generators not listed in the table. The list is complete through the 95-stem, except that:*

- (1) $d_3(h_0^7 h_6)$ equals either $h_0 \cdot \Delta^2 h_3^2$ or $h_0 \cdot \Delta^2 h_3^2 + h_0(\tau^2 M e_0 + h_1 \cdot \Delta x)$.
- (2) $d_3(\Delta^3 h_1 h_3)$ equals either $\tau^4 \Delta h_1 e_0^2 g$ or $\tau^4 \Delta h_1 e_0^2 g + \tau \Delta^2 h_0 d_0 e_0$.
- (3) $d_3(h_2 h_6 g)$ might equal $\tau h_1^3 h_4 Q_3$.
- (4) $d_3(\Delta^3 h_1^2 d_0)$ equals either $\tau^3 \Delta h_1 d_0^2 e_0^2$ or $\tau^3 \Delta h_1 d_0^2 e_0^2 + P \Delta^2 h_0 d_0 e_0$.
- (5) $d_3(x_{94,8})$ might equal $h_1 x_{92,10}$.
- (6) $d_3(\Delta^2 Mh_1)$ might equal $\tau^3 M d_0 e_0^2$.

PROOF. The d_3 differential on many multiplicative generators is zero. A few of these multiplicative generators appear in Table 6 because their proofs require further explanation. For the remaining majority of such multiplicative generators, the d_3 differential is zero because there are no possible non-zero values, because of comparison to the Adams spectral sequence for $C\tau$, or because the element is

already known to be a permanent cycle as shown in Table 5. These cases do not appear in Table 6.

The last column of Table 6 gives information on the proof of each differential. Most follow immediately by comparison to the Adams spectral sequence for $C\tau$. A few additional differentials follow by comparison to the classical Adams spectral sequence for tmf , or by comparison to the \mathbb{C} -motivic Adams spectral sequence for mmf .

If an element is listed in the last column of Table 6, then the corresponding differential can be deduced from a straightforward argument using a multiplicative relation. For example,

$$d_3(h_1 \cdot \tau P d_0 e_0) = P h_1 \cdot d_3(\tau d_0 e_0) = P^2 h_1 c_0 d_0,$$

so $d_3(\tau P d_0 e_0)$ must equal $P^2 c_0 d_0$.

If a d_4 differential is listed in the last column of Table 6, then the corresponding differential is forced by consistency with that later differential. In each case, a d_3 differential on an element x is forced by the existence of a later d_4 differential on τx . For example, Table 7 shows that there is a differential $d_4(\tau^2 e_0 g) = P d_0^2$. Therefore, $\tau e_0 g$ cannot survive to the E_4 -page. It follows that $d_3(\tau e_0 g) = c_0 d_0^2$.

In some cases, it is necessary to combine these different techniques to establish the differential.

The remaining more difficult computations are carried out in the following lemmas. \square

Table 6 lists the multiplicative generators of the Adams E_3 -page through the 95-stem whose d_3 differentials are non-zero, or whose d_3 differentials are zero for non-obvious reasons.

REMARK 5.8. Several of the uncertainties in the values of the d_3 differential are inconsequential because they do not affect the structure of later pages.

- (1) The uncertainty in $d_3(h_0^7 h_6)$ is inconsequential because $d_4(\tau^2 C')$ equals $h_0(\tau^2 M e_0 + h_1 \cdot \Delta x)$.
- (2) The uncertainty in $d_3(\Delta^3 h_1 h_3)$ is inconsequential because $d_3(\tau^2 d_0 B_5) = \Delta^2 h_0 d_0 e_0$.
- (3) The uncertainty in $d_3(\Delta^3 h_1^2 d_0)$ is inconsequential since $d_3(\tau^2 M h_0 d_0 k) = P \Delta^2 h_0 d_0 e_0$.

REMARK 5.9. Note that $d_3(x_{94,8})$ is non-zero in the Adams spectral sequence for $C\tau$. Therefore, either $d_2(x_{94,8})$ equals $\tau h_1^2 x_{91,8}$, or $d_3(x_{94,8})$ equals $h_1 x_{92,10}$. In either case, $x_{94,8}$ does not survive to the E_∞ -page, and either $\tau h_1^2 x_{91,8}$ or $h_1 x_{92,10}$ is hit by a differential.

REMARK 5.10. One other d_3 differential possesses a different kind of uncertainty. We know that either $\tau D'_3$ or $\tau D'_3 + \tau^2 h_2 G_0$ survives to the E_3 -page. In Lemma 5.15, we show that the surviving element supports a d_3 differential hitting $\tau^2 M h_2 g$. This is an uncertainty in the source of the differential, rather than the target, which arises from an uncertainty about the value of $d_2(D'_3)$.

PROPOSITION 5.11. *Some permanent cycles in the \mathbb{C} -motivic Adams spectral sequence are shown in Table 5.*

PROOF. The third column of the table gives information on the proof for each element. If a Toda bracket is given in the third column, then the Moss Convergence

Theorem 2.16 implies that the element must survive to detect that Toda bracket (see Table 11 for more information on how each Toda bracket is computed). If a product is given in the third column, then the element must survive to detect that product (see Table 21 for more information on how each product is computed). In a few cases, the third column refers to a specific lemma that gives a more detailed argument. \square

LEMMA 5.12.

- (1) $d_3(h_2h_5) = \tau h_1d_1$.
- (2) $d_3(Ph_2h_6) = \tau h_1h_4Q_2$.

PROOF. In the Adams spectral sequence for $C\tau$, there is an η extension from h_2h_5 to $\overline{h_1^2d_1}$. The element $\overline{h_1^2d_1}$ maps to $h_1^2d_1$ under projection from $C\tau$ to the top cell, so h_2h_5 must also map non-trivially under projection from $C\tau$ to the top cell. The only possibility is that h_2h_5 maps to h_1d_1 . Therefore, τh_1d_1 must be hit by a differential. This establishes the first differential.

The proof for the second differential is identical, using that there is an η extension from Ph_2h_6 to $\overline{h_1^2h_4Q_2}$ in the Adams spectral sequence for $C\tau$. \square

LEMMA 5.13. $d_3(\tau^2\Delta_1h_1^2) = \tau Mc_0$.

PROOF. The element MP maps to zero under inclusion of the bottom cell into $C\tau$. Therefore, MP is either hit by a differential, or it is the target of a hidden τ extension. There are no possible differentials, so there must be a hidden τ extension. The only possibility is that τMc_0 is zero in the E_∞ -page, and that there is a hidden τ extension from Mc_0 to MP . \square

LEMMA 5.14. $d_3(\tau h_0^3 \cdot \Delta g_2) = \tau^3 \Delta h_2^2 e_0 g$.

PROOF. Table 2 shows that the element h_0h_5i maps to $\Delta^2 h_2^2$ in the Adams spectral sequence for tmf .

Now $\Delta^2 h_2^2 d_0$ is not zero and not divisible by 2 in tmf . Therefore, $\kappa\{h_0h_5i\}$ must be non-zero and not divisible by 2 in $\pi_{68,36}$. The only possibility is that $\kappa\{h_0h_5i\}$ is detected by $Ph_2h_5j = d_0 \cdot h_0h_5i$, and that Ph_2h_5j is not an h_0 multiple in the E_∞ -page. Therefore, $\tau \Delta g_2 \cdot h_0^3$ cannot survive to the E_∞ -page. \square

LEMMA 5.15. *Either $d_3(\tau D'_3) = \tau^2 Mh_2g$ or $d_3(\tau D'_3 + \tau^2 h_2 G_0) = \tau^2 Mh_2g$, depending only on whether $\tau D'_3$ or $\tau D'_3 + \tau^2 h_2 G_0$ survives to the E_3 -page.*

PROOF. Table 11 shows that the Toda bracket $\langle 2, 8\sigma, 2, \sigma^2 \rangle$ contains $\tau\nu\bar{\kappa}$, which is detected by $\tau^2 h_2g$. Table 19 shows that Mh_2 detects $\nu\alpha$ for some α in $\pi_{45,24}$ detected by $h_3^2 h_5$. (Beware that there is a crossing extension, Mh_2 does not detect $\nu\alpha$ for every α that is detected by $h_3^2 h_5$.) It follows that $\tau^2 Mh_2g$ detects $\langle 2, 8\sigma, 2, \sigma^2 \rangle \alpha$.

This expression is contained in $\langle 2, 8\sigma, \langle 2, \sigma^2, \alpha \rangle \rangle$. Lemma 6.12 shows that the inner bracket equals $\{0, 2\tau\bar{\kappa}^3\}$.

The Toda bracket $\langle 2, 8\sigma, 0 \rangle$ in $\pi_{68,36}$ consists entirely of multiples of 2. The Toda bracket $\langle 2, 8\sigma, 2\tau\bar{\kappa}^3 \rangle$ contains $\langle 2, 8\sigma, 2 \rangle \tau\bar{\kappa}^3$. This last expression equals zero because

$$\langle 2, 8\sigma, 2 \rangle = \tau\eta \cdot 8\sigma = 0$$

by Corollary 6.2. Therefore, $\langle 2, 8\sigma, 2\tau\bar{\kappa}^3 \rangle$ equals its indeterminacy, which consists entirely of multiples of 2 in $\pi_{68,36}$.

We conclude that $\tau^2 Mh_2g$ is either hit by a differential, or is the target of a hidden 2 extension. Lemma 7.24 shows that there is no hidden 2 extension from h_3A' to $\tau^2 Mh_2g$, and there are no other possible extensions to $\tau^2 Mh_2g$.

Therefore, $\tau^2 Mh_2g$ must be hit by a differential, and the only possible sources of this differential are $\tau D'_3$ or $\tau D'_3 + \tau^2 h_2G_0$, depending on which element survives to the E_3 -page. \square

LEMMA 5.16. $d_3(\tau^2 Mh_0l) = \Delta^2 h_0 d_0^2$.

PROOF. Table 7 shows that $d_4(\tau^2 d_0 e_0 + h_0^7 h_5)$ equals $P^2 d_0$, so $d_4(\tau^3 Mh_1 d_0 e_0)$ equals $\tau M P^2 h_1 d_0$. We have the relation $h_0 \cdot \tau^2 Mh_0l = \tau^3 Mh_1 d_0 e_0$, but the element $\tau M P^2 h_1 d_0$ is not divisible by h_0 . Therefore, $\tau^2 Mh_0l$ cannot survive to the E_4 -page.

By comparison to the Adams spectral sequence for tmf , the value of $d_3(\tau^2 Mh_0l)$ cannot be $\tau^3 \Delta h_1 e_0^3 + \Delta^2 h_0 d_0^2$ or $\tau^3 \Delta h_1 e_0^3$. The only remaining possibility is that $d_3(\tau^2 Mh_0l)$ equals $\Delta^2 h_0 d_0^2$. \square

LEMMA 5.17. $d_3(h_0^3 x_{78,10}) = \tau^6 e_0 g^3$.

PROOF. Suppose that $h_0^3 x_{78,10}$ were a permanent cycle. Then it would map under inclusion of the bottom cell to the element $h_0^3 x_{78,10}$ in the Adams E_∞ -page for $C\tau$.

There is a hidden ν extension from $h_0^3 x_{78,10}$ to $\Delta^3 h_1^2 h_3$ in the Adams E_∞ -page for $C\tau$. Then $\Delta^3 h_1^2 h_3$ would also have to be in the image of inclusion of the bottom cell. The only possible pre-image is the element $\Delta^3 h_1^2 h_3$ in the Adams spectral sequence for the sphere, but this element does not survive by Lemma 5.44.

By contradiction, we have shown that $h_0^3 x_{78,10}$ must support a differential. The only possibility is that $d_3(h_0^3 x_{78,10})$ equals $\tau^6 e_0 g^3$. \square

LEMMA 5.18. $d_3(x_1) = \tau h_1 m_1$.

PROOF. This follows from the interaction between algebraic squaring operations and classical Adams differentials [6, Theorem 2.2]. The theorem says that

$$d_* \text{Sq}^1 e_1 = \text{Sq}^3 d_3 e_1 + h_1 \text{Sq}^3 e_1.$$

The notation means that there is an Adams differential on $\text{Sq}^1 e_1$ hitting either $\text{Sq}^3 d_3 e_1$ or $h_1 \text{Sq}^3 e_1$, depending on which element has lower Adams filtration. Therefore $d_3 \text{Sq}^1 e_1 = h_1 \text{Sq}^3 e_1$.

Finally, we observe from [8] that $\text{Sq}^1 e_1 = x_1$ and $\text{Sq}^3 e_1 = m_1$. \square

LEMMA 5.19. *The element $\Delta^2 d_1$ is a permanent cycle.*

PROOF. The element $\Delta^2 d_1$ in the Adams E_∞ -page for $C\tau$ must map to zero under the projection from $C\tau$ to the top cell. The only possible value in sufficiently high filtration is $\tau^2 \Delta h_1 e_0^2 g$. However, comparison to mmf shows that this element is not annihilated by τ , and therefore cannot be in the image of projection to the top cell.

Therefore, $\Delta^2 d_1$ must be in the image of the inclusion of the bottom cell into $C\tau$. The element $\Delta^2 d_1$ is the only possible pre-image in the Adams E_∞ -page for the sphere in sufficiently low filtration. \square

LEMMA 5.20. $d_3(\Delta^3 h_1 h_3)$ equals either $\tau^4 \Delta h_1 e_0^2 g$ or $\tau^4 \Delta h_1 e_0^2 g + \tau \Delta^2 h_0 d_0 e_0$.

PROOF. There is a relation $Ph_1 \cdot \Delta^3 h_1 h_3 = \tau \Delta^3 h_1^3 d_0$ in the Adams E_2 -page. Because of the differential $d_2(\Delta^3 h_1 e_0) = \Delta^3 h_1^3 d_0 + \tau^5 e_0^2 g m$, we have the relation $Ph_1 \cdot \Delta^3 h_1 h_3 = \tau^6 e_0^2 g m$ in the E_3 -page.

There is a differential $d_4(\tau^6 e_0^2 g m) = \tau^4 d_0^4 l$. But $\tau^4 d_0^4 l$ is not divisible by Ph_1 , so $\tau^6 e_0^2 g m$ cannot be divisible by Ph_1 in the E_4 -page. Therefore, $d_3(\Delta^3 h_1 h_3)$ must be non-zero.

The same argument shows that $d_3(\Delta^3 h_1 h_3 + \tau^3 d_0 B_5)$ must also be non-zero. Because of Lemma 5.21, the only possibilities are that $d_3(\Delta^3 h_1 h_3)$ equals either $\tau^4 \Delta h_1 e_0^2 g$ or $\tau^4 \Delta h_1 e_0^2 g + \tau \Delta^2 h_0 d_0 e_0$. \square

LEMMA 5.21. $d_3(\tau^2 d_0 B_5) = \Delta^2 h_0 d_0 e_0$.

PROOF. The element $\Delta^2 h_0 d_0 e_0$ is a permanent cycle because there are no possible differentials that it could support. Moreover, it must map to zero under the inclusion of the bottom cell into $C\tau$ because there are no elements in the Adams E_∞ -page for $C\tau$ of sufficiently high filtration. Therefore, $\Delta^2 h_0 d_0 e_0$ is either hit by a differential, or it is the target of a hidden τ extension.

The only possible hidden τ extension has source $h_1^3 x_{76,6}$. However, Table 13 shows that $h_1^3 x_{76,6}$ is in the image of projection from $C\tau$ to the top cell. Therefore, it cannot support a hidden extension.

We now know that $\Delta^2 h_0 d_0 e_0$ must be hit by a differential. Lemma 5.19 rules out one possible source for this differential. The only remaining possibility is that $d_3(\tau^2 d_0 B_5)$ equals $d_5(\Delta^2 d_1)$. \square

LEMMA 5.22.

- (1) $d_3(h_2 h_4 h_6) = 0$.
- (2) $d_3(P^2 h_2 h_6) = 0$.

PROOF. The value of $d_3(h_2 h_4 h_6)$ is not $h_2 h_6 d_0$ nor $h_2 h_6 d_0 + \tau h_1 x_1$ by comparison to the Adams spectral sequence for $C\tau$.

It remains to show that $d_3(h_2 h_4 h_6)$ cannot equal $\tau h_1 x_1$. Suppose that the differential did occur. Then there would be no possible targets for a hidden τ extension on $h_1 x_1$, so the η extension from $h_1 x_1$ to $h_1^2 x_1$ would be detected by projection from $C\tau$ to the top cell. But there is no such η extension in the homotopy groups of $C\tau$. This establishes the first formula.

The proof of the second formula is essentially the same, using that the η extension from $\Delta^2 h_1 d_1$ to $\Delta^2 h_1^2 d_1$ cannot be detected by projection from $C\tau$ to the top cell. \square

LEMMA 5.23. $d_3(\Delta^2 p) = 0$.

PROOF. Suppose that $d_3(\Delta^2 p)$ were equal to $\tau^3 h_1 M e_0^2$. In the Adams E_4 -page, the Massey product $\langle \tau^2 M g, \tau h_1 d_0, d_0 \rangle$ would equal $\tau^2 M \Delta h_2^2 g$, with no indeterminacy, because of the Adams differential $d_3(\Delta h_2^2) = \tau h_1 d_0^2$ and because $d_0 \cdot \Delta^2 p = 0$. By Moss's higher Leibniz rule 2.27, $d_4(\tau^2 M \Delta h_2^2 g)$ would be a linear combination of multiples of $\tau^2 M g$ and d_0 . But Table 7 shows that $d_4(\tau^2 M \Delta h_2^2 g)$ equals $M P \Delta h_0^2 e_0$, which is not such a linear combination in the Adams E_4 -page. \square

LEMMA 5.24. $d_3(\tau h_6 g + \tau h_2 e_2) = 0$.

PROOF. In the Adams E_3 -page, we have the matrix Massey product

$$\tau h_6 g + \tau h_2 e_2 = \left\langle \begin{bmatrix} \tau g & \tau h_2 \end{bmatrix}, \begin{bmatrix} h_5^2 \\ x_1 \end{bmatrix}, h_0 \right\rangle$$

because of the Adams differentials $d_2(h_6) = h_0h_5^2$ and $d_2(e_2) = h_0x_1$, as well as the relation $\tau g \cdot h_5^2 + \tau h_2x_1$ in the Adams E_2 -page. Moss's higher Leibniz rule 2.27 implies that $d_3(\tau h_6g + \tau h_2e_2)$ belongs to

$$\left\langle [0 \ 0], \begin{bmatrix} h_5^2 \\ x_1 \end{bmatrix}, h_0 \right\rangle + \left\langle [\tau g \ \tau h_2], \begin{bmatrix} 0 \\ \tau h_1m_1 \end{bmatrix}, h_0 \right\rangle + \left\langle [\tau g \ \tau h_2], \begin{bmatrix} h_5^2 \\ x_1 \end{bmatrix}, 0 \right\rangle$$

since $d_3(x_1) = \tau h_1m_1$, where the Massey products are formed in the Adams E_3 -page using the d_2 differential. This expression simplifies to $\left\langle [\tau g \ \tau h_2], \begin{bmatrix} 0 \\ \tau h_1m_1 \end{bmatrix}, h_0 \right\rangle$,

which equals $\{0, \tau h_0^2h_4Q_3\}$.

Table 19 shows that there is a hidden ν extension from $h_0^2h_4Q_3$ to $Ph_1x_{76,6}$. The element $\tau Ph_1x_{76,6}$ is non-zero in the Adams E_∞ -page. Therefore, $h_0^2h_4Q_3$ supports a (hidden or not hidden) τ extension whose target is in Adams filtration at most 10. The only possibility is that $\tau h_0^2h_4Q_3$ is non-zero in the Adams E_∞ -page. \square

LEMMA 5.25. $d_3(x_{85,6}) = 0$.

PROOF. Suppose that $d_3(x_{85,6})$ equaled τh_2gD_3 . Then h_2gD_3 could not support a hidden τ extension. The only possible target would be $M\Delta h_1d_0$, but that is eliminated by the hypothetical hidden τ extension on Δh_1j_1 given in Remark 7.6.

Table 19 shows that there is a hidden ν extension from h_2gD_3 to B_6d_1 . This hidden extension would be detected by projection from $C\tau$ to the top cell. But there is no such ν extension in the homotopy of $C\tau$. \square

LEMMA 5.26. $d_3(\tau^2Mh_0d_0k) = P\Delta^2h_0d_0e_0$.

PROOF. Table 7 shows that $d_4(\tau^2Mh_1e_0) = MP^2h_1$. Multiply by τd_0^2 to see that $d_4(\tau^3Mh_1d_0^2e_0) = \tau MP^2h_1d_0^2$. We have the relation $h_2 \cdot \tau^2Mh_0d_0k = \tau^3Mh_1d_0^2e_0$, but $\tau MP^2h_1d_0^2$ is not divisible by h_2 . Therefore, $\tau^2Mh_0d_0k$ cannot survive to the E_4 -page. By comparison to mmf , there is only one possible value for $d_3(\tau^2Mh_0d_0k)$. \square

LEMMA 5.27. $d_3(h_2B_5g) = Mh_1c_0e_0^2$.

PROOF. First observe the relation $d_0 \cdot h_2B_5g = \tau Mh_1e_0g^2$. This relation follows from [18, Theorem 1.1], modulo a possible error term $Ph_1^7h_6c_0e_0$. However, multiplication by h_1 eliminates this possibility.

Table 6 shows that $d_3(\tau e_0g^2) = c_0d_0e_0^2$. Therefore, $d_3(\tau Mh_1e_0g^2)$ equals $Mh_1c_0d_0e_0^2$. Observing that $Mh_1c_0d_0e_0^2$ is in fact non-zero in the Adams E_3 -page, we conclude that $d_3(h_2B_5g)$ must equal $Mh_1c_0e_0^2$. \square

LEMMA 5.28. *The element M^2 is a permanent cycle.*

PROOF. Table 3 shows that the Massey product $\langle Mh_1, h_0, h_0^2g_2 \rangle$ equals M^2h_1 . Therefore, M^2h_1 detects the Toda bracket $\langle \eta\theta_{4,5}, 2, \sigma^2\theta_4 \rangle$. The indeterminacy consists entirely of multiples of $\eta\theta_{4,5}$. The Toda bracket contains $\theta_4\langle \eta\theta_{4,5}, 2, \sigma^2 \rangle$. Now $\langle \eta\theta_{4,5}, 2, \sigma^2 \rangle$ is zero because $\pi_{61,33}$ is zero.

We have now shown that M^2h_1 detects a multiple of η . In fact, it detects a non-zero multiple of η because M^2h_1 cannot be hit by a differential by comparison to the Adams spectral sequence for $C\tau$.

Therefore, there exists a non-zero element of $\pi_{90,48}$ that is detected in Adams filtration at most 12. The only possibility is that M^2 survives. \square

LEMMA 5.29. $d_3(\Delta h_2^2 h_6) = \tau h_1 h_6 d_0^2$.

PROOF. In the Adams E_3 -page, $\Delta h_2^2 h_6$ equals $\langle \Delta h_2^2, h_5^2, h_0 \rangle$, with no indeterminacy, because of the Adams differential $d_2(h_6) = h_0 h_5^2$. Using that $d_3(\Delta h_2^2) = \tau h_1 d_0^2$, Moss's higher Leibniz rule 2.27 implies that $d_3(\Delta h_2^2 h_6)$ is contained in

$$\langle \tau h_1 d_0^2, h_5^2, h_0 \rangle + \langle \Delta h_2^2, 0, h_0 \rangle + \langle \Delta h_2^2, h_5^2, 0 \rangle.$$

All of these brackets have no indeterminacy, and the last two equal zero. The first bracket equals $\tau h_1 h_6 d_0^2$, using the Adams differential $d_2(h_6) = h_0 h_5^2$. \square

LEMMA 5.30. $d_3(P^2 h_6 d_0) = 0$.

PROOF. In the Adams E_3 -page, the element $P^2 h_6 d_0$ equals the Massey product $\langle P^2 d_0, h_5^2, h_0 \rangle$, with no indeterminacy, because of the Adams differential $d_2(h_6) = h_0 h_5^2$. Moss's higher Leibniz rule 2.27 implies that $d_3(P^2 h_6 d_0)$ is a linear combination of multiples of h_0 and of $P^2 d_0$. The only possibility is that $d_3(P^2 h_6 d_0)$ is zero. \square

LEMMA 5.31. $d_3(\tau^2 M P h_0 d_0 j) = P^2 \Delta^2 h_0 d_0^2$.

PROOF. Table 7 shows that $d_4(\tau^2 P d_0 e_0) = P^3 d_0$. Multiplication by $\tau M P h_1$ shows that $d_4(\tau^3 M P^2 h_1 d_0 e_0)$ equals $\tau M P^4 h_1 d_0$. But $\tau^3 M P^2 h_1 d_0 e_0$ equals $h_0 \cdot \tau^2 M P h_0 d_0 j$, while $\tau M P^4 h_1 d_0$ is not divisible by h_0 . Therefore, $\tau^2 M P h_0 d_0 j$ cannot survive to the E_4 -page.

The possible values for $d_3(\tau^2 M P h_0 d_0 j)$ are linear combinations of $P^2 \Delta^2 h_0 d_0^2$ and $\tau^3 P \Delta h_1 d_0^3 e_0$. Comparison to the Adams spectral sequence for tmf shows that the term $\tau^3 P \Delta h_1 d_0^3 e_0$ cannot appear. \square

LEMMA 5.32. *The element $P^3 h_6 c_0$ is a permanent cycle.*

PROOF. Table 17 shows that $P^3 c_0$ detects the product $\eta \rho_{31}$. Using the Moss Convergence Theorem 2.16 and the Adams differential $d_2(h_6) = h_0 h_5^2$, the element $P^3 h_6 c_0$ must survive to detect the Toda bracket $\langle \eta \rho_{31}, 2, \theta_5 \rangle$. \square

REMARK 5.33. We suspect that $P^3 h_6 c_0$ detects the product $\eta_6 \rho_{31}$. However, the argument of Lemma 7.151 cannot be completed because the Toda bracket $\langle \eta \rho_{31}, 2, \theta_5 \rangle$ might have indeterminacy in lower Adams filtration.

5.3. The Adams d_4 differential

THEOREM 5.34. *Table 7 lists some values of the Adams d_4 differential on multiplicative generators. The Adams d_4 differential is zero on all multiplicative generators not listed in the table. The list is complete through the 95-stem, except that:*

- (1) $d_4(\Delta^2 M h_1^2)$ might equal $M P \Delta h_0^2 e_0$.

PROOF. The d_4 differential on many multiplicative generators is zero. A few of these multiplicative generators appear in Table 7 because their proofs require further explanation. For the remaining majority of such multiplicative generators, the d_4 differential is zero because there are no possible non-zero values, or because of comparison to the Adams spectral sequences for $C\tau$, tmf , or mmf . In a few cases, the multiplicative generator is already known to be a permanent cycle as shown in Table 5. These cases do not appear in Table 7.

The last column of Table 7 gives information on the proof of each differential. Most follow immediately by comparison to the Adams spectral sequence for $C\tau$, or by comparison to the classical Adams spectral sequence for tmf , or by comparison to the \mathbb{C} -motivic Adams spectral sequence for mmf .

If an element is listed in the last column of Table 7, then the corresponding differential can be deduced from a straightforward argument using a multiplicative relation. For example,

$$d_4(d_0 \cdot \tau^2 e_0 g^2) = d_4(e_0^2 \cdot \tau^2 e_0 g) = e_0^2 \cdot P d_0^2 = d_0^5,$$

so $d_4(\tau^2 e_0 g^2)$ must equal d_0^4 .

The remaining more difficult computations are carried out in the following lemmas. \square

Table 7 lists the multiplicative generators of the Adams E_4 -page through the 95-stem whose d_4 differentials are non-zero, or whose d_4 differentials are zero for non-obvious reasons.

$$\text{LEMMA 5.35. } d_4(\tau h_1 \cdot \Delta x) = \tau^2 \Delta h_2^2 d_0 e_0.$$

For completeness, we repeat the argument from [44, Remark 11.2].

PROOF. Table 7 shows that $\tau^3 \Delta h_2^2 g^2$ supports a d_4 differential, and Table 5 shows that $\tau \Delta^2 h_1^2 g + \tau^3 \Delta h_2^2 g^2$ is a permanent cycle. Therefore, $\tau \Delta^2 h_1^2 g$ also supports a d_4 differential.

On the other hand, we have

$$h_1 \cdot \tau \Delta^2 h_1 g = P h_1 \cdot \Delta x = \Delta x \langle h_1, h_0^3 h_3, h_0 \rangle.$$

This expression equals $\langle h_1 \cdot \Delta x, h_0^3 h_3, h_0 \rangle$ by inspection of indeterminacies. Therefore, the Toda bracket $\langle \{\tau h_1 \cdot \Delta x\}, 8\sigma, 2 \rangle$ cannot be well-formed, since otherwise it would be detected by $\tau \Delta^2 h_1^2 g$. The only possibility is that $\tau h_1 \cdot \Delta x$ is not a permanent cycle, and the only possible differential is that $d_4(\tau h_1 \cdot \Delta x)$ equals $\tau^2 \Delta h_2^2 d_0 e_0$. \square

$$\text{LEMMA 5.36. } d_4(\Delta^2 h_3^2) = 0.$$

PROOF. Table 8 shows that $d_5(\tau h_1^2 \cdot \Delta x)$ equals $\tau^3 d_0^2 e_0^2$. The element $\tau^3 d_0^2 e_0^2$ is not divisible by h_1 in the E_5 -page, so $\tau h_1^2 \cdot \Delta x$ cannot be divisible by h_1 in the E_4 -page.

If $d_4(\Delta^2 h_3^2)$ equaled $\tau^2 \Delta h_2^2 d_0 e_0$, then $\Delta^2 h_3^2 + \tau h_1 \cdot \Delta x$ would survive to the E_5 -page, and $\tau^2 h_1^2 \cdot \Delta x$ would be divisible by h_1 in the E_5 -page. \square

$$\text{LEMMA 5.37. } d_4(\tau X_2) = \tau M h_2 d_0.$$

PROOF. Table 7 shows that $d_4(C')$ equals $M h_2 d_0$. Therefore, either τX_2 or $\tau X_2 + \tau C'$ is non-zero on the E_∞ -page. The inclusion of the bottom cell into $C\tau$ takes this element to $\overline{h_5 d_0 e_0}$.

In the homotopy of $C\tau$, there is a ν extension from $\overline{h_5 d_0 e_0}$ to τB_5 , and inclusion of the bottom cell into $C\tau$ takes $\tau h_2 C'$ to τB_5 .

It follows that there must be a ν extension with target $\tau h_2 C'$. The only possibility is that $\tau X_2 + \tau C'$ is non-zero on the E_∞ -page, and therefore $d_4(\tau X_2)$ equals $d_4(\tau C')$. \square

$$\text{LEMMA 5.38. } d_4(h_0 d_2) = X_3.$$

PROOF. The element X_3 is a permanent cycle. The only possible target for a differential is $\tau^2 d_0 e_0 m$, but this is ruled out by comparison to tmf .

The element X_3 must map to zero under the inclusion of the bottom cell into $C\tau$. Therefore, X_3 is the target of a hidden τ extension, or it is hit by a differential.

The only possible hidden τ extension would have source $\overline{h_1 \cdot \Delta_1 h_3^2}$. In $C\tau$, there is an η extension from $h_0 d_2$ to $\overline{h_1^2 \cdot \Delta_1 h_3^2}$. Since $\overline{h_1^2 \Delta_1 h_3^2}$ maps non-trivially (to $h_1^2 \cdot \Delta_1 h_3^2$) under projection to the top cell of $C\tau$, it follows that $h_0 d_2$ also maps non-trivially under projection. The only possibility is that $h_0 d_2$ maps to $h_1 \cdot \Delta_1 h_3^2$, and therefore $\overline{h_1 \cdot \Delta_1 h_3^2}$ does not support a hidden τ extension.

Therefore, X_3 must be hit by a differential, and there is just one possibility. \square

LEMMA 5.39. $d_4(Mh_2g) = 0$.

PROOF. Table 3 shows that the Massey product $\langle h_2g, h_0^3, g_2 \rangle$ equals Mh_2g . The Moss Convergence Theorem 2.16 shows that Mh_2g must survive to detect the Toda bracket $\langle \{h_2g\}, 8, \overline{\kappa}_2 \rangle$. \square

LEMMA 5.40. $d_4(h_2^2 G_0) = \tau g^2 n$.

PROOF. Table 15 shows that there is a hidden 2 extension from $h_0 h_3 g_2$ to τgn . Therefore, τgn detects $4\sigma \overline{\kappa}_2$.

Table 3 shows that $\langle h_1^3 h_4, h_1, \tau gn \rangle$ consists of the two elements $\tau g^2 n$ and $\tau g^2 n + Mh_2g \cdot h_0 h_2$. Then the Toda bracket $\langle \eta^2 \eta_4, \eta, 4\sigma \overline{\kappa}_2 \rangle$ is detected by either $\tau g^2 n$ or $\tau g^2 n + Mh_2g \cdot h_0 h_2$. But $Mh_2g \cdot h_0 h_2$ is hit by an Adams d_2 differential, so $\tau g^2 n$ detects the Toda bracket.

The Toda bracket has no indeterminacy, so it equals $\langle \eta^2 \eta_4, \eta, 2 \rangle 2\sigma \overline{\kappa}_2$. This last expression must be zero.

We have shown that $\tau g^2 n$ must be hit by some differential. The only possibility is that $d_4(h_2^2 G_0) = \tau g^2 n$. \square

LEMMA 5.41. $d_4(\Delta h_0^2 h_3 g_2) = \tau Mh_1 d_0^2$.

PROOF. Table 8 shows that $d_5(A') = \tau Mh_1 d_0$. Now $d_0 A'$ is zero in the E_5 -page, so $\tau Mh_1 d_0^2$ must also be zero in the E_5 -page. \square

LEMMA 5.42. $d_4(\Delta^2 h_1 h_3 g) = \tau \Delta h_2^2 d_0^2 e_0$.

PROOF. Table 11 shows that the element $\Delta h_2^2 d_0^2 e_0$ detects the Toda bracket $\langle \tau \eta \kappa \overline{\kappa}^2, \eta, \eta^2 \eta_4 \rangle$. Now shuffle to obtain

$$\tau \langle \tau \eta \kappa \overline{\kappa}^2, \eta, \eta^2 \eta_4 \rangle = \langle \tau, \tau \eta \kappa \overline{\kappa}^2, \eta \rangle \eta^2 \eta_4.$$

Table 11 shows that $\langle \tau, \tau \eta \kappa \overline{\kappa}^2, \eta \rangle$ is detected by $h_0 h_2 h_5 i$. It follows that the expression $\langle \tau, \tau \eta \kappa \overline{\kappa}^2, \eta \rangle \eta^2 \eta_4$ is zero, so $\tau \Delta h_2^2 d_0^2 e_0$ must be hit by some differential. The only possibility is that $d_4(\Delta^2 h_1 h_3 g)$ equals $\tau \Delta h_2^2 d_0^2 e_0$. \square

LEMMA 5.43. $d_4(h_0 e_2) = \tau h_1^3 x_{76,6}$.

PROOF. Table 21 shows that $\sigma^2 \theta_5$ is detected by $h_0 h_4 A$. Since $\nu \sigma = 0$, the element $h_0 h_2 h_4 A = \tau h_1^3 x_{76,6}$ must be hit by a differential. The only possibility is that $d_4(h_0 e_2)$ equals $\tau h_1^3 x_{76,6}$. \square

LEMMA 5.44. $d_4(\Delta^3 h_1^2 h_3) = \tau^4 d_0 e_0^2 l$.

PROOF. Table 17 shows that there is a hidden η extension from $\tau^2\Delta h_1g^2$ to $\tau^2d_0e_0m$. Therefore, there is also a hidden η extension from $\tau^2\Delta h_1e_0^2g$ to $\tau^2d_0e_0^2l$.

Also, $\tau^2\Delta h_1e_0^2g$ detects an element in $\pi_{79,43}$ that is annihilated by τ^2 . Therefore, $\tau^4d_0e_0^2l$ must be hit by some differential. Moreover, comparison to *mmf* shows that $\tau^3d_0e_0^2l$ is not hit by a differential.

The hidden η extension from $\tau^3\Delta h_1e_0^2g$ to $\tau^3d_0e_0^2l$ is detected by projection from $C\tau$ to the top cell. The only possibility is that this hidden η extension is the image of the h_1 extension from $\Delta^3h_1h_3$ to $\Delta^3h_1^2h_3$ in the Adams E_∞ -page for $C\tau$.

Therefore, $\Delta^3h_1^2h_3$ maps non-trivially under projection from $C\tau$ to the top cell. Consequently, $\Delta^3h_1^2h_3$ cannot be a permanent cycle in the Adams spectral sequence for the sphere. \square

LEMMA 5.45. $d_4(\Delta j_1) = \tau Mh_0e_0g$.

PROOF. Otherwise, both Δj_1 and $\tau gC'$ would survive to the E_∞ -page, and neither could be the target of a hidden τ extension. They would both map non-trivially under inclusion of the bottom cell into $C\tau$. But there are not enough elements in $\pi_{83,45}C\tau$ for this to occur. \square

LEMMA 5.46. *The element τh_1f_2 is a permanent cycle.*

PROOF. Let α be an element of $\pi_{66,35}$ that is detected by $\tau h_2C'$. Then $\nu\alpha$ is detected by $\tau h_2^2C'$, and $\tau\nu\alpha$ is zero.

Let $\bar{\alpha}$ be an element of $\pi_{70,36}C\tau$ that is detected by $h_0^2h_3h_6$. Projection from $C\tau$ to the top cell takes $\bar{\alpha}$ to $\nu\alpha$. Moreover, in the homotopy of $C\tau$, the Toda bracket $\langle 2, \sigma^2, \bar{\alpha} \rangle$ is detected by $h_0^3c_3$.

Now projection from $C\tau$ to the top cell takes $\langle 2, \sigma^2, \bar{\alpha} \rangle$ to $\langle 2, \sigma^2, \nu\alpha \rangle$, which equals zero by Lemma 6.23. Therefore, $h_0^3c_3$ maps to zero under projection to the top cell of $C\tau$, so it must be in the image of inclusion of the bottom cell. The only possibility is that τh_1f_2 survives and maps to $h_0^3c_3$ under inclusion of the bottom cell. \square

REMARK 5.47. In the proof of Lemma 5.46, we have used that $d_5(\tau p_1 + h_0^2h_3h_6)$ equals $\tau^2h_2^2C'$ in order to conclude that $\tau\nu\alpha$ is zero. This differential depends on work in preparation [5].

However, we can also prove Lemma 5.46 independently of [5]. Lemma 5.57 shows that the other possible value of $d_5(\tau p_1 + h_0^2h_3h_6)$ is $\tau^2h_2^2C' + \tau h_3(\Delta e_1 + C_0)$. In this case, let β be an element of $\pi_{62,33}$ that is detected by $\Delta e_1 + C_0$. Then $\nu\alpha + \sigma\beta$ is detected by $\tau h_2^2C' + h_3(\Delta e_1 + C_0)$, and $\tau(\nu\alpha + \sigma\beta)$ is zero.

Projection from $C\tau$ to the top cell takes $\bar{\alpha}$ to $\nu\alpha + \sigma\beta$, and takes $\langle 2, \sigma^2, \bar{\alpha} \rangle$ to $\langle 2, \sigma^2, \nu\alpha + \sigma\beta \rangle$, which equals zero by Lemmas 6.23 and 6.24. As in the proof of Lemma 5.46, the only possibility is that τh_1f_2 survives and maps to $h_0^3c_3$ under inclusion of the bottom cell of $C\tau$.

LEMMA 5.48. $d_4(h_1c_3) = \tau h_0h_2h_4Q_3$.

PROOF. Lemma 7.155 shows that there exists an element α in $\pi_{67,36}$ that is detected by $h_0Q_3 + h_0n_1$ such that $\tau\nu\alpha$ equals $(\eta\sigma + \epsilon)\theta_5$.

Table 11 shows that the Toda bracket $\langle \nu, \sigma, 2\sigma \rangle$ is detected by h_2h_4 , so the element $\tau h_0h_2h_4Q_3$ detects $\tau\alpha\langle \nu, \sigma, 2\sigma \rangle$, which is contained in $\langle \tau\nu\alpha, \sigma, 2\sigma \rangle$. The indeterminacy in these expressions is zero because $\tau\nu\alpha \cdot \pi_{15,8}$ and $2\sigma \cdot \pi_{78,41}$ are both zero.

We now know that $\tau h_0 h_2 h_4 Q_3$ detects the Toda bracket $\langle (\epsilon + \eta\sigma)\theta_5, \sigma, 2\sigma \rangle$. This bracket contains $\theta_5 \langle \epsilon + \eta\sigma, \sigma, 2\sigma \rangle$. Lemma 6.6 shows that the bracket $\langle \epsilon + \eta\sigma, \sigma, 2\sigma \rangle$ contains 0, so $\theta_5 \langle \epsilon + \eta\sigma, \sigma, 2\sigma \rangle$ equals zero.

Finally, we have shown that $\tau h_0 h_2 h_4 Q_3$ detects zero, so it must be hit by some differential. \square

LEMMA 5.49. $d_4(x_{87,7}) = 0$.

PROOF. Consider the exact sequence

$$\pi_{87,45} \rightarrow \pi_{87,45} C\tau \rightarrow \pi_{86,46}.$$

The middle term $\pi_{87,45} C\tau$ is isomorphic to $(\mathbb{Z}/2)^4$. The elements of $\pi_{87,45}$ that are not divisible by τ are detected by $P^2 h_6 c_0$, and possibly $x_{87,7}$ and $\tau \Delta h_1 H_1$. On the other hand, the elements of $\pi_{86,46}$ that are annihilated by τ are detected by $\tau^3 \Delta c_0 e_0^2 g$ and possibly $M \Delta h_0^2 e_0$.

In order for the possibility $M \Delta h_0^2 e_0$ to occur, either $x_{87,7}$ or $\tau \Delta h_1 H_1$ would have to support a differential hitting $\tau M \Delta h_0^2 e_0$, in which case one of those possibilities could not occur.

If $d_4(x_{87,7})$ equaled $\tau^3 g G_0$, then there would not be enough elements to make the above sequence exact. \square

LEMMA 5.50. $d_4(\tau \Delta h_1 H_1) = 0$.

PROOF. The element $\Delta^2 h_2^2 d_1$ is a permanent cycle that cannot be hit by any differential because $h_2 \cdot \Delta^2 h_2^2 d_1$ cannot be hit by a differential. The element $\Delta^2 h_2^2 d_1$ cannot be in the image of projection from $C\tau$ to the top cell, and it cannot support a hidden τ extension. Therefore, $\tau \Delta^2 h_2^2 d_1$ cannot be hit by a differential. \square

LEMMA 5.51. $d_4(\tau h_2 B_5 g) = M h_1 d_0^3$.

PROOF. Table 21 shows that $M d_0$ detects $\kappa \theta_{4,5}$. Therefore, $M d_0^3$ detects $\kappa^3 \theta_{4,5}$, which equals $\eta^2 \bar{\kappa}^2 \theta_{4,5}$ because Table 17 shows that there is a hidden η extension from $\tau^2 h_1 g^2$ to d_0^3 .

Now $\eta^2 \bar{\kappa}^2 \theta_{4,5}$ is zero because $\eta^2 \bar{\kappa} \theta_{4,5}$ is zero. Therefore, $M d_0^3$ and $M h_1 d_0^3$ must both be hit by differentials.

There are several possible differentials that can hit $M h_1 d_0^3$. The element $h_1 x_{88,10}$ cannot be the source of this differential because Table 5 shows that $x_{88,10}$ is a permanent cycle. The element $\tau h_2^2 g C'$ cannot be the source of the differential because $h_2^2 g C'$ is a permanent cycle by comparison to mmf . The element $\Delta h_1 g_2 g$ cannot be the source because it equals $h_3(\Delta e_1 + C_0)g$. The only remaining possibility is that $d_4(\tau h_2 B_5 g)$ equals $M h_1 d_0^3$. \square

LEMMA 5.52. $d_4(\Delta h_2^2 A') = 0$.

PROOF. In the Adams E_4 -page, the element $\Delta h_2^2 A'$ equals the Massey product $\langle A', h_1, \tau d_0^2 \rangle$, with no indeterminacy because of the Adams differential $d_3(\Delta h_2^2) = \tau h_1 d_0^2$. Moss's higher Leibniz rule 2.27 implies that $d_4(\Delta h_2^2 A')$ is contained in

$$\langle 0, h_1, \tau d_0^2 \rangle + \langle A', 0, \tau d_0^2 \rangle + \langle A', h_1, 0 \rangle,$$

so it is a linear combination of multiples of A' and τd_0^2 . The only possibility is that $d_4(\Delta h_2^2 A')$ is zero. \square

LEMMA 5.53. $d_4(h_4^2 h_6) = h_0^3 g_3$.

PROOF. By comparison to the Adams spectral sequence for $C\tau$, the value of $d_4(h_4^2 h_6)$ is either $h_0^3 g_3$ or $h_0^3 g_3 + \tau h_1 h_4^2 D_3$.

Table 21 shows that $h_0^2 g_3$ detects the product $\theta_4 \theta_5$. Since $2\theta_4 \theta_5$ equals zero, $h_0^3 g_3$ must be hit by a differential. \square

5.4. The Adams d_5 differential

THEOREM 5.54. *Table 8 lists some values of the Adams d_5 differential on multiplicative generators. The Adams d_5 differential is zero on all multiplicative generators not listed in the table. The list is complete through the 95-stem, except that:*

- (1) $d_5(\Delta^2 g_2)$ might equal $\tau^2 \Delta^2 h_2 g^2$.

PROOF. The d_5 differential on many multiplicative generators is zero. For the majority of such multiplicative generators, the d_5 differential is zero because there are no possible non-zero values, or by comparison to the Adams spectral sequence for $C\tau$, or by comparison to *tmf* or *mmf*. In a few cases, the multiplicative generator is already known to be a permanent cycle; $h_1 h_6$ is one such example. A few additional cases appear in Table 8 because their proofs require further explanation.

The last column of Table 8 gives information on the proof of each differential. Many computations follow immediately by comparison to the Adams spectral sequence for $C\tau$.

If an element is listed in the last column of Table 8, then the corresponding differential can be deduced from a straightforward argument using a multiplicative relation. For example,

$$d_5(\tau \cdot gA') = d_5(\tau g \cdot A') = \tau g \cdot \tau M h_1 d_0 = \tau^2 M h_1 e_0^2,$$

so $d_5(gA')$ must equal $\tau M h_1 e_0^2$.

A few of the more difficult computations appear in [5]. The remaining more difficult computations are carried out in the following lemmas. \square

Table 8 lists the multiplicative generators of the Adams E_5 -page through the 95-stem whose d_5 differentials are non-zero, or whose d_5 differentials are zero for non-obvious reasons.

$$\text{LEMMA 5.55. } d_5(\tau h_1^2 \cdot \Delta x) = \tau^3 d_0^2 e_0^2.$$

PROOF. The element $\tau^2 d_0^2 e_0^2$ cannot be hit by a differential. There is a hidden η extension from $\tau \Delta h_2^2 d_0 e_0$ to $\tau^2 d_0^2 e_0^2$ because of the hidden τ extensions from $\tau h_1 g^3 + h_1^5 h_5 c_0 e_0$ to $\Delta h_2^2 d_0 e_0$ and from $h_1^6 h_5 c_0 e_0$ to $d_0^2 e_0^2$. This shows that $\tau^3 d_0^2 e_0^2$ must be hit by some differential.

This hidden η extension is detected by projection from $C\tau$ to the top cell. Since $\overline{Ph_5 c_0 d_0}$ in $C\tau$ maps to $\tau \Delta h_2^2 d_0 e_0$ under projection to the top cell, it follows that $\overline{Ph_1 h_5 c_0 d_0}$ in $C\tau$ maps to $\tau^2 d_0^2 e_0^2$ under projection to the top cell.

If $\tau h_1^2 \cdot \Delta x$ survived, then it could not be the target of a hidden τ extension and it could not be hit by a differential. Also, it could not map non-trivially under inclusion of the bottom cell into $C\tau$, since the only possible value $\overline{Ph_1 h_5 c_0 d_0}$ has already been accounted for in the previous paragraph. \square

$$\text{LEMMA 5.56. } d_5(h_5 d_0 i) = \tau \Delta h_1 d_0^3.$$

PROOF. We showed in Lemma 5.14 that $Ph_2 h_5 j$ cannot be divisible by h_0 in the E_∞ -page. Therefore, $h_5 d_0 i$ must support a differential. \square

LEMMA 5.57. $d_5(\tau p_1 + h_0^2 h_3 h_6)$ equals either $\tau^2 h_2^2 C'$ or $\tau^2 h_2^2 C' + \tau h_3(\Delta e_1 + C_0)$.

PROOF. Projection to the top cell of $C\tau$ takes $h_4 D_2$ to $\tau^3 d_1 g^2$. Moreover, there is a ν extension in the homotopy of $C\tau$ from $h_0^2 h_3 h_6$ to $h_4 D_2$. Therefore, this ν extension must be in the image of projection to the top cell.

Table 19 shows that there is a hidden ν extension from $\tau h_2^2 C'$ to $\tau^3 d_1 g^2$. Therefore, either $\tau h_2^2 C'$ or $\tau h_2^2 C' + h_3(\Delta e_1 + C_0)$ is in the image of projection to the top cell, so $\tau^2 h_2^2 C'$ or $\tau^2 h_2^2 C' + \tau h_2(\Delta e_1 + C_0)$ is hit by a differential. The element $\tau p_1 + h_0^2 h_3 h_6$ is the only possible source for this differential. \square

LEMMA 5.58. $d_5(h_1 x_{71,6}) = 0$.

PROOF. Table 14 shows that there is a hidden τ extension from $Mh_1^2 h_3 g$ to $Mh_1 d_0^2$. Therefore, $Mh_2^2 g$ must also support a τ extension. This shows that $\tau Mh_2^2 g$ cannot be the target of a differential. \square

LEMMA 5.59. $d_5(h_4 D_2) = \tau^4 d_1 g^2$.

PROOF. Suppose for sake of contradiction that $h_4 D_2$ survived, and let α be an element of $\pi_{73,38}$ that is detected by it. Table 14 shows that there is a hidden τ extension from $h_1^2 h_6 c_0$ to $h_0 h_4 D_2$. Therefore, $h_0 h_4 D_2$ detects both 2α and $\tau\eta\epsilon\eta_6$. However, it is possible that the difference between these two elements is detected by $\tau^2 M d_0^2$ or by $\tau^3 \Delta h_1 d_0 e_0^2$. We will handle of each of these cases.

First, suppose that 2α equals $\tau\eta\epsilon\eta_6$. Then the Toda bracket

$$\left\langle \eta, \left[\begin{array}{cc} 2 & \tau\eta\epsilon \end{array} \right], \left[\begin{array}{c} \alpha \\ \eta_6 \end{array} \right] \right\rangle$$

is well-defined. Inclusion of the bottom cell into $C\tau$ takes this bracket to

$$\left\langle \eta, \left[\begin{array}{cc} 2 & 0 \end{array} \right], \left[\begin{array}{c} \alpha \\ \eta_6 \end{array} \right] \right\rangle = \langle \eta, 2, \alpha \rangle,$$

so $\langle \eta, 2, \alpha \rangle$ is in the image of inclusion of the bottom cell.

On the other hand, in the homotopy of $C\tau$, the bracket $\langle \eta, 2, \alpha \rangle$ is detected by $h_1^3 h_6 c_0$, with indeterminacy generated by $h_1^2 h_4 Q_2$. These elements map non-trivially under projection to the top cell, which contradicts that they are in the image of inclusion of the bottom cell.

Next, suppose that $2\alpha + \tau\eta\epsilon\eta_6$ is detected by $\tau^3 \Delta h_1 d_0 e_0^2$. Then the Toda bracket

$$\left\langle \eta, \left[\begin{array}{ccc} 2 & \tau\eta\epsilon & \tau^2 \beta \end{array} \right], \left[\begin{array}{c} \alpha \\ \eta_6 \\ \bar{\kappa} \end{array} \right] \right\rangle$$

is well-defined, where β is an element of $\pi_{53,29}$ that is detected by $\Delta h_1 d_0^2$. The same argument involving inclusion of the bottom cell into $C\tau$ applies to this Toda bracket.

Finally, assume that $2\alpha + \tau\eta\epsilon\eta_6$ is detected by $\tau^2 M d_0^2$. Table 21 shows that Md_0 detects $\kappa\theta_{4,5}$, so $\tau^2 M d_0^2$ detects $\tau^2 \kappa^2 \theta_{4,5}$. Then $2\alpha + \tau\eta\epsilon\eta_6$ equals either $\tau^2 \kappa^2 \theta_{4,5}$ or $\tau^2 \kappa^2 \theta_{4,5} + \tau^2 \beta \bar{\kappa}$. We can apply the same argument to the Toda bracket

$$\left\langle \eta, \left[\begin{array}{ccc} 2 & \tau\eta\epsilon & \tau^2 \kappa \theta_{4,5} \end{array} \right], \left[\begin{array}{c} \alpha \\ \eta_6 \\ \kappa \end{array} \right] \right\rangle,$$

or to the Toda bracket

$$\left\langle \eta, \left[\begin{array}{cccc} 2 & \tau\eta\epsilon & \tau^2\kappa\theta_{4.5} & \tau^2\beta \end{array} \right], \left[\begin{array}{c} \alpha \\ \eta_6 \\ \kappa \\ \overline{\kappa} \end{array} \right] \right\rangle.$$

We have now shown by contradiction that h_4D_2 does not survive. After ruling out other possibilities by comparison to $C\tau$ and to mmf , the only remaining possibility is that $d_5(h_4D_2)$ equals $\tau^4d_1g^2$. \square

LEMMA 5.60. $d_5(\tau^3gG_0) = \tau M\Delta h_1^2d_0$.

PROOF. Suppose for sake of contradiction that the element τ^3gG_0 survived. It cannot be the target of a hidden τ extension, and it cannot be hit by a differential. Therefore, it maps non-trivially under inclusion of the bottom cell into $C\tau$, and the only possible image is $\Delta^2e_1 + \overline{\tau\Delta h_2e_1g}$.

Let α be an element of $\pi_{86,45}$ that is detected by τ^3gG_0 . Consider the Toda bracket $\langle \alpha, 2\nu, \nu \rangle$. Lemma 4.12 implies that this Toda bracket is detected by $e_0x_{76,9}$, or is detected in higher Adams filtration.

On the other hand, under inclusion of the bottom cell into $C\tau$, the Toda bracket is detected by $\Delta^2h_1g_2$. This is inconsistent with the conclusion of the previous paragraph, since inclusion of the bottom cell can only increase Adams filtrations.

We now know that τ^3gG_0 does not survive. After eliminating other possibilities by comparison to mmf , the only remaining possibility is that $d_5(\tau^3gG_0)$ equals $\tau M\Delta h_1^2d_0$. \square

LEMMA 5.61. $d_5(g_3) = h_6d_0^2$.

PROOF. Table 11 shows that h_1h_6 detects the Toda bracket $\langle \eta, 2, \theta_5 \rangle$. Therefore, $h_1h_6d_0^2$ detects $\kappa^2\langle \eta, 2, \theta_5 \rangle$. Now consider the shuffle

$$\tau\kappa^2\langle \eta, 2, \theta_5 \rangle = \langle \tau\kappa^2, \eta, 2 \rangle \theta_5.$$

Lemma 6.7 shows that the last bracket is zero. Therefore, $h_1h_6d_0^2$ does not support a hidden τ extension, so it is either hit by a differential or in the image of projection from $C\tau$ to the top cell.

In the Adams spectral sequence for $C\tau$, the element $h_0^3h_4^2h_6$ detects the Toda bracket $\langle \theta_4, 2, \theta_5 \rangle$. Therefore, $h_0^3h_4^2h_6$ must be in the image of inclusion of the bottom cell into $C\tau$. In particular, $h_0^3h_4^2h_6$ cannot map to $h_1h_6d_0^2$ under projection from $C\tau$ to the top cell.

Now $h_1h_6d_0^2$ cannot be in the image of projection from $C\tau$ to the top cell, so it must be hit by some differential. The only possibility is that $d_5(h_1g_3)$ equals $h_1h_6d_0^2$. \square

LEMMA 5.62. $d_5(e_0x_{76,9}) = M\Delta h_1c_0d_0$.

PROOF. If $M\Delta h_1c_0d_0$ were a permanent non-zero cycle, then it could not support a hidden τ extension because Lemma 5.80 shows that $MP\Delta h_1d_0$ is hit by some differential. Therefore, it would lie in the image of projection from $C\tau$ to the top cell, and the only possible pre-image is the element $\Delta^2h_1g_2$ in the E_∞ -page of the Adams spectral sequence for $C\tau$.

There is a σ extension from $\Delta^2e_1 + \overline{\tau\Delta h_2e_1g}$ to $\Delta^2h_1g_2$ in the Adams spectral sequence for $C\tau$. Then $M\Delta h_1c_0d_0$ would also have to be the target of a σ extension. The only possible source for this extension would be $M\Delta h_1^2d_0$.

Table 17 shows that Mh_1 detects $\eta\theta_{4.5}$, so $M\Delta h_1^2 d_0$ detects $\eta\theta_{4.5}\{\Delta h_1 d_0\}$. The product $\eta\sigma\theta_{4.5}\{\Delta h_1 d_0\}$ equals zero because $\sigma\{\Delta h_1 d_0\}$ is zero. Therefore, $M\Delta h_1^2 d_0$ cannot support a hidden σ extension to $M\Delta h_1 c_0 d_0$.

We have now shown that $M\Delta h_1 c_0 d_0$ must be hit by some differential, and the only possibility is that equals $d_5(e_0 x_{76,9})$. \square

5.5. Higher differentials

THEOREM 5.63. *Table 9 lists some values of the Adams d_r differential on multiplicative generators of the E_r -page, for $r \geq 6$. For $r \geq 6$, the Adams d_r differential is zero on all multiplicative generators of the E_r -page not listed in the table. The list is complete through the 90-stem, except that:*

- (1) $d_9(\tau h_6 g + \tau h_2 e_2)$ might equal $\Delta^2 h_2 n$.
- (2) $d_9(x_{85,6})$ might equal $M\Delta h_1 d_0$.
- (3) $d_9(h_1 x_{85,6})$ might equal $M\Delta h_1^2 d_0$.
- (4) $d_{10}(h_2 h_6 g)$ or $d_{10}(h_2 h_6 g + h_1^2 f_2)$ might equal $M\Delta h_1^2 d_0$.
- (5) $d_6(\tau \Delta h_1 H_1)$ might equal $\tau M\Delta h_0^2 e_0$.
- (6) $d_9(x_{87,7})$ might equal $\tau M\Delta h_0^2 e_0$.
- (7) $d_6(\Delta^2 f_1)$ might equal $\tau^2 M d_0^3$.

PROOF. The d_r differential on many multiplicative generators is zero. For the majority of such multiplicative generators, the d_r differential is zero because there are no possible non-zero values, or by comparison to the Adams spectral sequence for $C\tau$, or by comparison to tmf or mmf . In a few cases, the multiplicative generator is already known to be a permanent cycle, as shown in Table 5. A few additional cases appear in Table 9 because their proofs require further explanation.

Some of the more difficult computations appear in [5]. The remaining more difficult computations are carried out in the following lemmas. \square

Table 9 lists the multiplicative generators of the Adams E_r -page, for $r \geq 6$, through the 90-stem whose d_r differentials are non-zero, or whose d_r differentials are zero for non-obvious reasons.

REMARK 5.64. Because $d_6(\Delta g_2 g)$ equals $M d_0^3$, the uncertainty about $d_6(\Delta^2 f_1)$ is inconsequential. Either $\Delta^2 f_1$ or $\Delta^2 f_1 + \tau^2 \Delta g_2 g$ survives to the E_∞ -page.

LEMMA 5.65.

- (1) $d_6(\tau Q_3 + \tau n_1) = 0$.
- (2) $d_6(g Q_3) = 0$

PROOF. Several possible differentials on these elements are eliminated by comparison to the Adams spectral sequences for $C\tau$ and for tmf . The only remaining possibility is that $d_6(\tau Q_3 + \tau n_1)$ might equal $\tau^2 M h_1 g$, and that $d_6(g Q_3)$ might equal $\tau M h_1 g^2$.

The element $M\Delta h_0^2 e_0$ is not hit by any differential because Table 5 shows that $h_1^2 c_3$ is a permanent cycle, and Table 11 shows that $\tau^2 g Q_3 = h_4^2 Q_2$ must survive to detect the Toda bracket $\langle \theta_4, \tau \bar{\kappa}, \{t\} \rangle$.

Lemma 6.25 shows that $M\Delta h_0^2 e_0$ detects the Toda bracket $\langle \tau \eta \bar{\kappa}^2, 2, 4 \bar{\kappa}_2 \rangle$, which contains $\tau \bar{\kappa}^2 \langle \eta, 2, 4 \bar{\kappa}_2 \rangle$. Lemma 6.10 shows that this expression contains zero. We now know that $M\Delta h_0^2 e_0$ detects an element in the indeterminacy of the bracket $\langle \tau \eta \bar{\kappa}^2, 4, 2 \bar{\kappa}_2 \rangle$. In fact, it must detect a multiple of $\tau \eta \bar{\kappa}^2$ since $2 \bar{\kappa}_2 \cdot \pi_{42,22}$ is zero.

The only possibility is that $M\Delta h_0^2 e_0$ detects $\bar{\kappa}$ times an element detected by $\tau^2 Mh_1 g$. Therefore, $\tau^2 Mh_1 g$ cannot be hit by a differential. This shows that $\tau Q_3 + \tau n_1$ is a permanent cycle.

We also know that $M\Delta h_0^2 e_0$ is the target of a hidden τ extension, since it detects a multiple of τ . The element $\tau^2 Mh_1 g^2$ is the only possible source of this hidden τ extension, so it cannot be hit by a differential. This shows that $d_6(gQ_3)$ cannot equal $\tau Mh_1 g^2$. \square

LEMMA 5.66.

- (1) $d_6(h_2^2 H_1) = Mc_0 d_0$.
- (2) $d_7(\tau h_2^2 H_1) = MPd_0$.

PROOF. Table 3 shows that MPd_0 equals the Massey product $\langle Pd_0, h_0^3, g_2 \rangle$. This implies that MPd_0 detects the Toda bracket $\langle \tau\eta^2 \bar{\kappa}, 8, \bar{\kappa}_2 \rangle$. Lemma 6.16 shows that this Toda bracket consists entirely of multiples of $\tau\eta^2 \bar{\kappa}$.

We now know that MPd_0 detects a multiple of $\tau\eta^2 \bar{\kappa}$. The only possibility is that MPd_0 detects η times an element detected by $\tau^2 Mh_1 g$.

We will show in Lemma 7.108 that $\tau^2 Mh_1 g$ is the target of a ν extension, so $\tau^2 Mh_1 g$ cannot support a hidden η extension. Therefore, MPd_0 must be hit by some differential. The only possibility is that $d_7(\tau h_2^2 H_1)$ equals MPd_0 . Then $h_2^2 H_1$ cannot survive to the E_7 -page, so $d_6(h_2^2 H_1)$ equals $Mc_0 d_0$. \square

LEMMA 5.67. *The element $\tau h_1 p_1$ is a permanent cycle.*

PROOF. Lemma 5.57, together with results of [5], show that $\tau h_1 p_1$ survives to the E_6 -page. We must eliminate possible higher differentials.

Table 14 shows that there is a hidden τ extension from $\tau h_2^2 C''$ to $\Delta^2 h_1^2 h_4 c_0$. This means that $\tau h_2 C''' + h_1 h_3 (\Delta e_1 + C_0)$ must also support a hidden τ extension.

The two possible targets for this hidden τ extension are $\Delta^2 h_2 c_1$ and $\tau \Delta^2 h_1^2 g + \tau^3 \Delta h_2^2 g^2$. The second possibility is ruled out by comparison to tmf , so $\Delta^2 h_2 c_1$ cannot be hit by a differential. \square

LEMMA 5.68. *The element $Ph_0 h_2 h_6$ is a permanent cycle.*

PROOF. First note that projection from $C\tau$ to the top cell takes $Ph_2 h_6$ to a non-zero element. If $Ph_0 h_2 h_6$ were not a permanent cycle in the Adams spectral sequence for the sphere, then projection from $C\tau$ to the top cell would also take $Ph_0 h_2 h_6$ to a non-zero element. Then the 2 extension from $Ph_2 h_6$ to $Ph_0 h_2 h_6$ in $\pi_{74,38} C\tau$ would project to a 2 extension in $\pi_{73,39}$. However, there are no possible 2 extensions in $\pi_{73,39}$. \square

LEMMA 5.69. $d_7(m_1) = 0$.

PROOF. The only other possibility is that $d_7(m_1)$ equals $\tau^2 g^2 t$. If that were the case, then the ν extension from $\tau g^2 t$ to $\tau^2 c_1 g^3$ would be detected by projection from $C\tau$ to the top cell. However, the homotopy groups of $C\tau$ have no such ν extension. \square

LEMMA 5.70. $d_8(h_1 x_1) = 0$.

PROOF. Table 5 shows that $\tau h_1 x_1$ is a permanent cycle. Then $d_8(\tau h_1 x_1)$ cannot equal $\tau^2 Me_0^2$, and $d_8(h_1 x_1)$ cannot equal τMe_0^2 . \square

LEMMA 5.71. $d_6(h_2 h_4 h_6) = 0$.

PROOF. Table 5 shows that $h_2^2 h_4 h_6$ is a permanent cycle. Therefore, the Adams differential $d_6(h_2^2 h_4 h_6)$ does not equal $\tau h_2 c_1 A'$, and $d_6(h_2 h_4 h_6)$ does not equal $\tau c_1 A'$. \square

LEMMA 5.72. $d_7(x_{87,7}) = 0$.

PROOF. If $\tau \Delta^2 h_2^2 d_1$ were hit by a differential, then the ν extension from $\Delta^2 h_2^2 d_1$ to $\Delta^2 h_1^2 h_3 d_1$ would be detected by projection from $C\tau$ to the top cell. But the homotopy of $C\tau$ has no such ν extension. \square

LEMMA 5.73. *The element $x_{88,10}$ is a permanent cycle.*

PROOF. In the Adams spectral sequence for $C\tau$, there is a hidden η extension from $h_1^2 x_{85,6}$ to $x_{88,10}$. Therefore, $x_{88,10}$ lies in the image of inclusion of the bottom cell into $C\tau$. The only possible pre-image is the element $x_{88,10}$ in the Adams spectral sequence in the sphere, so $x_{88,10}$ must survive. \square

LEMMA 5.74.

- (1) $d_6(h_2^2 g H_1) = M c_0 e_0^2$.
- (2) $d_7(\tau h_2^2 g H_1) = 0$.

PROOF. If $M c_0 e_0^2$ is non-zero in the E_∞ -page, then it detects an element that is annihilated by τ because Lemma 5.75 shows that the only possible target of such an extension is hit by a differential. Then $M c_0 e_0^2$ would be in the image of projection from $C\tau$ to the top cell. The only possible pre-image would be the element $\Delta g_2 g$ of the Adams spectral sequence for $C\tau$.

In the Adams spectral sequence for $C\tau$, there is a σ extension from $g A'$ to $\Delta g_2 g$. Projection from $C\tau$ to the top cell would imply that there is a hidden σ extension in the homotopy groups of the sphere, from $M h_1 e_0^2$ to $M c_0 e_0^2$, because $g A'$ maps to $M h_1 e_0^2$ under projection from $C\tau$ to the top cell.

But $M h_1 e_0^2$ detects $\eta \theta_{4,5} \{e_0^2\}$, which cannot support a σ extension. This establishes the first formula.

For the second formula, if $d_7(\tau h_2^2 g H_1)$ were equal to $\tau^2 \Delta h_2^2 e_0 g^2$, then the same argument would apply, with $\tau \Delta h_2^2 e_0 g^2$ substituted for $M c_0 e_0^2$. \square

LEMMA 5.75. $d_6(\Delta g_2 g) = M d_0^3$.

PROOF. The proof of Lemma 5.51 shows that $M d_0^3$ must be hit by a differential. The only possibility is that $d_6(\Delta g_2 g)$ equals $M d_0^3$.

Alternatively, Lemma 5.66 shows that $d_7(\tau h_2^2 H_1) = M P d_0$. Note that $\tau g \cdot \tau h_2^2 H_1 = 0$ in the E_7 -page. Therefore, $\tau M d_0^3 = \tau g \cdot M P d_0$ must already be zero in the E_7 -page. The only possibility is that $d_6(\tau \Delta g_2 g) = \tau M d_0^3$, and then $d_6(\Delta g_2 g) = M d_0^3$. \square

LEMMA 5.76. *The element $h_0 g_3$ is a permanent cycle.*

PROOF. In the homotopy of $C\tau$, the product $\theta_4 \theta_5$ is detected by $h_0^2 g_3$. In the sphere, the product $\theta_4 \theta_5$ is therefore non-zero and detected in Adams filtration at most 6.

Table 11 shows that the Toda bracket $\langle 2, \theta_4, \theta_4, 2 \rangle$ contains θ_5 . Therefore, the product $\theta_4 \theta_5$ is contained in

$$\theta_4 \langle 2, \theta_4, \theta_4, 2 \rangle = \langle \theta_4, 2, \theta_4, \theta_4 \rangle 2.$$

(Note that the sub-bracket $\langle \theta_4, \theta_4, 2 \rangle$ is zero because $\pi_{61,32}$ is zero.) Therefore, $\theta_4\theta_5$ is divisible by 2. It follows that $\theta_4\theta_5$ is detected by $h_0^2g_3$, and h_0g_3 is a permanent cycle that detects $\langle \theta_4, 2, \theta_4, \theta_4 \rangle$. \square

LEMMA 5.77. $d_6(\Delta_1h_1^2e_1) = 0$.

PROOF. Consider the element $\overline{\tau Mh_2^2g^2}$ in the Adams spectral sequence for $C\tau$. This element cannot be in the image of inclusion of the bottom cell into $C\tau$. Therefore, it must map non-trivially under projection from $C\tau$ to the top cell. The only possibility is that $\tau Mh_2^2g^2$ is the image. Therefore, $\tau Mh_2^2g^2$ cannot be the target of a differential. \square

LEMMA 5.78. $d_7(x_{92,10})$ does not equal $\tau^2\Delta^2h_2g^2$.

PROOF. If $\tau^2\Delta^2h_2g^2$ were hit by a differential, then the 2 extension from $\tau\Delta^2h_2g^2$ to $\tau\Delta^2h_0h_2g^2$ would be detected by projection from $C\tau$ to the top cell. But the homotopy of $C\tau$ has no such 2 extension. \square

LEMMA 5.79. $d_8(\Delta_1h_2e_1) = 0$.

PROOF. Consider the element $e_0x_{76,9}$ in the Adams E_∞ -page for $C\tau$. It cannot be in the image of inclusion of the bottom cell into $C\tau$, so it must project to a non-zero element in the top cell. The only possible image is $M\Delta h_1^3g$. Therefore, $M\Delta h_1^3g$ cannot be the target of a differential. \square

LEMMA 5.80. *The element $MP\Delta h_1d_0$ is hit by some differential.*

PROOF. Table 14 shows that there is a hidden τ extension from $\Delta h_1c_0d_0$ to $P\Delta h_1d_0$. Therefore, $P\Delta h_1d_0$ detects $\tau\epsilon\{\Delta h_1d_0\}$. On the other hand, Tables 17 and 21 show that $P\Delta h_1d_0$ also detects $\tau\eta\kappa\{\Delta h_1h_3\}$. Since there are no elements in higher Adams filtration, we have that $\tau\epsilon\{\Delta h_1d_0\}$ equals $\tau\eta\kappa\{\Delta h_1h_3\}$.

Table 21 shows that MP detects $\tau\epsilon\theta_{4.5}$, so $MP\Delta h_1d_0$ detects $\tau\epsilon\{\Delta h_1d_0\}\theta_{4.5}$, which equals $\tau\eta\kappa\{\Delta h_1h_3\}\theta_{4.5}$. But $\tau\eta\kappa\theta_{4.5}$ is zero because all elements of $\pi_{60,32}$ are detected by tmf . This shows that $MP\Delta h_1d_0$ detects zero, so it must be hit by a differential. \square

CHAPTER 6

Toda brackets

The purpose of this chapter is to establish various Toda brackets that are used elsewhere in this manuscript. Many Toda brackets can be easily computed from the Moss Convergence Theorem 2.16. These are summarized in Table 11 without further discussion. However, some brackets require more complicated arguments. Those arguments are collected in this chapter.

We will need the following \mathbb{C} -motivic version of a theorem of Toda [38, Theorem 3.6] that applies to symmetric Toda brackets.

THEOREM 6.1. *Let α be an element of $\pi_{s,w}$, with s even. There exists an element α^* in $\pi_{2s+1,2w}$ such that $\langle \alpha, \beta, \alpha \rangle$ contains the product $\beta\alpha^*$ for all β such that $\alpha\beta$.*

COROLLARY 6.2. *If $2\beta = 0$, then $\langle 2, \beta, 2 \rangle$ contains $\tau\eta\beta$.*

PROOF. Apply Theorem 6.1 to $\alpha = 2$. We need to find the value of α^* . Table 3 shows that the Massey product $\langle h_0, h_1, h_0 \rangle$ equals τh_1^2 . The Moss Convergence Theorem 2.16 then shows that $\langle 2, \eta, 2 \rangle$ equals $\tau\eta^2$. It follows that α^* equals $\tau\eta$. \square

THEOREM 6.3. *Table 11 lists some Toda brackets in the \mathbb{C} -motivic stable homotopy groups.*

PROOF. The fourth column of the table gives information about the proof of each Toda bracket.

If the fourth column shows a Massey product, then the Toda bracket follows from the Moss Convergence Theorem 2.16. If the fourth column shows an Adams differential, then the Toda bracket follows from the Moss Convergence Theorem 2.16, using the mentioned differential.

A few Toda brackets are established elsewhere in the literature; specific citations are given in these cases.

Additional more difficult cases are established in the following lemmas. \square

Table 11 lists information about some Toda brackets. The third column of Table 11 gives an element of the Adams E_∞ -page that detects an element of the Toda bracket. The fourth column of Table 11 gives partial information about indeterminacies, again by giving detecting elements of the Adams E_∞ -page. We have not completely analyzed the indeterminacies of all brackets when the details are inconsequential for our purposes. The fifth column indicates the proof of the Toda bracket, and the sixth column shows where each specific Toda bracket is used in the manuscript.

LEMMA 6.4. *The Toda bracket $\langle \kappa, 2, \eta \rangle$ contains zero, with indeterminacy generated by $\eta\rho_{15}$.*

PROOF. Using the Adams differential $d_3(h_0h_4) = h_0d_0$, the Moss Convergence Theorem 2.16 shows that the Toda bracket is detected in filtration at least 3. The only element in sufficiently high filtration is Pc_0 , which detects the product $\eta\rho_{15}$. This product lies in the indeterminacy, so the bracket must contain zero. \square

LEMMA 6.5. *The Toda bracket $\langle \kappa, 2, \eta, \nu \rangle$ is detected by τg .*

PROOF. The subbracket $\langle 2, \eta, \nu \rangle$ is strictly zero, since $\pi_{5,3}$ is zero. The subbracket $\langle \kappa, 2, \eta \rangle$ contains zero by Lemma 6.4. Therefore, the fourfold bracket $\langle \kappa, 2, \eta, \nu \rangle$ is well-defined.

Shuffle to obtain

$$\langle \kappa, 2, \eta, \nu \rangle \eta^2 = \kappa \langle 2, \eta, \nu, \eta^2 \rangle.$$

Table 11 shows that ϵ is contained in the Toda bracket $\langle \eta^2, \nu, \eta, 2 \rangle$, so the latter expression equals $\epsilon\kappa$, which is detected by c_0d_0 . It follows that $\langle \kappa, 2, \eta, \nu \rangle$ must be detected by τg . \square

LEMMA 6.6. *The Toda bracket $\langle \epsilon + \eta\sigma, \sigma, 2\sigma \rangle$ contains zero, with indeterminacy generated by $4\nu\bar{\kappa}$ in $\{Ph_1d_0\}$.*

PROOF. Consider the shuffle

$$\langle \epsilon + \eta\sigma, \sigma, 2\sigma \rangle \eta = (\epsilon + \eta\sigma) \langle \sigma, 2\sigma, \eta \rangle.$$

Table 11 shows that h_1h_4 detects $\langle \sigma, 2\sigma, \eta \rangle$, so $h_1h_4c_0$ detects the product $\epsilon \langle \sigma, 2\sigma, \eta \rangle$. On the other hand, Table 21 shows that $h_1h_4c_0$ also detects $\eta\sigma \langle \sigma, 2\sigma, \eta \rangle$. Therefore, $(\epsilon + \eta\sigma) \langle \sigma, 2\sigma, \eta \rangle$ is detected in filtration greater than 5. Then h_4c_0 cannot detect $\langle \epsilon + \eta\sigma, \sigma, 2\sigma \rangle$.

The shuffle

$$2 \langle \epsilon + \eta\sigma, \sigma, 2\sigma \rangle = \langle 2, \epsilon + \eta\sigma, \sigma \rangle 2\sigma = 0$$

shows that no elements of the Toda bracket can be detected by τh_2g or τh_0h_2g .

The element $4\nu\bar{\kappa}$ generates the indeterminacy because it equals $\tau\eta\kappa(\epsilon + \eta\sigma)$. \square

LEMMA 6.7. *The Toda bracket $\langle \tau\kappa^2, \eta, 2 \rangle$ equals zero, with no indeterminacy.*

PROOF. The Adams differential $d_3(\Delta h_2^2) = \tau h_1d_0^2$ implies that the bracket is detected by $h_0 \cdot \Delta h_2^2$, which equals zero. Therefore, the Toda bracket is detected in Adams filtration at least 7, but there are no elements in the Adams E_∞ -page in sufficiently high filtration.

The indeterminacy can be computed by inspection. \square

LEMMA 6.8. *The Toda bracket $\langle \eta^2, \theta_4, \eta^2 \rangle$ contains zero, with indeterminacy generated by $\eta^3\eta_5$.*

PROOF. If the bracket were detected by h_2d_1 , then

$$\nu \langle \eta^2, \theta_4, \eta^2 \rangle = \langle \nu, \eta^2, \theta_4 \rangle \eta^2$$

would be detected by $h_2^2d_1$. However, $h_2^2d_1$ does not detect a multiple of η^2 .

The bracket cannot be detected by $\tau h_1e_0^2$ by comparison to tmf .

By inspection, the only remaining possibility is that the bracket contains zero. The indeterminacy can be computed by inspection. \square

LEMMA 6.9. *The Toda bracket $\langle \tau, \eta^2\kappa_1, \eta \rangle$ is detected by t , with indeterminacy generated by $\eta^3\mu_{33}$.*

PROOF. There is a relation $h_1 \cdot \overline{h_1^2 d_1} = t$ in the homotopy of $C\tau$. Using the connection between Toda brackets and cofibers as described in [16, Section 3.1.1], this shows that t detects the Toda bracket. \square

The indeterminacy is computed by inspection. \square

LEMMA 6.10. *The Toda bracket $\langle \eta, 2, 4\overline{\kappa}_2 \rangle$ contains zero.*

PROOF. The Massey product $Mh_1 = \langle h_1, h_0, h_0^2 g_2 \rangle$ shows that Mh_1 detects the Toda bracket. Table 17 shows that Mh_1 , $\Delta h_2 c_1$, and $\tau d_0 l + \Delta c_0 d_0$ are all targets of hidden η extensions. (Beware that the hidden η extension from $h_3^2 h_5$ to Mh_1 is a crossing extension in the sense of Section 2.1, but that does not matter.) Therefore, Mh_1 detects only multiples of η , so the Toda bracket contains a multiple of η . This implies that it contains zero, since multiples of η belong to the indeterminacy. \square

LEMMA 6.11. *The Toda bracket $\langle \tau\overline{\kappa}_2, \sigma^2, 2 \rangle$ equals zero.*

PROOF. No elements of the bracket can be detected by $\tau^2 \Delta h_1 d_0 g$ by comparison to tmf .

Consider the shuffle

$$\langle \tau\overline{\kappa}_2, \sigma^2, 2 \rangle \kappa = \tau\overline{\kappa}_2 \langle \sigma^2, 2, \kappa \rangle.$$

The bracket $\langle \sigma^2, 2, \kappa \rangle$ is zero because it is contained in $\pi_{29,16} = 0$. On the other hand, while $\{\tau Md_0\}\kappa$ is non-zero and detected by τMd_0^2 . Therefore, no elements of $\langle \tau\overline{\kappa}_2, \sigma^2, 2 \rangle$ can be detected by τMd_0 . \square

LEMMA 6.12. *For every α that is detected by $h_3^2 h_5$, the Toda bracket $\langle 2, \sigma^2, \alpha \rangle$ contains zero. The indeterminacy is generated by $2\tau\overline{\kappa}_3$, which is detected by $\tau^2 d_0^2 l$.*

PROOF. The product $\sigma^2 \alpha$ is zero for every α that is detected by $h_3^2 h_5$. For degree reasons, the only elements that could detect this product either support η extensions or are detected by tmf . Therefore, the bracket is defined for all α .

By comparison to tmf , the bracket cannot be detected by $\tau^4 g^3$. Table 15 shows that $\tau^2 d_0^2 l$ is the target of a hidden 2 extension, so it detects an element in the indeterminacy. Since there are no other possibilities, the bracket must contain zero. \square

REMARK 6.13. This result is consistent with Table 23 of [16], which claims that the bracket $\langle 2, \sigma^2, \theta_{4,5} \rangle$ contains an element that is detected by B_3 . The element B_3 is now known to be zero in the Adams E_∞ -page, so this just means that the bracket contains an element detected in Adams filtration strictly greater than the filtration of B_3 .

LEMMA 6.14. *The Toda bracket $\langle \theta_4, \eta^2, \theta_4 \rangle$ equals zero.*

PROOF. Theorem 6.1 says that there exists an element θ_4^* in $\pi_{61,32}$ such that $\langle \theta_4, \eta^2, \theta_4 \rangle$ contains $\eta^2 \theta_4^*$. The group $\pi_{61,32}$ is zero, so θ_4^* must be zero, and the bracket must contain zero.

In order to compute the indeterminacy of $\langle \theta_4, \eta^2, \theta_4 \rangle$, we must consider the product of θ_4 with elements of $\pi_{33,18}$. There are several cases to consider.

First consider $\{\Delta h_1^2 h_3\}$. The product $\theta_4 \{\Delta h_1^2 h_3\}$ is detected in Adams filtration at least 10, but there are no elements in sufficiently high filtration.

Next consider $\nu\theta_4$ detected by p . The product θ_4^2 is zero [46], so $\nu\theta_4^2$ is also zero.

Finally, consider $\eta\eta_5$ detected by $h_1^2h_5$. Table 11 shows that $\langle\eta, 2, \theta_4\rangle$ detects η_5 . Shuffle to obtain

$$\eta\eta_5\theta_4 = \eta\langle\eta, 2, \theta_4\rangle\theta_4 = \eta^2\langle 2, \theta_4, \theta_4\rangle.$$

The bracket $\langle 2, \theta_4, \theta_4\rangle$ is zero because it is contained in $\pi_{61,32} = 0$. \square

LEMMA 6.15. *The Toda bracket $\langle\eta^2, \theta_4, \eta^2, \theta_4\rangle$ is detected by $\Delta_1h_3^2$.*

PROOF. Table 3 shows that $\Delta_1h_3^2$ equals $\langle h_1^2, h_4^2, h_1^2, h_4^2\rangle$. Therefore, $\Delta_1h_3^2$ detects $\langle\eta^2, \theta_4, \eta^2, \theta_4\rangle$, if the Toda bracket is well-defined.

In order to show that the Toda bracket is well-defined, we need to know that the subbrackets $\langle\eta^2, \theta_4, \eta^2\rangle$ and $\langle\theta_4, \eta^2, \theta_4\rangle$ contain zero. These are handled by Lemmas 6.8 and 6.14. \square

LEMMA 6.16. *The Toda bracket $\langle\tau\eta^2\bar{\kappa}, 8, \bar{\kappa}_2\rangle$ contains zero, and its indeterminacy is generated by multiples of $\tau\eta^2\bar{\kappa}$.*

PROOF. The bracket $\langle\tau\eta^2\bar{\kappa}, 8, \bar{\kappa}_2\rangle$ contains $\tau\eta\bar{\kappa}\langle\eta, 2, 4\bar{\kappa}_2\rangle$. Lemma 6.10 shows that this expression contains zero.

It remains to show that $\bar{\kappa}_2 \cdot \pi_{23,12}$ equals zero. There are several cases to consider.

First, the product $\tau\sigma\eta_4\bar{\kappa}_2$ in $\pi_{60,32}$ could only be detected by τ^4g^3 or $\tau^2d_0^2l$. Comparison to *tmf* rules out both possibilities. Therefore, $\tau\sigma\eta_4\bar{\kappa}_2$ is zero.

Second, the product $\bar{\kappa}\bar{\kappa}_2$ in $\pi_{64,35}$ must be detected in filtration at least 9, since τgg_2 equals zero, so it could only be detected by $h_1^2(\Delta e_1 + C_0)$. This implies that $\tau\nu\bar{\kappa}\bar{\kappa}_2$ is zero.

Third, we must consider the product $\rho_{23}\bar{\kappa}_2$. Table 11 shows that the Toda bracket $\langle\sigma, 16, 2\rho_{15}\rangle$ detects ρ_{23} . Then $\rho_{23}\bar{\kappa}_2$ is contained in

$$\langle\sigma, 16, 2\rho_{15}\rangle\bar{\kappa}_2 = \sigma\langle 16, 2\rho_{15}, \bar{\kappa}_2\rangle.$$

The latter bracket is contained in $\pi_{60,32}$. As above, comparison to *tmf* shows that the expression is zero. \square

LEMMA 6.17. *The Toda bracket $\langle\eta, \nu, \tau\theta_{4,5}\bar{\kappa}\rangle$ is detected by $\tau h_1D_3'$.*

PROOF. Table 6 (see also Remark 5.10) shows that either $d_3(\tau D_3')$ or $d_3(\tau D_3' + \tau^2h_2G_0)$ equals τ^2Mh_2g . In either case, the Moss Convergence Theorem 2.16 shows that $\tau h_1D_3'$ detects the Toda bracket. \square

LEMMA 6.18. *There exists an element α in $\pi_{66,34}$ detected by τ^2h_2C' such that the Toda bracket $\langle\eta, \nu, \alpha\rangle$ is defined and detected by τh_1p_1 .*

PROOF. The differential $d_5(\tau p_1 + h_0^2h_3h_6) = \tau^2h_2^2C'$ and the Moss Convergence Theorem 2.16 establish that the Toda bracket is detected by τh_1p_1 , provided that the Toda bracket is well-defined.

Let α be an element of $\pi_{66,34}$ that is detected by τ^2h_2C' . Then $\nu\alpha$ does not necessarily equal zero; it could be detected in higher filtration by $\tau^2h_2B_5 + h_2D_2'$. Then we can adjust our choice of α by an element detected by $\tau^2B_5 + D_2'$ to ensure that $\nu\alpha$ is zero. \square

LEMMA 6.19. *The Toda bracket $\langle\sigma^2, 2, \{t\}, \tau\bar{\kappa}\rangle$ is detected by $h_4Q_2 + h_3^2D_2$.*

PROOF. The subbracket $\langle \sigma^2, 2, \{t\} \rangle$ contains zero by comparison to $C\tau$, and its indeterminacy is generated by $\sigma^3\theta_4 = 4\sigma\bar{\kappa}_2$ detected by τgn . The subbracket $\langle 2, \{t\}, \tau\bar{\kappa} \rangle$ is strictly zero because it cannot be detected by $h_0h_2h_5i$ by comparison to tmf . This shows that the desired four-fold Toda bracket is well-defined.

Consider the relation

$$\eta\langle \sigma^2, 2, \{t\}, \tau\bar{\kappa} \rangle \subseteq \langle \langle \eta, \sigma^2, 2 \rangle, \{t\}, \tau\bar{\kappa} \rangle.$$

Let α be any element of $\langle \eta, \sigma^2, 2 \rangle$. Table 11 shows that α is detected by h_1h_4 and equals either η_4 or $\eta_4 + \eta\rho_{15}$. By inspection, the indeterminacy of $\langle \alpha, \{t\}, \tau\bar{\kappa} \rangle$ equals $\tau\bar{\kappa} \cdot \pi_{53,29}$, which is detected in Adams filtration at least 14. (In fact, the indeterminacy is non-zero, since it contains both $\tau\bar{\kappa} \cdot \{Mc_0\}$ detected by τMd_0^2 and also $\tau\bar{\kappa} \cdot \{\Delta h_1 d_0^2\}$ detected by $\tau^2 \Delta h_1 d_0 e_0^2$.)

Table 11 shows that $\langle \alpha, \{t\}, \tau\bar{\kappa} \rangle$ is detected by $h_1h_4Q_2$. Together with the partial analysis of the indeterminacy in the previous paragraph, this shows that $\langle \alpha, \{t\}, \tau\bar{\kappa} \rangle$ does not contain zero.

Then $\eta\langle \sigma^2, 2, \{t\}, \tau\bar{\kappa} \rangle$ also does not contain zero, and the only possibility is that $\langle \sigma^2, 2, \{t\}, \tau\bar{\kappa} \rangle$ is detected by $h_4Q_2 + h_3^2D_2$. \square

LEMMA 6.20. *The Toda bracket $\langle \theta_4, \theta_4, \kappa \rangle$ equals zero.*

PROOF. The Massey product $\langle h_4^2, h_4^2, d_0 \rangle$ equals zero, since

$$h_1^2 \langle h_4^2, h_4^2, d_0 \rangle = \langle h_1^2, h_4^2, h_4^2 \rangle d_0 = 0,$$

while $h_1^2 x_{75,7}$ is not zero. The Moss Convergence Theorem 2.16 then implies that $\langle \theta_4, \theta_4, \kappa \rangle$ is detected in Adams filtration at least 8.

The only element in sufficiently high filtration is $Ph_1^4 h_6$. However,

$$\eta^2 \langle \theta_4, \theta_4, \kappa \rangle = \langle \eta^2, \theta_4, \theta_4 \rangle \kappa = 0,$$

while $h_1^2 \cdot Ph_1^4 h_6$ is not zero. Then $\langle \theta_4, \theta_4, \kappa \rangle$ must contain zero because there are no remaining possibilities.

The indeterminacy can be computed by inspection, using that $\theta_4\theta_{4.5}$ is zero by comparison to $C\tau$. \square

LEMMA 6.21. *The Toda bracket $\langle \kappa, 2, \theta_5 \rangle$ is detected by h_6d_0 .*

PROOF. The differential $d_3(h_0h_4) = h_0d_0$ implies that $\langle \kappa, 2, \theta_5 \rangle$ is detected by $h_0h_4 \cdot h_5^2 = 0$ in filtration 4. In other words, the Toda bracket is detected in Adams filtration at least 5.

The element $h_1h_6d_0$ detects $\langle \eta\kappa, 2, \theta_5 \rangle$, using the Adams differential $d_2(h_6) = h_0h_5^2$. This expression contains $\eta\langle \kappa, 2, \theta_5 \rangle$, which shows that $\langle \kappa, 2, \theta_5 \rangle$ is detected in filtration at most 5.

The only possibility is that the Toda bracket is detected by h_6d_0 . \square

LEMMA 6.22. *The Toda bracket $\langle 2, \eta, \tau\eta\{h_1x_{76,6}\} \rangle$ is detected by τh_1x_1 .*

PROOF. Let α be an element of $\pi_{77,41}$ that is detected by $h_1x_{76,6}$. First we must show that the Toda bracket is well-defined.

Note that 2α is zero because there are no 2 extensions in $\pi_{77,41}$ in sufficiently high Adams filtration. Now consider the shuffle

$$\tau\eta^2\alpha = \langle 2, \eta, 2 \rangle \alpha = 2\langle \eta, 2, \alpha \rangle.$$

Table 11 shows that $\langle \eta, 2, \alpha \rangle$ is detected by $h_0 h_2 x_{76,6}$, but this element does not support a hidden 2 extension. This shows that $\tau \eta^2 \alpha$ is zero and that the Toda bracket is well-defined.

Finally, use the Adams differential $d_4(h_0 e_2) = \tau h_1^3 x_{76,6}$ and the relation $h_0 \cdot h_0 e_2 = \tau h_1 x_1$ to compute the Toda bracket. \square

LEMMA 6.23. *The Toda bracket $\langle 2, \sigma^2, \{\tau h_2^2 C'\} \rangle$ equals zero, with no indeterminacy.*

PROOF. Let α be an element of $\pi_{66,35}$ that is detected by $\tau h_2 C'$, so $\nu \alpha$ is the unique element that is detected by $\tau h_2^2 C'$. We consider the Toda bracket $\langle 2, \sigma^2, \nu \alpha \rangle$. By inspection, the indeterminacy is zero, so the bracket equals $\langle 2, \sigma^2, \nu \alpha \rangle$, which equals $\langle \alpha, 2, \sigma^2 \rangle \nu$.

Apply the Moss Convergence Theorem 2.16 with the Adams d_2 differential to see that the Toda bracket $\langle \alpha, 2, \sigma^2 \rangle$ is detected by 0 in Adams filtration 9, but it could be detected by a non-zero element in higher filtration. However, this shows that $\langle \alpha, 2, \sigma^2 \rangle \nu$ is zero by inspection. \square

LEMMA 6.24. *The Toda bracket $\langle 2, \sigma^2, \{h_3(\Delta e_1 + C_0)\} \rangle$ equals zero, with no indeterminacy.*

PROOF. Let β be an element of $\pi_{62,33}$ that is detected by $\Delta e_1 + C_0$, so $\sigma \beta$ is the unique element that is detected by $h_3(\Delta e_1 + C_0)$. We consider the Toda bracket $\langle 2, \sigma^2, \sigma \beta \rangle$. By inspection, the indeterminacy is zero, so the bracket equals $\langle 2, \sigma^2, \beta \rangle \sigma$.

Apply the Moss Convergence Theorem 2.16 with the Adams d_2 differential to see that the Toda bracket $\langle 2, \sigma^2, \beta \rangle$ is detected by 0 in Adams filtration 9, but it could be detected by a non-zero element in higher filtration. Then the only possible non-zero value for $\langle 2, \sigma^2, \beta \rangle \sigma$ is $\{M \Delta h_1 h_3\} \sigma$. Table 21 shows that $M \Delta h_1 h_3$ detects $\{\Delta h_1 h_3\} \theta_{4.5}$, so $\sigma \{M \Delta h_1 h_3\}$ equals $\sigma \{\Delta h_1 h_3\} \theta_{4.5}$, which equals zero. \square

LEMMA 6.25. *The Toda bracket $\langle \tau \eta \bar{\kappa}^2, 2, 4 \bar{\kappa}_2 \rangle$ is detected by $M \Delta h_0^2 e_0$.*

PROOF. Table 3 shows that the Massey product $\langle \Delta h_0^2 e_0, h_0^2, h_0 g_2 \rangle$ equals the element $M \Delta h_0^2 e_0$. Now apply the Moss Convergence Theorem 2.16, using that Table 17 shows that $\Delta h_0^2 e_0$ detects $\tau \eta \bar{\kappa}^2$. \square

LEMMA 6.26. *There exists an element α in $\pi_{67,36}$ that is detected by $h_0 Q_3 + h_0 n_1$ such that $h_1^2 c_3$ detects the Toda bracket $\langle \tau \alpha, \nu_4, \eta \rangle$.*

PROOF. A consequence of the proof of Lemma 5.48 is that there exists α in $\pi_{67,36}$ that is detected by $h_0 Q_3 + h_0 n_1$ such that the product $\tau \nu_4 \alpha$ is zero. Therefore, $h_1^2 c_3$ detects the Toda bracket $\langle \tau \alpha, \nu_4, \eta \rangle$ because of the Adams differential $d_4(h_1 c_3) = \tau h_0 h_2 h_4 Q_3$. \square

Hidden extensions

In this chapter, we will discuss hidden extensions in the E_∞ -page of the Adams spectral sequence. We methodically explore hidden extensions by τ , 2 , η , and ν , and we study other miscellaneous hidden extensions that are relevant for specific purposes.

7.1. Hidden τ extensions

In order to study hidden τ extensions, we will use the long exact sequence

$$(7.1) \quad \cdots \longrightarrow \pi_{p,q+1} \xrightarrow{\tau} \pi_{p,q} \longrightarrow \pi_{p,q}C\tau \longrightarrow \pi_{p-1,q+1} \xrightarrow{\tau} \pi_{p-1,q} \longrightarrow \cdots$$

extensively. This sequence governs hidden τ extensions in the following sense. An element α in $\pi_{p,q}$ is divisible by τ if and only if it maps to zero in $\pi_{p,q}C\tau$, and an element α in $\pi_{p-1,q+1}$ supports a τ extension if and only if it is not in the image of $\pi_{p,q}C\tau$. Therefore, we need to study the maps $\pi_{*,*} \rightarrow \pi_{*,*}C\tau$ and $\pi_{*,*}C\tau \rightarrow \pi_{*-1,*+1}$ induced by inclusion of the bottom cell into $C\tau$ and by projection from $C\tau$ to the top cell.

The E_∞ -pages of the Adams spectral sequences for $S^{0,0}$ and $C\tau$ give associated graded objects for the homotopy groups that are the sources and targets of these maps. Naturality of the Adams spectral sequence induces maps on associated graded objects.

These maps on associated graded objects often detect the values of the maps on homotopy groups. For example, the element h_0 in the Adams spectral sequence for the sphere is mapped to the element h_0 in the Adams spectral sequence for $C\tau$. In homotopy groups, this means that inclusion of the bottom cell into $C\tau$ takes the element 2 in $\pi_{0,0}$ to the element 2 in $\pi_{0,0}C\tau$.

On the other side, the element $\overline{h_1^4}$ in the Adams spectral sequence for $C\tau$ is mapped to the element h_1^4 in Adams spectral sequence for the sphere. In homotopy groups, this means that projection from $C\tau$ to the top cell takes the element $\{\overline{h_1^4}\}$ in $\pi_{5,3}C\tau$ to the element η^4 in $\pi_{4,4}$.

However, some values of the maps on homotopy groups can be hidden in the map of associated graded objects. This situation is rare in low stems but becomes more and more common in higher stems. The first such example occurs in the 30-stem. The element Δh_2^2 is a permanent cycle in the Adams spectral sequence for $C\tau$, so $\{\Delta h_2^2\}$ is an element in $\pi_{30,16}C\tau$. Now Δh_2^2 maps to zero in the E_∞ -page of the Adams spectral sequence for the sphere, but $\{\Delta h_2^2\}$ does not map to zero in $\pi_{29,17}$. In fact $\{\Delta h_2^2\}$ maps to $\eta\kappa^2$, which is detected by $h_1d_0^2$. This demonstrates that projection from $C\tau$ to the top cell has a hidden value.

We refer the reader to Section 2.1 for a precise discussion of these issues.

THEOREM 7.1.

- (1) Through the 90-stem, Table 12 lists all hidden values of inclusion of the bottom cell into $C\tau$, except that:
- (a) If $h_1x_{85,6}$ does not survive, then $\tau h_1x_{85,6}$ maps to $\Delta^2e_1 + \overline{\tau\Delta h_2e_1g}$.
 - (b) If h_2h_6g or $h_2h_6g + h_1^2f_2$ does not survive, then τh_2h_6g or $\tau h_2h_6g + \tau h_1^2f_2$ maps to $\Delta^2e_1 + \overline{\tau\Delta h_2e_1g}$.
 - (c) If $\tau\Delta h_1H_1$ survives, then it maps to $\overline{\Delta h_1B_7}$.
- (2) Through the 90-stem, Table 13 lists all hidden values of projection from $C\tau$ to the top cell, except that:
- (a) If $\tau h_6g + \tau h_2e_2$ does not survive, then τh_6g maps to $\tau(\Delta e_1 + C_0)g$.
 - (b) If $x_{85,6}$ does not survive, then $x_{85,6}$ maps to Δh_1j_1 .
 - (c) If $h_1x_{85,6}$ does not survive, then $\overline{h_1x_{85,6}}$ maps to τMh_0g^2 .
 - (d) If $h_1x_{85,6}$ survives, then $\Delta^2e_1 + \overline{\tau\Delta h_2e_1g}$ maps to $M\Delta h_1^2d_0$.
 - (e) If h_2h_6g or $h_2h_6g + h_1^2f_2$ does not survive, then h_2h_6g maps to τMh_0g^2 .
 - (f) If $x_{87,7}$ does not survive, then $x_{87,7}$ maps to $M\Delta h_0^2e_0$.
 - (g) If $\tau\Delta h_1H_1$ does not survive, then $\overline{\Delta h_1B_7}$ maps to $M\Delta h_0^2e_0$.

PROOF. The values of inclusion of the bottom cell and projection to the top cell are almost entirely determined by inspection of Adams E_∞ -pages. Taking into account the multiplicative structure, there are no other combinatorial possibilities. For example, consider the exact sequence

$$\pi_{30,16} \rightarrow \pi_{30,16}C\tau \rightarrow \pi_{29,17}.$$

In the Adams E_∞ -page for $C\tau$, h_4^2 and Δh_2^2 are the only two elements in the 30-stem with weight 16. In the Adams E_∞ -page for the sphere, h_4^2 is the only element in the 30-stem with weight 16, and $h_1d_0^2$ is the only element in the 29-stem with weight 17. The only possibility is that h_4^2 maps to h_4^2 under inclusion of the bottom cell, and Δh_2^2 maps to $h_1d_0^2$ under projection to the top cell.

One case, given below in Lemma 7.7, requires a more complicated argument. \square

REMARK 7.2. Through the 90-stem, inclusion of the bottom cell into $C\tau$ has only one hidden value with target indeterminacy. Namely, h_2c_1A' is the hidden value of $\overline{h_1gB_7}$, with target indeterminacy generated by Δj_1 . Through the 90-stem, projection from $C\tau$ to the top cell has no hidden values with target indeterminacy.

REMARK 7.3. Through the 90-stem, inclusion of the bottom cell into $C\tau$ has no crossing values. On the other hand, projection from $C\tau$ to the top cell does have crossing values in this range. These occurrences are described in the fourth column of Table 13. Each can be verified by direct inspection.

THEOREM 7.4. Through the 90-stem, Table 14 lists all hidden τ extensions in \mathbb{C} -motivic stable homotopy groups, except that:

- (1) if Δ^2h_2n is not hit by a differential, then there is a hidden τ extension from $\tau(\Delta e_1 + C_0)g$ to Δ^2h_2n .
- (2) if $M\Delta h_1d_0$ is not hit by a differential, then there is a hidden τ extension from Δh_1j_1 to $M\Delta h_1d_0$.
- (3) if $M\Delta h_1^2d_0$ is not hit by a differential, then there is a hidden τ extension from τMh_0g^2 to $M\Delta h_1^2d_0$.

In this range, the only crossing extension is:

- (1) the hidden τ extension from $h_1^2 h_6 c_0$ to $h_0 h_4 D_2$, and the not hidden τ extension on $\tau h_2^2 Q_3$.

PROOF. Almost all of these hidden τ extensions follow immediately from the values of the maps in the long exact sequence (7.1) given in Tables 12 and 13.

For example, consider the element Pd_0 in the Adams E_∞ -page for the sphere, which belongs to the 22-stem with weight 12. Now $\pi_{22,12} C\tau$ is zero because there are no elements in that degree in the Adams E_∞ -page for $C\tau$, so inclusion of the bottom cell takes $\{Pd_0\}$ to zero. Therefore, $\{Pd_0\}$ must be in the image of multiplication by τ . The only possibility is that there is a hidden τ extension from $c_0 d_0$ to Pd_0 . \square

REMARK 7.5. A straightforward analysis of the sequence (7.1) shows that $\Delta^2 h_2 n$ is not hit by a differential if and only if the possible τ extension on $\tau(\Delta e_1 + C_0)g$ occurs. Thus this uncertain hidden τ extension is entirely determined by a corresponding uncertainty in a value of an Adams differential.

REMARK 7.6. If $M\Delta h_1 d_0$ and $M\Delta h_1^2 d_0$ are not hit by differentials, then a straightforward analysis of the sequence (7.1) shows that the possible τ extensions on $\Delta h_1 j_1$ and $\tau M h_0 g^2$ must occur. Thus these uncertainties are entirely determined by corresponding uncertainties in values of the Adams differentials.

LEMMA 7.7.

- (1) The element $h_1 h_3(\Delta e_1 + C_0) + \tau h_2 C''$ maps to $\overline{h_1^4 c_0 Q_2}$ under inclusion of the bottom cell into $C\tau$.
- (2) There is a hidden τ extension from $d_1 e_1$ to $h_1 h_3(\Delta e_1 + C_0)$.

PROOF. Consider the exact sequence $\pi_{70,38} \rightarrow \pi_{70,38} C\tau \rightarrow \pi_{69,39}$. For combinatorial reasons, one of the following two possibilities must occur:

- (a) the element $h_1 h_3(\Delta e_1 + C_0) + \tau h_2 C''$ maps to $\overline{h_1^4 c_0 Q_2}$ under inclusion of the bottom cell into $C\tau$, and there is a hidden τ extension from $d_1 e_1$ to $h_1 h_3(\Delta e_1 + C_0)$.
- (b) the element $h_1 h_3(\Delta e_1 + C_0)$ maps to $\overline{h_1^4 c_0 Q_2}$ under inclusion of the bottom cell into $C\tau$, and there is a hidden τ extension from $d_1 e_1$ to $h_1 h_3(\Delta e_1 + C_0) + \tau h_2 C''$.

We will show that there cannot be a hidden τ extension from $d_1 e_1$ to $h_1 h_3(\Delta e_1 + C_0) + \tau h_2 C''$.

Lemma 7.157 shows that $\tau\nu\{d_1 e_1\}$ equals $\tau\eta\sigma\{k_1\}$. Since there is no hidden τ extension on $h_1 k_1$, there must exist an element α in $\{k_1\}$ such that $\tau\eta\alpha = 0$. Therefore, $\tau\nu\{d_1 e_1\}$ must be zero.

If there were a τ extension from $d_1 e_1$ to $h_1 h_3(\Delta e_1 + C_0) + \tau h_2 C''$, then $\tau\nu\{d_1 e_1\}$ would be detected by

$$h_2 \cdot (h_1 h_3(\Delta e_1 + C_0) + \tau h_2 C'') = \tau h_2^2 C'',$$

and in particular would be non-zero. \square

7.2. Hidden 2 extensions

THEOREM 7.8. Table 15 lists some hidden extensions by 2.

PROOF. Many of the hidden extensions follow by comparison to $C\tau$. For example, there is a hidden 2 extension from $h_0 h_2 g$ to $h_1 c_0 d_0$ in the Adams spectral sequence for $C\tau$. Pulling back along inclusion of the bottom cell into $C\tau$, there

must also be a hidden 2 extension from h_0h_2g to $h_1c_0d_0$ in the Adams spectral sequence for the sphere. This type of argument is indicated by the notation $C\tau$ in the fourth column of Table 15.

Next, Table 14 shows a hidden τ extension from $h_1c_0d_0$ to Ph_1d_0 . Therefore, there is also a hidden 2 extension from τh_0h_2g to Ph_1d_0 . This type of argument is indicated by the notation τ in the fourth column of Table 15.

Many cases require more complicated arguments. In stems up to approximately dimension 62, see [16, Section 4.2.2 and Tables 27–28] [45], and [46]. The higher-dimensional cases are handled in the following lemmas. \square

REMARK 7.9. Through the 90-stem, there are no crossing 2 extensions.

REMARK 7.10. The hidden 2 extension from $h_0h_3g_2$ to τgn is proved in [45], which relies on the “ $\mathbb{R}P^\infty$ -method” to establish a hidden σ extension from τh_3d_1 to Δh_2c_1 and a hidden η extension from τh_1g_2 to Δh_2c_1 . We now have easier proofs for these η and σ extensions, using the hidden τ extension from $h_1^2g_2$ to Δh_2c_1 given in Table 14, as well as the relation $h_2^2d_1 = h_1^2g_2$.

REMARK 7.11. Comparison to synthetic homotopy gives additional information about some possible hidden 2 extensions, including:

- (1) there is a hidden 2 extension from h_0h_5i to $\tau^4e_0^2g$.
- (2) there is no hidden 2 extension from $Px_{76,6}$ to $M\Delta h_1d_0$.

See [5] for more details. We are grateful to John Rognes for pointing out a mistake in [16, Lemma 4.56 and Table 27] concerning the hidden 2 extension on h_0h_5i . Lemma 7.19 shows that the extension occurs but does not determine its target precisely.

REMARK 7.12. The first correct proof of the relation $2\theta_5 = 0$ appeared in [46]. Earlier claims in [28] and [25] were based upon a mistaken understanding of the Toda bracket $\langle \sigma^2, 2, \theta_4 \rangle$. See [16, Table 23] for the correct value of this bracket.

REMARK 7.13. If $M\Delta h_1^2d_0$ is non-zero in the E_∞ -page, then there is a hidden τ extension from τMh_0g^2 to $M\Delta h_1^2d_0$. This implies that there must be a hidden 2 extension from τ^2Mg^2 to $M\Delta h_1^2d_0$.

THEOREM 7.14. *Table 16 lists all unknown hidden 2 extensions, through the 90-stem.*

PROOF. Many possibilities are eliminated by comparison to $C\tau$, to tmf , or to mmf . For example, there cannot be a hidden 2 extension from $h_2^2h_4$ to τh_1g by comparison to $C\tau$.

Many additional possibilities are eliminated by consideration of other parts of the multiplicative structure. For example, there cannot be a hidden 2 extension from Ph_1h_5 to τ^3g^2 because τ^3g^2 supports an h_1 extension and 2η equals zero.

Several cases are a direct consequence of Proposition 7.17.

Some possibilities are eliminated by more complicated arguments. These cases are handled in the following lemmas. \square

REMARK 7.15. The element τh_5^2g detects the product $\overline{\kappa}\theta_5$, so it cannot support a hidden 2 extension since $2\theta_5$ is zero. If $\tau h_6g + \tau h_2e_2$ survives, then there is a hidden τ extension from $\tau(\Delta e_1 + C_0)g$ to Δ^2h_2n . Then there could not be a hidden 2 extension from h_5^2g to $\tau(\Delta e_1 + C_0)g$.

REMARK 7.16. If $M\Delta h_1^2 d_0$ is not zero in the E_∞ -page, then $M\Delta h_1 d_0$ supports an h_1 multiplication, and there cannot be a hidden 2 extension from $Px_{76,6}$ to $M\Delta h_1 d_0$.

PROPOSITION 7.17. *Suppose that 2α and $\tau\eta\alpha$ are both zero. Then $2\langle\alpha, 2, \theta_5\rangle$ is zero.*

PROOF. Consider the shuffle

$$2\langle\alpha, 2, \theta_5\rangle = \langle 2, \alpha, 2\rangle\theta_5.$$

Since $2\theta_5$ is zero, this expression has no indeterminacy. Corollary 6.2 implies that it equals $\tau\eta\alpha\theta_5$, which is zero by assumption. \square

REMARK 7.18. Proposition 7.17 eliminates possible hidden 2 extensions on several elements, including $h_2^2 h_6$, $h_0^3 h_3 h_6$, $h_3^2 h_6$, $h_6 c_1$, $h_2^2 h_4 h_6$, $h_0^5 h_6 i$, and $h_2^2 h_6 g$.

LEMMA 7.19. *There is a hidden 2 extension from $h_0 h_5 i$ to either $\tau M P h_1$ or to $\tau^4 e_0^2 g$.*

PROOF. Table 2 shows that $h_0 h_5 i$ maps to $\Delta^2 h_2^2$ in the homotopy of tmf . The element $\Delta^2 h_2^2$ supports a hidden 2 extension, so $h_0 h_5 i$ must support a hidden 2 extension as well. \square

LEMMA 7.20.

- (1) *There is a hidden 2 extension from $\tau h_1 H_1$ to $\tau h_1(\Delta e_1 + C_0)$.*
- (2) *There is no hidden 2 extension on $\tau X_2 + \tau C'$.*
- (3) *There is a hidden 2 extension from $\tau h_1 h_3 H_1$ to $\tau h_1 h_3(\Delta e_1 + C_0)$.*

PROOF. Table 17 shows that there is an η extension from $\tau h_1 H_1$ to $h_3 Q_2$. Let α be any element of $\pi_{63,33}$ that is detected by $\tau h_1 H_1$. Then $\tau\eta^2\alpha$ is non-zero and detected by $\tau h_1 h_3 Q_2$. Note that $\tau h_1 h_3 Q_2$ cannot be the target of a hidden 2 extension because there are no possibilities.

If 2α were zero, then we would have the shuffling relation

$$\tau\eta^2\alpha = \langle 2, \eta, 2\rangle\alpha = 2\langle\eta, 2, \alpha\rangle.$$

But this would contradict the previous paragraph.

We now know that 2α must be non-zero for every possible choice of α . The only possibility is that there is a hidden 2 extension from $\tau h_1 H_1$ to $\tau h_1(\Delta e_1 + C_0)$, and that there is no hidden 2 extension on $\tau X_2 + \tau C'$. This establishes the first two parts.

The third part follows immediately from the first part by multiplication by h_3 . \square

LEMMA 7.21. *There is a hidden 2 extension from $h_1 h_6$ to $\tau h_1^2 h_5^2$.*

PROOF. Table 11 shows that $h_1 h_6$ detects $\langle\eta, 2, \theta_5\rangle$. Now shuffle to obtain

$$2\langle\eta, 2, \theta_5\rangle = \langle 2, \eta, 2\rangle\theta_5 = \tau\eta^2\theta_5.$$

\square

LEMMA 7.22. *There is no hidden 2 extension on $\Delta_1 h_3^2$.*

PROOF. Table 11 shows that $\Delta_1 h_3^2$ detects the Toda bracket $\langle \eta^2, \theta_4, \eta^2, \theta_4 \rangle$. We have

$$2\langle \eta^2, \theta_4, \eta^2, \theta_4 \rangle \subseteq \langle \langle 2, \eta^2, \theta_4 \rangle, \eta^2, \theta_4 \rangle.$$

Table 11 shows that

$$\nu\theta_4 = \langle 2, \eta, \eta\theta_4 \rangle = \langle 2, \eta^2, \theta_4 \rangle,$$

so we must compute $\langle \nu\theta_4, \eta^2, \theta_4 \rangle$.

This bracket contains $\nu\langle \theta_4, \eta^2, \theta_4 \rangle$, which equals zero by Lemma 6.14. Therefore, we only need to compute the indeterminacy of $\langle \nu\theta_4, \eta^2, \theta_4 \rangle$.

The only possible non-zero element in the indeterminacy is the product $\theta_4\{t\}$. Table 11 shows that $\{t\} = \langle \nu, \eta, \eta\theta_4 \rangle$. Now

$$\theta_4\{t\} = \langle \nu, \eta, \eta\theta_4 \rangle\theta_4 = \nu\langle \eta, \eta\theta_4, \theta_4 \rangle.$$

This last expression is well-defined because θ_4^2 is zero [46], and it must be zero because $\pi_{63,34}$ consists entirely of multiples of η . \square

LEMMA 7.23. *There is no hidden 2 extension on $h_0Q_3 + h_2^2D_3$.*

PROOF. By comparison to the homotopy of $C\tau$, there is no hidden extension with value h_2^2A' . Table 17 shows that $\tau^2\Delta h_2^2e_0g$ supports a hidden η extension. Therefore, it cannot be the target of a 2 extension. \square

LEMMA 7.24. *There is no hidden 2 extension on h_3A' .*

PROOF. Table 11 shows that h_3A' detects the Toda bracket $\langle \sigma, \kappa, \tau\eta\theta_{4.5} \rangle$. Shuffle to obtain

$$\langle \sigma, \kappa, \tau\eta\theta_{4.5} \rangle 2 = \sigma\langle \kappa, \tau\eta\theta_{4.5}, 2 \rangle.$$

The bracket $\langle \kappa, \tau\eta\theta_{4.5}, 2 \rangle$ is zero because it is contained in $\pi_{61,32} = 0$. \square

LEMMA 7.25. *There is no hidden 2 extension on p' .*

PROOF. Table 21 shows that p' detects the product $\sigma\theta_5$, and $2\theta_5$ is already known to be zero [46]. \square

LEMMA 7.26. *There is no hidden 2 extension on $\tau h_1D_3'$.*

PROOF. Table 11 shows that $\tau h_1D_3'$ detects the Toda bracket $\langle \eta, \nu, \tau\theta_{4.5}\bar{\kappa} \rangle$. Now shuffle to obtain

$$2\langle \eta, \nu, \tau\theta_{4.5}\bar{\kappa} \rangle = \langle 2, \eta, \nu \rangle \tau\theta_{4.5}\bar{\kappa},$$

which equals zero because $\langle 2, \eta, \nu \rangle$ is contained in $\pi_{5,3} = 0$. \square

LEMMA 7.27. *There is no hidden 2 extension on $h_1h_3h_6$.*

PROOF. Table 11 shows that h_1h_6 detects the Toda bracket $\langle \eta, 2, \theta_5 \rangle$. Let α be an element of this bracket. Then $h_1h_3h_6$ detects $\sigma\alpha$, and

$$2\sigma\alpha = 2\sigma\langle \eta, 2, \theta_5 \rangle = \sigma\langle 2, \eta, 2 \rangle\theta_5 = \tau\eta^2\sigma\theta_5.$$

Table 11 also shows that h_6c_0 detects the Toda bracket $\langle \epsilon, 2, \theta_5 \rangle$. Let β be an element of this bracket. As in the proof of Lemma 7.28, we compute that 2β equals $\tau\eta\epsilon\theta_5$.

Now consider the element $\sigma\alpha + \beta$, which is also detected by $h_1h_3h_6$. Then

$$2(\sigma\alpha + \beta) = \tau\eta^2\sigma\theta_5 + \tau\eta\epsilon\theta_5 = \tau\nu^3\theta_5,$$

using Toda's relation $\eta^2\sigma + \nu^3 = \eta\epsilon$ [38].

Table 19 shows that there is a hidden ν extension from $h_2h_5^2$ to τh_1Q_3 . Therefore, τh_1Q_3 detects $\nu^2\theta_5$.

This does not yet imply that $\nu^3\theta_5$ is zero, because $\nu^2\theta_5 + \eta\{\tau Q_3 + \tau n_1\}$ might be detected h_3A' or Ph_2h_5j in higher filtration. However, h_3A' does not support a hidden ν extension by Lemma 7.113. Also, Table 2 shows that Ph_2h_5j maps non-trivially to tmf , while $\nu^2\theta_5 + \eta\{\tau Q_3 + \tau n_1\}$ maps to zero. This is enough to conclude that $\nu^3\theta_5$ is zero.

We have now shown that $2(\sigma\alpha + \beta)$ is zero in $\pi_{71,37}$. Since $h_1h_3h_6$ detects $\sigma\alpha + \beta$, it follows that $h_1h_3h_6$ does not support a hidden 2 extension. \square

LEMMA 7.28. *There is a hidden 2 extension from h_6c_0 to $\tau h_1^2p'$.*

PROOF. Table 11 shows that h_6c_0 detects the Toda bracket $\langle \epsilon, 2, \theta_5 \rangle$. Now shuffle to obtain

$$2\langle \epsilon, 2, \theta_5 \rangle = \langle 2, \epsilon, 2 \rangle \theta_5 = \tau \eta \epsilon \theta_5.$$

Finally, $\tau \eta \epsilon \theta_5$ is detected by $\tau h_1^2p'$ because of the relation $h_1^2p' = h_1h_5^2c_0$. \square

LEMMA 7.29. *There is no hidden 2 extension on τh_1p_1 .*

PROOF. Lemma 6.18 shows that τh_1p_1 detects $\langle \eta, \nu, \alpha \rangle$ for some α detected by τ^2h_2C' . Now shuffle to obtain

$$2\langle \eta, \nu, \alpha \rangle = \langle 2, \eta, \nu \rangle \alpha,$$

which is zero because $\langle 2, \eta, \nu \rangle$ is contained in $\pi_{5,3} = 0$. \square

LEMMA 7.30. *There is a hidden 2 extension from $h_2^3H_1$ to τMh_2^2g .*

PROOF. Table 19 shows that there are hidden ν extensions from $h_2^3H_1$ to h_3C'' , and from τMh_2^2g to $Mh_1d_0^2$. Table 15 shows that there is also a hidden 2 extension from h_3C'' to $Mh_1d_0^2$. The only possibility is that there must also be a hidden 2 extension on $h_1^2h_3H_1$. \square

LEMMA 7.31. *There is no hidden 2 extension on Ph_1h_6 .*

PROOF. Table 11 shows that Ph_1h_6 detects the Toda bracket $\langle \mu_9, 2, \theta_5 \rangle$. Shuffle to obtain

$$2\langle \mu_9, 2, \theta_5 \rangle = \langle 2, \mu_9, 2 \rangle \theta_5 = \tau \eta \mu_9 \theta_5.$$

Table 11 also shows that μ_9 is contained in the Toda bracket $\langle \eta, 2, 8\sigma \rangle$. Shuffle again to obtain

$$\tau \eta \mu_9 \theta_5 = \langle \eta, 2, 8\sigma \rangle \tau \eta \theta_5 = \tau \eta^2 \langle 2, 8\sigma, \theta_5 \rangle.$$

Table 11 shows that $h_0^3h_3h_6$ detects $\langle 2, 8\sigma, \theta_5 \rangle$.

By inspection, the product $\eta^2\{h_0^3h_3h_6\}$ can only be detected by $\Delta^2h_1h_4c_0$. However, this cannot occur by comparison to $C\tau$. Therefore, $\eta^2\{h_0^3h_3h_6\}$, and also $\tau \eta^2\{h_0^3h_3h_6\}$, must be zero. \square

LEMMA 7.32. *There is no hidden 2 extension on $h_4Q_2 + h_3^2D_2$.*

PROOF. Table 11 shows that the element $h_4Q_2 + h_3^2D_2$ detects the Toda bracket $\langle \sigma^2, 2, \{t\}, \tau \bar{\kappa} \rangle$. Consider the relation

$$2\langle \sigma^2, 2, \{t\}, \tau \bar{\kappa} \rangle \subseteq \langle \langle 2, \sigma^2, 2 \rangle, \{t\}, \tau \bar{\kappa} \rangle.$$

Corollary 6.2 shows that the Toda bracket $\langle 2, \sigma^2, 2 \rangle$ contains zero since $\eta\sigma^2$ is zero. Therefore, it consists of even multiples of ρ_{15} ; let $2k\rho_{15}$ be any such element in $\pi_{15,8}$.

The Toda bracket $\langle 2k\rho_{15}, \{t\}, \tau\bar{\kappa} \rangle$ contains $k\rho_{15}\langle 2, \{t\}, \tau\bar{\kappa} \rangle$, which equals zero as discussed in the proof of Lemma 6.19. Moreover, its indeterminacy is equal to $\tau\bar{\kappa} \cdot \pi_{52,28}$, which is detected in Adams filtration at least 12. This implies that $\langle 2k\rho_{15}, \{t\}, \tau\bar{\kappa} \rangle$ is detected in Adams filtration at least 12, and that the target of a hidden 2 extension on $h_4Q_2 + h_3^2D_2$ must have Adams filtration at least 12.

The remaining possible targets with Adams filtration at least 12 are eliminated by comparison to $C\tau$ or to mmf . \square

REMARK 7.33. The proof of Lemma 7.32 might be simplified by considering the shuffle

$$2\langle \sigma^2, 2, \{t\}, \tau\bar{\kappa} \rangle = \langle 2, \sigma^2, 2, \{t\} \rangle \tau\bar{\kappa}.$$

However, the latter four-fold bracket may not exist, since both three-fold subbrackets have indeterminacy. See [15] for a discussion of the analogous difficulty with Massey products.

LEMMA 7.34. *There is no hidden 2 extension on $h_2^2Q_3$.*

PROOF. The element $\tau h_2^2Q_3$ detects $\nu^2\{\tau Q_3 + \tau n_1\}$, so it cannot support a hidden 2 extension. This rules out all possible 2 extensions on $h_2^2Q_3$. \square

LEMMA 7.35. *There is no hidden 2 extension on $h_0h_4D_2$.*

PROOF. Table 14 shows that there is a hidden τ extension from $h_1^2h_6c_0$ to $h_0h_4D_2$. Therefore, $h_0h_4D_2$ detects either $\tau\eta\epsilon\eta_6$ or $\tau\eta\epsilon\eta_6 + \nu^2\{\tau Q_3 + \tau n_1\}$, because of the presence of $\tau h_2^2Q_3$ in higher filtration. In either case, $h_0h_4D_2$ cannot support a hidden 2 extension. \square

LEMMA 7.36. *There is a hidden 2 extension from $h_3(\tau Q_3 + \tau n_1)$ to $\tau x_{74,8}$.*

PROOF. Table 21 shows that $\tau x_{74,8}$ detects $\tau\bar{\kappa}_2\theta_4$. Table 11 shows that θ_4 equals the Toda bracket $\langle \sigma^2, 2, \sigma^2, 2 \rangle$.

Now consider the shuffle

$$\tau\bar{\kappa}_2\theta_4 = \tau\bar{\kappa}_2\langle \sigma^2, 2, \sigma^2, 2 \rangle = \langle \tau\bar{\kappa}_2, \sigma^2, 2, \sigma^2 \rangle 2.$$

Lemma 6.11 shows that the latter bracket is well-defined. This implies that $\tau x_{74,8}$ is the target of a hidden 2 extension, and $h_3(\tau Q_3 + \tau n_1)$ is the only possible source. \square

LEMMA 7.37. *There is no hidden 2 extension on h_6d_0 .*

PROOF. Table 11 shows that h_6d_0 detects the Toda bracket $\langle \kappa, 2, \theta_5 \rangle$. Now shuffle to obtain

$$2\langle \kappa, 2, \theta_5 \rangle = \langle 2, \kappa, 2 \rangle \theta_5 = \tau\eta\kappa\theta_5.$$

Lemma 7.156 shows that this product equals $\tau\eta\sigma^2\theta_5$, which equals zero because $\eta\sigma^2$ is zero. \square

LEMMA 7.38. *There is a hidden 2 extension from e_0A' to $M\Delta h_1^2h_3$.*

PROOF. Let α be an element of $\pi_{76,40}$ that is detected by $x_{76,9}$. Table 17 shows that there is a hidden η extension from $x_{76,9}$ to $M\Delta h_1h_3$, so $\tau\eta^2\alpha$ is detected by $\tau M\Delta h_1^2h_3$. Now shuffle to obtain

$$\tau\eta^2\alpha = \langle 2, \eta, 2 \rangle \alpha = 2\langle \eta, 2, \alpha \rangle.$$

This shows that $\tau M\Delta h_1^2h_3$ must be the target of a hidden 2 extension.

Moreover, the source of this hidden 2 extension must be in Adams filtration at least 10, since the Adams differential $d_2(\tau x_{77,8}) = h_0 x_{76,9}$ implies that $\langle \eta, 2, \alpha \rangle$ is detected by $h_1 x_{77,8} = 0$ in filtration 9. The only possible source is $e_0 A'$. \square

LEMMA 7.39. *There is no hidden 2 extension on $h_1 h_4 h_6$.*

PROOF. Table 11 shows that $h_1 h_4 h_6$ detects the Toda bracket $\langle \eta_4, 2, \theta_5 \rangle$. Now shuffle to obtain

$$2\langle \eta_4, 2, \theta_5 \rangle = \langle 2, \eta_4, 2 \rangle \theta_5,$$

which equals $\tau \eta \eta_4 \theta_5$ by Table 11. We will show that this product is zero.

There are several elements in the Adams E_∞ -page that might detect $\eta_4 \theta_5$. The possibilities $h_1 h_6 d_0$ and $x_{78,9}$ are ruled out by comparison to $C\tau$. The possibility $\tau e_0 A'$ is ruled out because Table 15 shows that $e_0 A'$ supports a hidden 2 extension.

Two possibilities remain. If $\eta_4 \theta_5$ is detected by $\tau M \Delta h_1^2 h_3$, then $\tau \eta \eta_4 \theta_5$ must be zero because there are no elements in sufficiently high Adams filtration.

Finally, suppose that $\eta_4 \theta_5$ is detected by $\tau h_1^2 x_{76,6}$. Let α be an element of $\pi_{77,41}$ that is detected by $h_1 x_{76,6}$. If $\eta_4 \theta_5 + \tau \eta \alpha$ is not zero, then it is detected in higher filtration. It cannot be detected by $\tau e_0 A'$ because of the hidden 2 extension on $e_0 A'$. If it is detected by $\tau M \Delta h_1^2 h_3$, then we may change the choice of α to ensure that $\eta_4 \theta_5 + \tau \eta \alpha$ is zero.

We have now shown that $\tau \eta \eta_4 \theta_5$ equals $\tau^2 \eta^2 \alpha$. Shuffle to obtain

$$\tau^2 \eta^2 \alpha = \tau \alpha \langle 2, \eta, 2 \rangle = \langle \alpha, 2, \eta \rangle 2\tau.$$

Table 11 shows that $\langle \alpha, 2, \eta \rangle$ is detected by $h_0 h_2 x_{76,6}$, and Lemma 7.40 shows that this element does not support a hidden 2 extension. Therefore, $\langle \alpha, 2, \eta \rangle 2\tau$ is zero. \square

LEMMA 7.40. *There is no hidden 2 extension on $h_0 h_2 x_{76,6}$.*

PROOF. Let α be an element of $\pi_{76,40}$ that is detected by $h_4 A$. Then $\nu \alpha$ is detected by $h_0 h_2 x_{76,6}$.

Table 21 shows that $h_0 h_4 A$ detects $\sigma^2 \theta_5$. Then $2\alpha + \sigma^2 \theta_5$ equals zero, or is detected by $x_{76,9}$. This implies that $2\nu \alpha$ either equals zero, or is detected by $h_1 x_{78,9}$.

However, there cannot be a hidden 2 extension from $h_0 h_2 x_{76,6}$ to $h_1 x_{78,9}$ by comparison to $C\tau$. \square

LEMMA 7.41. *There is no hidden 2 extension on $Ph_6 c_0$.*

PROOF. Table 21 shows that $Ph_6 c_0$ detects the product $\rho_{15} \eta_6$. Table 11 shows that η_6 is contained in the Toda bracket $\langle \eta, 2, \theta_5 \rangle$. Now shuffle to obtain

$$2\rho_{15} \eta_6 = 2\rho_{15} \langle \eta, 2, \theta_5 \rangle = \rho_{15} \theta_5 \langle 2, \eta, 2 \rangle,$$

which equals $\tau \eta^2 \rho_{15} \theta_5$ by Table 11.

Table 21 shows that $\rho_{15} \theta_5$ is detected by either $h_0 x_{77,7}$ or $\tau^2 m_1$. First suppose that it is detected by $h_0 x_{77,7}$. Table 15 shows that $h_0 x_{77,7}$ is the target of a 2 extension. Then $\rho_{15} \theta_5$ equals 2α modulo higher filtration. In any case, $\tau \eta^2 \rho_{15} \theta_5$ is zero.

Next suppose that $\rho_{15} \theta_5$ is detected by $\tau^2 m_1$. Then $\rho_{15} \theta_5$ equals $\tau^2 \alpha$ modulo higher filtration for some element α detected by m_1 . Table 14 shows that there is a hidden τ extension from $h_1 m_1$ to $M \Delta h_1^2 h_3$. This implies that $\tau \eta \alpha$ is detected by $M \Delta h_1^2 h_3$. Finally, $\tau^3 \eta^2 \alpha = \tau \eta^2 \rho_{15} \theta_5$ must be zero. \square

LEMMA 7.42. *There is no hidden 2 extension on ΔB_6 .*

PROOF. Table 17 shows that there is a hidden η extension from $h_0^6 h_4 h_6$ to $\tau \Delta B_6$. Therefore, $\tau \Delta B_6$ cannot be the source of a hidden 2 extension, so there cannot be a hidden 2 extension from ΔB_6 to $\tau^2 M e_0^2$. \square

LEMMA 7.43. *There is no hidden 2 extension on $h_0^2 h_6 g$.*

PROOF. The element $h_0^2 h_6 g$ equals $h_2^2 h_6 d_0$, so it detects $\nu^2 \{h_6 d_0\}$. \square

LEMMA 7.44. *There is no hidden 2 extension on $\tau h_1 h_6 g$.*

PROOF. If $\tau h_6 g + \tau h_2 e_2$ is a permanent cycle, then $\tau h_1 h_6 g$ detects a multiple of η and cannot support a hidden 2 extension. However, we need a more complicated argument because we do not know if $\tau h_6 g + \tau h_2 e_2$ survives.

The element $\tau h_1 h_6 g$ detects $\eta_6 \bar{\kappa}$. Table 11 shows that η_6 is contained in the Toda bracket $\langle \eta, 2, \theta_5 \rangle$. Shuffle to obtain

$$2\eta_6 \bar{\kappa} = 2\langle \eta, 2, \theta_5 \rangle \bar{\kappa} = \langle 2, \eta, 2 \rangle \theta_5 \bar{\kappa} = \tau \eta^2 \theta_5 \bar{\kappa}.$$

The product $\theta_5 \bar{\kappa}$ is detected by $\tau h_5^2 g$.

There are several possible values for $\eta \{h_5^2 g\}$, but they are either ruled out by comparison to $C\tau$, or they are multiples of h_2 . In all cases, $\eta^2 \{h_5^2 g\}$ must be zero. This implies that $\tau \eta^2 \theta_5 \bar{\kappa}$ is also zero. \square

LEMMA 7.45. *There is a hidden 2 extension from $\tau h_1 f_2$ to $\tau^2 h_2 h_4 Q_3$.*

PROOF. Table 19 shows that $\tau h_1 Q_3$ detects $\nu^2 \theta_5$. The Moss Convergence Theorem 2.16 implies that $\tau^2 h_2 h_4 Q_3$ detects the Toda bracket $\langle 2, \eta_4, \tau \nu^2 \theta_5 \rangle$, using the differential $d_2(h_0^3 c_3) = \tau^2 h_1^2 h_4 Q_3$.

On the other hand, this bracket contains $\langle 2, \eta_4, \nu \rangle \tau \nu \theta_5$. The bracket $\langle 2, \eta_4, \nu \rangle$ in $\pi_{20,11}$ must contain zero by comparison to tmf .

We have now shown that $\tau^2 h_2 h_4 Q_3$ detects a linear combination of a multiple of 2 and a multiple of $\tau \nu^2 \theta_5$. However, $\tau \nu^2 \theta_5 \cdot \pi_{17,9}$ is zero, so $\tau^2 h_2 h_4 Q_3$ detects a multiple of 2. The only possibility is that there is a hidden 2 extension from $\tau h_1 f_2$ to $\tau^2 h_2 h_4 Q_3$. \square

LEMMA 7.46. *There is no hidden 2 extension on $\tau h_2 h_4 Q_3$.*

PROOF. There cannot be a hidden 2 extension from $\tau h_2 h_4 Q_3$ to $\tau P h_1 x_{76,6}$ because there is no hidden τ extension from $h_0 h_2 h_4 Q_3$ to $P h_1 x_{76,6}$.

Table 17 shows that $\tau^3 M g^2$ supports a hidden η extension. Therefore, it cannot be the target of a hidden 2 extension. \square

LEMMA 7.47. *Neither $h_6 c_0 d_0$ nor $P h_6 d_0$ support hidden 2 extensions.*

PROOF. Table 17 shows that both elements are targets of hidden η extensions. \square

LEMMA 7.48. *There is no hidden 2 extension on $h_4 h_6 c_0$.*

PROOF. Table 21 shows that $h_4 h_6 c_0$ detects the product $\sigma \{h_1 h_4 h_6\}$, and the element $h_1 h_4 h_6$ does not support a hidden 2 extension by Lemma 7.39. \square

LEMMA 7.49. *There is no hidden 2 extension on $\tau h_2 g C'$.*

PROOF. The possible target $\tau^3 e_0^3 m$ is ruled out by comparison to mmf . The possible target $Ph_1^7 h_6$ is ruled out by comparison to $C\tau$.

It remains to eliminate the possible target $\tau^2 Mh_1 g^2$. Table 14 shows that there are hidden τ extensions from $\tau h_2 g C'$ and $\tau^2 Mh_1 g^2$ to $\Delta^2 h_2^2 d_1$ and $M\Delta h_0^2 e_0$ respectively. However, there is no hidden 2 extension from $\Delta^2 h_2^2 d_1$ to $M\Delta h_0^2 e_0$, so there cannot be a 2 extension from $\tau h_2 g C'$ to $\tau^2 Mh_1 g^2$. \square

LEMMA 7.50. *There is no hidden 2 extension on $h_1^2 c_3$.*

PROOF. Table 11 shows that the Toda bracket $\langle \tau\{h_0 Q_3 + h_0 n_1\}, \nu_4, \eta \rangle$ is detected by $h_1^2 c_3$. Shuffle to obtain

$$\langle \tau\{h_0 Q_3 + h_0 n_1\}, \nu_4, \eta \rangle 2 = \tau\{h_0 Q_3 + h_0 n_1\} \langle \nu_4, \eta, 2 \rangle.$$

These expressions have no indeterminacy because $\tau\{h_0 Q_3 + h_0 n_1\}$ does not support a 2 extension. Finally, the bracket $\langle \nu_4, \eta, 2 \rangle$ contains zero by comparison to tmf . \square

LEMMA 7.51. *If $\tau\Delta h_1 H_1$ survives, then it supports a hidden 2 extension to either $\tau\Delta^2 h_3 d_1$ or to $\tau^2 \Delta^2 c_1 g$.*

PROOF. Suppose that $\tau\Delta h_1 H_1$ survives. Let α be an element of $\pi_{87,45}$ that is detected by $\tau\Delta h_1 H_1$. Comparison to $C\tau$ shows that $\eta\alpha$ is detected by $\Delta^2 f_1$ or $\Delta^2 f_1 + \tau^2 \Delta g_2 g$, depending only on which one survives. Then $\tau\Delta^2 h_1 f_1$ detects $\tau\eta^2 \alpha$.

If 2α were zero, then we could shuffle to obtain

$$\tau\eta^2 \alpha = \langle 2, \eta, 2 \rangle \alpha = 2 \langle \eta, 2, \alpha \rangle.$$

Considering the Massey product $\langle h_1, h_0, \tau\Delta h_1 H_1 \rangle$, the Moss Convergence Theorem 2.16 would imply that $\langle \eta, 2, \alpha \rangle$ is detected in filtration at least 11. But there are no possible 2 extensions whose source is in filtration at least 11 and whose target is $\tau\Delta^2 h_1 f_1$. \square

LEMMA 7.52. *There is no hidden 2 extension on $P^2 h_6 c_0$.*

PROOF. Table 21 shows that $P^2 h_6 c_0$ detects the product $\rho_{23} \eta_6$. Table 11 shows that η_6 is contained in the Toda bracket $\langle \eta, 2, \theta_5 \rangle$. Shuffle to obtain

$$2\rho_{23} \eta_6 = 2\rho_{23} \langle \eta, 2, \theta_5 \rangle = \langle 2, \eta, 2 \rangle \rho_{23} \theta_5 = \tau\eta^2 \rho_{23} \theta_5.$$

The product $\rho_{23} \theta_5$ is detected in Adams filtration at least 13, and then $\tau\eta^2 \rho_{23} \theta_5$ is detected in filtration at least 16. This rules out all possible targets for a hidden 2 extension on $P^2 h_6 c_0$. \square

LEMMA 7.53. *There is no hidden 2 extension on M^2 .*

PROOF. Table 21 shows that M^2 detects $\theta_{4.5}^2$. Graded commutativity implies that $2\theta_{4.5}^2$ is zero. \square

7.3. Hidden η extensions

THEOREM 7.54. *Table 17 lists some hidden extensions by η .*

PROOF. Many of the hidden extensions follow by comparison to $C\tau$. For example, there is a hidden η extension from $\tau h_1 g$ to $c_0 d_0$ in the Adams spectral sequence for $C\tau$. Pulling back along inclusion of the bottom cell into $C\tau$, there must also be a hidden η extension from $\tau h_1 g$ to $c_0 d_0$ in the Adams spectral sequence for the

sphere. This type of argument is indicated by the notation $C\tau$ in the fourth column of Table 17.

Next, Table 14 shows a hidden τ extension from c_0d_0 to Pd_0 . Therefore, there is also a hidden η extension from τ^2h_1g to Pd_0 . This type of argument is indicated by the notation τ in the fourth column of Table 17.

The proofs of several of the extensions in Table 17 rely on analogous extensions in mmf . Extensions in mmf have not been rigorously analyzed [17]. However, the specific extensions from mmf that we need are easily deduced from extensions in tmf , together with the multiplicative structure. For example, there is a hidden η extension in tmf from an to τd_0^2 . Therefore, there is a hidden η extension in mmf from ang to τd_0^2g , and also a hidden η extension from $\Delta h_2^2e_0$ to $\tau d_0e_0^2$ in the homotopy groups of the sphere spectrum. Note that mmf really is required here, since ang and d_0^2g equal zero in the homotopy of tmf .

Many cases require more complicated arguments. In stems up to approximately dimension 62, see [16, Section 4.2.3 and Tables 29–30] and [45]. The higher-dimensional cases are handled in the following lemmas. \square

REMARK 7.55. The hidden η extension from τC to τ^2gn is proved in [45], which relies on the “ $\mathbb{R}P^\infty$ -method” to establish a hidden σ extension from τh_3d_1 to Δh_2c_1 and a hidden η extension from τh_1g_2 to Δh_2c_1 . We now have easier proofs for these η and σ extensions, using the hidden τ extension from $h_1^2g_2$ to Δh_2c_1 given in Table 14, as well as the relation $h_3^2d_1 = h_1^2g_2$.

REMARK 7.56. By comparison to $C\tau$, $h_0^6h_4h_6$ must support a hidden η extension. The only possible targets are $\tau\Delta B_6$ and Δ^2n . If Δ^2h_2n is not hit by a differential, then the target must be $\tau\Delta B_6$.

REMARK 7.57. If $x_{85,6}$ survives, then there is a hidden τ extension from Δh_1j_1 to $M\Delta h_1d_0$. It follows that there must be a hidden η extension from $\tau\Delta j_1 + \tau^2gC'$ to $M\Delta h_1d_0$.

REMARK 7.58. The last column of Table 17 indicates the crossing η extensions.

THEOREM 7.59. *Table 18 lists all unknown hidden η extensions, through the 90-stem.*

PROOF. Many possible extensions can be eliminated by comparison to $C\tau$, to tmf , or to mmf . For example, there cannot be a hidden η extension from τMd_0 to τ^4g^3 because τ^4g^3 maps to a non-zero element in $\pi_{60}tmf$ that is not divisible by η .

Other possibilities are eliminated by consideration of other parts of the multiplicative structure. For example, there cannot be a hidden η extension whose target supports a multiplication by 2, since 2η equals zero.

Many cases are eliminated by more complicated arguments. These are handled in the following lemmas. \square

REMARK 7.60. If $\tau h_6g + \tau h_2e_2$ survives, then $\tau(\Delta e_1 + C_0)g$ supports a hidden τ extension. It follows that $h_1^3h_4h_6$ cannot support a hidden η extension because $\tau\eta^3\{h_1h_4h_6\} = 4\nu\{h_1h_4h_6\}$ must be zero.

REMARK 7.61. If $\tau h_6g + \tau h_2e_2$ survives, then there is a hidden τ extension from $\tau(\Delta e_1 + C_0)g$ to Δ^2h_2n . Then the possible extension from τgD_3 to $\tau(\Delta e_1 + C_0)g$ occurs if and only if the possible extension from τ^2gD_3 to Δ^2h_2n occurs.

REMARK 7.62. If $\tau h_6 g + \tau h_2 e_2$ survives, then $\tau h_2 h_6 g + \tau h_1^2 f_2$ is a multiple of h_2 . This implies that $\tau h_2 h_6 g + \tau h_1^2 f_2$ cannot support a hidden η extension. On the other hand, if $h_2 h_6 g$ or $h_2 h_6 g + h_1^2 f_2$ does not survive, then $\tau h_2 h_6 g + \tau h_1^2 f_2$ maps to $\Delta^2 e_1 + \tau \Delta h_2 e_1 g$ in the Adams E_∞ -page for $C\tau$. This element supports an h_1 multiplication, so $\tau h_2 h_6 g + \tau h_1^2 f_2$ must support a hidden η extension.

REMARK 7.63. If $h_1 x_{85,6}$ does not survive, then $\tau h_1 x_{85,6}$ maps to $\Delta^2 e_1 + \tau \Delta h_2 e_1 g$ in the Adams E_∞ -page for $C\tau$. This element supports an h_1 multiplication, so $\tau h_1 x_{85,6}$ must support a hidden η extension.

LEMMA 7.64. *There is no hidden η extension on $\tau h_1 Q_2$.*

PROOF. There cannot be a hidden η extension from $\tau h_1 Q_2$ to $\tau^2 \Delta h_1 d_0 g$ by comparison to tmf . It remains to show that there cannot be a hidden η extension from $\tau h_1 Q_2$ to $\tau M d_0$.

Note that $h_1 d_0 Q_2 = \tau^3 d_1 g^2$, so $\kappa\{h_1 Q_2\}$ is detected by $\tau^3 d_1 g^2$. Therefore, $\kappa\{h_1 Q_2\} + \tau \kappa_1 \bar{\kappa}^2$ is detected in higher filtration. The only possibility is $\tau^3 e_0 g m$, but that cannot occur by comparison to mmf . Therefore, $\kappa\{h_1 Q_2\} + \tau \kappa_1 \bar{\kappa}^2$ is zero.

Now $\tau \eta \kappa_1 \bar{\kappa}^2$ is zero because $\tau \eta \kappa_1 \bar{\kappa}$ cannot be detected by $\Delta h_1 d_0^2$ by comparison to tmf . Therefore, $\eta \kappa\{h_1 Q_2\}$ is zero, so $\tau \eta \kappa\{h_1 Q_2\}$ is also zero.

On the other hand, $\tau \kappa\{M d_0\}$ is non-zero because it is detected by $\tau M d_0^2$. Therefore $\tau \eta\{h_1 Q_2\}$ cannot be detected by $\tau M d_0$. \square

LEMMA 7.65. *There is no hidden η extension on $\tau h_1^2 h_5^2$.*

PROOF. Table 15 shows that $\tau h_1^2 h_5^2$ is the target of a hidden 2 extension. \square

LEMMA 7.66. *There is a hidden η extension from $\tau^2 h_1 X_2$ to $\tau^2 M h_0 g$.*

PROOF. Table 17 shows that there is a hidden η extension from $\tau h_1 X_2$ to $c_0 Q_2$. Since $c_0 Q_2$ does not support a hidden τ extension, there exists an element β in $\pi_{65,35}$ that is detected by $c_0 Q_2$ such that $\tau \beta = 0$.

Projection from $C\tau$ to the top cell takes $c_0 Q_2$ and $P(A + A')$ to $c_0 Q_2$ and $\tau M h_0 h_2 g$ respectively. Since $h_2 \cdot c_0 Q_2 = P(A + A')$ in the Adams spectral sequence for $C\tau$, it follows that $\nu \beta$ is non-zero and detected by $\tau M h_0 h_2 g$.

Let α be an element of $\pi_{63,33}$ that is detected by $\tau X_2 + \tau C'$, and consider the sum $\eta^2 \alpha + \beta$. Both terms are detected by $c_0 Q_2$, but the sum could be detected in higher filtration. In fact, the sum is non-zero because $\nu(\eta^2 \alpha + \beta)$ is non-zero.

It follows that $\eta^2 \alpha + \beta$ is detected by $\tau M h_0 g$, and that $\tau \eta^2 \alpha$ is detected by $\tau^2 M h_0 g$. \square

LEMMA 7.67. *There is no hidden η extension on $\tau h_1^3 h_6$.*

PROOF. The element $\tau \eta^2 \eta_6$ is detected by $\tau h_1^3 h_6$. Table 11 shows that η_6 is contained in the Toda bracket $\langle \eta, 2, \theta_5 \rangle$. Now shuffle to obtain

$$\eta \cdot \tau \eta^2 \eta_6 = 4\nu \eta_6 = 4\nu \langle \eta, 2, \theta_5 \rangle = 4 \langle \nu, \eta, 2 \rangle \theta_5,$$

which equals zero because $2\theta_5$ is zero. \square

LEMMA 7.68. *There is no hidden η extension from $\tau \Delta_1 h_3^2$ to $h_2^2 A'$.*

PROOF. Table 19 shows that $h_2^2 A'$ supports a hidden ν extension, so it cannot be the target of a hidden η extension. \square

LEMMA 7.69. *There is no hidden η extension on $\tau h_1 Q_3$.*

PROOF. Table 19 shows that $\tau h_1 Q_3$ is the target of a hidden ν extension. Therefore, it cannot be the source of a hidden η extension. \square

LEMMA 7.70. *There is a hidden η extension from $h_3 A'$ to $h_3(\Delta e_1 + C_0)$.*

PROOF. Comparison to $C\tau$ shows that there is a hidden η extension from $h_3 A'$ to either $\tau h_2^2 C' + h_3(\Delta e_1 + C_0)$ or $h_3(\Delta e_1 + C_0)$. Table 19 shows that $\tau h_2^2 C' + h_3(\Delta e_1 + C_0)$ supports a hidden ν extension. Therefore, it cannot be the target of a hidden η extension. \square

LEMMA 7.71. *There is a hidden η extension from $h_2^2 h_6$ to $\tau h_0 h_2 Q_3$.*

PROOF. Table 3 gives the Massey product $h_0 h_2 = \langle h_1, h_0, h_1 \rangle$. Therefore,

$$\langle \tau h_1 Q_3, h_0, h_1 \rangle = \{ \tau h_0 h_2 Q_3, \tau h_0 h_2 Q_3 + \tau h_1 h_3 H_1 \}.$$

Table 19 shows that there is a hidden ν extension from $h_2 h_5^2$ to $\tau h_1 Q_3$, so $\nu^2 \theta_5$ is detected by $\tau h_1 Q_3$. Therefore, the Toda bracket $\langle \nu^2 \theta_5, 2, \eta \rangle$ is detected by $\tau h_0 h_2 Q_3$ or by $\tau h_0 h_2 Q_3 + \tau h_1 h_3 H_1$.

Now $\langle \nu^2 \theta_5, 2, \eta \rangle$ contains $\nu^2 \langle \theta_5, 2, \eta \rangle$. This expression equals $\nu \theta_5 \langle 2, \eta, \nu \rangle$, which equals zero because $\langle 2, \eta, \nu \rangle$ is contained in $\pi_{5,3} = 0$.

We now know that $\langle \nu^2 \theta_5, 2, \eta \rangle$ equals its own determinacy, so $\tau h_0 h_2 Q_3$ or $\tau h_0 h_2 Q_3 + \tau h_1 h_3 H_1$ detects a multiple of η . The only possibility is that there is a hidden η extension on $h_2^2 h_6$.

The target of this extension cannot be $\tau h_0 h_2 Q_3 + \tau h_1 h_3 H_1$ by comparison to $C\tau$. \square

LEMMA 7.72. *There is a hidden η extension from $\tau h_1 h_3 H_1$ to $h_3^2 Q_2$.*

PROOF. Table 17 shows that there is a hidden η extension from $\tau h_1 H_1$ to $h_3 Q_2$. Now multiply by h_3 . \square

LEMMA 7.73.

- (1) *There is no hidden η extension on $h_1 h_3(\Delta e_1 + C_0)$.*
- (2) *There is no hidden η extension on $\tau h_2 C''' + h_1 h_3(\Delta e_1 + C_0)$.*

PROOF. The element $\tau M h_2^2 g$ is the only possible target for such hidden η extensions. However, Table 19 shows that there is a hidden ν extension from $\tau M h_2^2 g$ to $M h_1 d_0^2$. \square

LEMMA 7.74. *There is no hidden η extension on $h_0^3 h_3 h_6$.*

PROOF. There are several possible targets for a hidden η extension on $h_0^3 h_3 h_6$. The element $\tau \Delta^2 h_2 g$ is ruled out because it supports an h_2 extension. The element $\Delta^2 h_4 c_0$ is ruled out by comparison to $C\tau$. The elements $\tau h_3^2 Q_2$ and $\tau d_0 Q_2$ are ruled out because Table 17 shows that they are targets of hidden η extensions from $\tau^2 h_1 h_3 H_1$ and $\tau^2 h_1 D_3'$ respectively.

The only remaining possibility is $\tau^2 l_1$. This case is more complicated.

Table 11 shows that $h_0^3 h_3 h_6$ detects the Toda bracket $\langle 8\sigma, 2, \theta_5 \rangle$. Now shuffle to obtain

$$\eta \langle 8\sigma, 2, \theta_5 \rangle = \langle \eta, 8\sigma, 2 \rangle \theta_5.$$

Table 11 shows that $\langle \eta, 8\sigma, 2 \rangle$ contains μ_9 and has indeterminacy generated by $\tau \eta^2 \sigma$ and $\tau \eta \epsilon$. Thus the expression $\langle \eta, 8\sigma, 2 \rangle \theta_5$ contains at most four elements.

The product $\mu_9 \theta_5$ is detected in filtration at least 8, so it is not detected by $\tau^2 l_1$. The product $(\mu_9 + \tau \eta^2 \sigma) \theta_5$ is detected by $\tau h_1^2 p'$ because Table 21 shows that there

is a hidden σ extension from h_5^2 to p' . The product $(\mu_9 + \tau\eta\epsilon)\theta_5$ is also detected by $\tau h_1^2 p' = \tau h_1 h_5^2 c_0$. Finally, the product $(\mu_9 + \tau\eta^2\sigma + \tau\eta\epsilon)\theta_5$ equals $(\mu_9 + \tau\nu^3)\theta_5$, which also must be detected in filtration at least 8. \square

LEMMA 7.75. *There is no hidden η extension on $h_2 Q_3$.*

PROOF. There cannot be a hidden η extension on $\tau h_2 Q_3$ because it is a multiple of h_2 . Therefore, the possible targets for an η extension on $h_2 Q_3$ must be annihilated by τ .

The element $h_1^3 h_3 H_1$ cannot be the target because Table 14 shows that it supports a hidden τ extension. The element $\tau M h_2^2 g$ cannot be the target because Table 19 shows that it supports a hidden ν extension to $M h_1 d_0^2$. \square

LEMMA 7.76. *There is no hidden η extension on $\tau h_1^2 p'$.*

PROOF. The element $\tau h_1^2 p'$ detects $\tau\eta^2\sigma\theta_5$ because Table 21 shows that there is a hidden σ extension from h_5^2 to p' . Then $\tau\eta^3\sigma\theta_5$ is zero since $\tau\eta^3\sigma$ is zero. \square

LEMMA 7.77. *There is no hidden η extension on $h_2^3 H_1$.*

PROOF. Table 21 shows that $M d_0$ detects the product $\kappa\theta_{4.5}$. Then Table 11 shows that $h_2^3 H_1$ detects the Toda bracket $\langle \nu, \epsilon, \kappa\theta_{4.5} \rangle$. Now shuffle to obtain

$$\eta\langle \nu, \epsilon, \kappa\theta_{4.5} \rangle = \langle \eta, \nu, \epsilon \rangle \kappa\theta_{4.5},$$

which is zero because $\langle \eta, \nu, \epsilon \rangle$ is contained in $\pi_{13,8} = 0$. \square

LEMMA 7.78. *There is a hidden η extension from $\tau h_1 h_6 c_0$ to $\tau^2 h_2^2 Q_3$.*

PROOF. The hidden τ extension from $h_1^2 h_6 c_0$ to $h_0 d_0 D_2$ implies that $\tau h_1 h_6 c_0$ must support a hidden η extension. However, this hidden τ extension crosses the τ extension from $\tau h_2^2 Q_3$ to $\tau^2 h_2^2 Q_3$. Therefore, the target of the hidden η extension is either $\tau^2 h_2^2 Q_3$ or $h_0 d_0 D_2$.

The element $\tau h_1 h_6 c_0$ detects the product $\tau\eta_6\epsilon$, so we want to compute $\tau\eta\eta_6\epsilon$. Table 11 shows that η_6 belongs to $\langle \theta_5, 2, \eta \rangle$. Shuffle to obtain

$$\tau\eta\eta_6\epsilon = \langle \theta_5, 2, \eta \rangle \tau\eta\epsilon = \theta_5 \langle 2, \eta, \tau\eta\epsilon \rangle.$$

Table 11 shows that $\langle 2, \eta, \tau\eta\epsilon \rangle$ contains ζ_{11} . Finally, $\theta_5\zeta_{11}$ is detected by $\tau^2 h_2^2 Q_3 = h_5^2 \cdot P h_2$. \square

LEMMA 7.79. *There is no hidden η extension on $h_1^3 p'$.*

PROOF. The element $h_1^3 p'$ does not support a hidden τ extension, while Table 14 shows that there is a hidden τ extension from $\tau h_2^2 C''$ to $\Delta^2 h_1^2 h_4 c_0$. Therefore, there cannot be a hidden η extension from $h_1^3 p'$ to $\tau h_2^2 C''$. \square

LEMMA 7.80. *There is a hidden η extension from $h_0 d_0 D_2$ to $\tau M d_0^2$.*

PROOF. Table 11 shows that $h_0 d_0 D_2$ detects the Toda bracket $\langle \tau\bar{\kappa}\theta_{4.5}, 2\nu, \nu \rangle$. Now shuffle to obtain

$$\langle \tau\bar{\kappa}\theta_{4.5}, 2\nu, \nu \rangle \eta = \tau\bar{\kappa}\theta_{4.5} \langle 2\nu, \nu, \eta \rangle.$$

Table 11 shows that the Toda bracket $\langle 2\nu, \nu, \eta \rangle$ contains ϵ . Finally, $\tau\bar{\kappa}\theta_{4.5}\epsilon$ is detected by $\tau M d_0^2$ because Table 21 shows that there is a hidden ϵ extension from $\tau M g$ to $M d_0^2$. \square

LEMMA 7.81. *There is a hidden η extension from $h_0 h_3 d_2$ to $\tau d_1 g_2$.*

PROOF. Table 11 shows that the Toda bracket $\langle \eta, \sigma^2, \eta, \sigma^2 \rangle$ equals κ_1 . We would like to consider the shuffle

$$\langle \eta, \sigma^2, \eta, \sigma^2 \rangle \tau \bar{\kappa}_2 = \eta \langle \sigma^2, \eta, \sigma^2, \tau \bar{\kappa}_2 \rangle,$$

but we must show that the Toda bracket $\langle \eta, \sigma^2, \tau \bar{\kappa}_2 \rangle$ is well-defined and contains zero. It is well-defined because $\sigma^2 \bar{\kappa}_2$ is detected by $h_3^2 g_2$ in $\pi_{58,32}$, and there are no τ extensions on this group. The bracket contains zero by comparison to tmf , since all non-zero elements of $\pi_{60,32}$ are detected by tmf .

We have now shown that $\tau \kappa_1 \bar{\kappa}_2$ is divisible by η . The only possibility is that there is a hidden η extension from $h_0 h_3 d_2$ to $\tau d_1 g_2$. \square

LEMMA 7.82. *There is no hidden η extension on $h_0 x_{77,7}$.*

PROOF. Table 15 shows that $h_0 x_{77,7}$ is the target of a hidden 2 extension. \square

LEMMA 7.83.

- (1) *There is no hidden η extension on $h_3^2 h_6$.*
- (2) *There is no hidden η extension on τm_1 .*

PROOF. Table 15 shows that $e_0 A'$ and $\tau e_0 A'$ support hidden 2 extensions, so they cannot be the targets of hidden η extensions. \square

LEMMA 7.84. *There is no hidden η extension on $h_1 h_6 d_0$.*

PROOF. Table 11 shows that $h_1 h_6$ detects the Toda bracket $\langle \theta_5, 2, \eta \rangle$, so $h_1 h_6 d_0$ detects $\langle \theta_5, 2, \eta \rangle \kappa$. Now shuffle to obtain

$$\langle \theta_5, 2, \eta \rangle \eta \kappa = \theta_5 \langle 2, \eta, \eta \kappa \rangle.$$

Table 11 shows that the Toda bracket $\langle 2, \eta, \eta \kappa \rangle$ equals $\nu \kappa$. Thus we need to compute the product $\nu \kappa \theta_5$. Lemma 7.156 shows that this product equals $\nu \sigma^2 \theta_5$, which equals zero. \square

LEMMA 7.85. *There is no hidden η extension on $\tau h_1^2 x_{76,6}$.*

PROOF. Let α be an element of $\pi_{77,41}$ that is detected by $h_1 x_{76,6}$. Then $\tau h_1^2 x_{76,6}$ detects $\tau \eta \alpha$. Now consider the shuffle

$$\tau \eta^2 \alpha = \langle 2, \eta, 2 \rangle \alpha = 2 \langle \eta, 2, \alpha \rangle.$$

Note that 2α is zero because there are no 2 extensions in $\pi_{77,41}$, so the second bracket is well-defined.

Finally, $2 \langle \eta, 2, \alpha \rangle$ must be zero because there are no 2 extensions in $\pi_{79,42}$ in sufficiently high filtration. \square

LEMMA 7.86. *If $h_2 h_4 h_6$ supports a hidden η extension, then its target is not $\tau h_2^2 x_{76,6}$.*

PROOF. Table 19 shows that $\tau h_2^2 x_{76,6}$ supports a hidden ν extension, so it cannot be the target of a hidden η extension. \square

LEMMA 7.87. *If $h_1^3 h_4 h_6$ supports a hidden extension, then its target is $\tau(\Delta e_1 + C_0)g$, and $\tau h_6 g + \tau h_2 e_2$ does not survive.*

PROOF. The element $\tau h_1^3 h_4 h_6$ is a multiple of h_0 , so it cannot support a hidden η extension. This eliminates all possible targets except for $\tau(\Delta e_1 + C_0)g$.

If $\tau h_6 g + \tau h_2 e_2$ survives, then Remark 7.5 shows that $\tau(\Delta e_1 + C_0)g$ supports a hidden τ extension. As in the previous paragraph, this eliminates $\tau(\Delta e_1 + C_0)g$ as a possible target. \square

LEMMA 7.88. *There is no hidden η extension on $h_3^2 n_1$.*

PROOF. The element $\tau h_3^2 n_1 = h_3^2(\tau Q_3 + \tau n_1)$ detects $\sigma^2\{\tau Q_3 + \tau n_1\}$. Then $\eta\sigma^2\{\tau Q_3 + \tau n_1\}$ is zero because $\eta\sigma^2$ is zero. \square

LEMMA 7.89. *There is no hidden η extension on $\Delta^2 p$.*

PROOF. Table 19 shows that $\Delta^2 p$ is the target of a hidden ν extension, so it cannot be the source of an η extension. \square

LEMMA 7.90. *There is no hidden η extension on $h_6 c_1$.*

PROOF. Table 11 shows that $h_6 c_1$ detects the Toda bracket $\langle \bar{\sigma}, 2, \theta_5 \rangle$. By inspection, all possible indeterminacy is in higher Adams filtration, so $h_6 c_1$ detects every element of the Toda bracket.

Shuffle to obtain

$$\eta\langle \bar{\sigma}, 2, \theta_5 \rangle = \langle \eta, \bar{\sigma}, 2 \rangle \theta_5.$$

The Toda bracket $\langle \eta, \bar{\sigma}, 2 \rangle$ is detected in filtration at least 5 since the Massey product $\langle h_1, c_1, h_0 \rangle$ is zero. Therefore, the Toda bracket equals $\{0, \eta \bar{\kappa}\}$.

We now know that $\eta\langle \bar{\sigma}, 2, \theta_5 \rangle$ contains zero, and therefore $h_6 c_1$ does not support a hidden η extension. \square

LEMMA 7.91. *There is no hidden η extension on $h_0 h_6 g$.*

PROOF. Table 11 shows that $h_0 h_6 g$ detects the Toda bracket $\langle \nu, \eta, \eta_6 \kappa \rangle$. Shuffle to obtain

$$\eta\langle \nu, \eta, \eta_6 \kappa \rangle = \langle \eta, \nu, \eta \rangle \eta_6 \kappa.$$

Table 11 shows that $\langle \eta, \nu, \eta \rangle$ equals ν^2 . Finally,

$$\nu^2 \eta_6 \kappa = \nu^2 \kappa \langle \eta, 2, \theta_5 \rangle = \nu \theta_5 \kappa \langle \nu, \eta, 2 \rangle,$$

which equals zero because $\langle \nu, \eta, 2 \rangle$ is contained in $\pi_{5,3} = 0$. \square

LEMMA 7.92. *There is no hidden η extension on $h_2 h_4 Q_3$.*

PROOF. We must eliminate $\tau h_2 g C'$ as a possible target. One might hope to use the homotopy of $C\tau$ in order to do this, but the homotopy of $C\tau$ has an η extension in the relevant degree that could possibly detect a hidden extension from $h_2 h_4 Q_3$ to $\tau h_2 g C'$.

If there were a hidden η extension from $h_2 h_4 Q_3$ to $\tau h_2 g C'$, then the hidden τ extension from $\tau h_2 g C'$ to $\Delta^2 h_2^2 d_1$ would imply that there is a hidden η extension from $\tau h_2 h_4 Q_3$ to $\Delta^2 h_2^2 d_1$. However, $\tau h_2 h_4 Q_3$ detects the product $\nu_4\{\tau Q_3 + \tau n_1\}$, and $\eta \nu_4$ is zero. Therefore, $\tau h_2 h_4 Q_3$ cannot support a hidden η extension. \square

LEMMA 7.93. *If $h_1 x_{85,6}$ (resp., $h_2 h_6 g$, $h_2 h_6 g + h_1^2 f_2$) does not survive, then there is a hidden η extension from $\tau h_1 x_{85,6}$ (resp., $\tau h_2 h_6 g$, $\tau h_2 h_6 g + \tau h_1^2 f_2$) to either $\tau^2 g Q_3$ or to $\Delta^2 h_3 d_1$.*

PROOF. If $h_1x_{85,6}$ does not survive, then Table 12 shows that $\tau h_1x_{85,6}$ maps to $\Delta^2e_1 + \tau\Delta h_2e_{1g}$ under inclusion of the bottom cell into $C\tau$. This element supports an h_1 multiplication to $\Delta^2h_3d_1$ in $C\tau$. Therefore, $\tau h_1x_{85,6}$ would also support a hidden η extension.

The arguments for h_2h_6g and $h_2h_6g + h_1^2f_2$ are identical. \square

LEMMA 7.94.

- (1) *There is no hidden η extension on $h_1h_6c_0d_0$.*
- (2) *There is no hidden η extension on $Ph_1h_6d_0$.*

PROOF. Table 15 shows that $h_1h_6c_0d_0$ and $Ph_1h_6d_0$ are targets of hidden 2 extensions, so they cannot be the sources of hidden η extensions. \square

LEMMA 7.95. *There is a hidden η extension from $h_0^3h_6i$ to $\tau^2\Delta^2c_{1g}$.*

PROOF. The Adams differential $d_2(\Delta^3h_3^2) = \Delta^2h_0^3x$ implies that $\tau^2\Delta^2c_{1g} = h_1 \cdot \Delta^3h_3^2$ detects the Toda bracket $\langle \eta, 2, \{\Delta^2h_0^2x\} \rangle$. However, the later Adams differential $d_5(h_0^2h_6i) = \Delta^2h_0^2x$ implies that 0 belongs to $\{\Delta^2h_0^2x\}$. Therefore, $\tau^2\Delta^2c_{1g}$ detects $\langle \eta, 2, 0 \rangle$, so $\tau^2\Delta^2c_{1g}$ detects a multiple of η . The only possibility is that there is a hidden η extension from $h_0^3h_6i$ to $\tau^2\Delta^2c_{1g}$. \square

LEMMA 7.96. *There is no hidden η extension on B_6d_1 .*

PROOF. Table 15 shows that B_6d_1 is the target of a hidden 2 extension, so it cannot be the source of a hidden η extension. \square

LEMMA 7.97. *There is no hidden η extension on $h_1^2h_4h_6c_0$.*

PROOF. Table 19 shows that $h_1^2h_6h_6c_0$ is the target of a hidden ν extension, so it cannot support a hidden η extension. \square

LEMMA 7.98. *There is a hidden η extension from $\Delta^2h_1f_1$ to $\tau\Delta^2h_2c_{1g}$.*

PROOF. The element $\tau\Delta^2h_2c_{1g}$ detects the product $\nu^2\{\Delta^2t\}$. Table 11 shows that ν^2 equals the Toda bracket $\langle \eta, \nu, \eta \rangle$. Shuffle to obtain

$$\langle \eta, \nu, \eta \rangle \{\Delta^2t\} = \eta \langle \nu, \eta, \{\Delta^2t\} \rangle.$$

This shows that $\tau\Delta^2h_2c_{1g}$ is the target of a hidden η extension. The only possible source for this extension is $\Delta^2h_1f_1$. \square

7.4. Hidden ν extensions

THEOREM 7.99. *Table 19 lists some hidden extensions by ν .*

PROOF. Many of the hidden extensions follow by comparison to $C\tau$. For example, there is a hidden ν extension from h_0^2g to $h_1c_0d_0$ in the Adams spectral sequence for $C\tau$. Pulling back along inclusion of the bottom cell into $C\tau$, there must also be a hidden ν extension from h_0^2g to $h_1c_0d_0$ in the Adams spectral sequence for the sphere. This type of argument is indicated by the notation $C\tau$ in the fourth column of Table 17.

Next, Table 14 shows a hidden τ extension from $h_1c_0d_0$ to Ph_1d_0 . Therefore, there is also a hidden ν extension from τh_0^2g to Ph_1d_0 . This type of argument is indicated by the notation τ in the fourth column of Table 17.

Some extensions can be resolved by comparison to tmf or to mmf . For example, Table 2 shows that the classical unit map $S \rightarrow tmf$ takes $\{\Delta h_1 h_3\}$ in π_{32} to a non-zero element α of $\pi_{32}tmf$ such that $\nu\alpha = \eta\kappa\bar{\kappa}$ in $\pi_{35}tmf$. Therefore, there must be a hidden ν extension from $\Delta h_1 h_3$ to $\tau h_1 e_0^2$.

The proofs of several of the extensions in Table 19 rely on analogous extensions in mmf . Extensions in mmf have not been rigorously analyzed [17]. However, the specific extensions from mmf that we need are easily deduced from extensions in tmf , together with the multiplicative structure. For example, there is a hidden ν extension in tmf from Δh_1 to τd_0^2 . Therefore, there is a hidden ν extension in mmf from $\Delta h_1 g$ to $\tau d_0^2 g$, and also a hidden ν extension from $\tau \Delta h_1 g$ to $\tau^2 d_0 e_0^2$ in the homotopy groups of the sphere spectrum. Note that mmf really is required here, since $d_0^2 g$ equals zero in the homotopy of tmf .

Many cases require more complicated arguments. In stems up to approximately dimension 62, see [16, Section 4.2.4 and Tables 31–32] and [45]. The higher-dimensional cases are handled in the following lemmas. \square

REMARK 7.100. The last column of Table 19 indicates which ν extensions are crossing, as well as which extensions have indeterminacy in the sense of Section 2.1.1.

REMARK 7.101. The hidden ν extension from $h_2 h_5 d_0$ to $\tau g n$ is proved in [45], which relies on the “ $\mathbb{R}P^\infty$ -method” to establish a hidden σ extension from $\tau h_3 d_1$ to $\Delta h_2 c_1$ and a hidden η extension from $\tau h_1 g_2$ to $\Delta h_2 c_1$. We now have easier proofs for these η and σ extensions, using the hidden τ extension from $h_1^2 g_2$ to $\Delta h_2 c_1$ given in Table 14, as well as the relation $h_3^2 d_1 = h_1^2 g_2$.

REMARK 7.102. If $M\Delta h_1^2 d_0$ is not hit by a differential, then there is a hidden τ extension from $\tau M h_0 g^2$ from $M\Delta h_1^2 d_0$. This implies that there must be a hidden ν extension from $\tau(\Delta e_1 + C_0)g$ to $M\Delta h_1^2 d_0$.

THEOREM 7.103. *Table 20 lists all unknown hidden ν extensions, through the 90-stem.*

PROOF. Many possible extensions can be eliminated by comparison to $C\tau$, to tmf , or to mmf . For example, there cannot be a hidden ν extension from $h_0 h_2 h_4$ to $\tau h_1 g$ by comparison to $C\tau$.

Other possibilities are eliminated by consideration of other parts of the multiplicative structure. For example, there cannot be a hidden ν extension whose target supports a multiplication by η , since $\eta\nu$ equals zero.

Many cases are eliminated by more complicated arguments. These are handled in the following lemmas. \square

REMARK 7.104. Comparison to synthetic homotopy eliminates several possible hidden ν extensions, including:

- (1) from $\tau h_1 p_1$ to $\tau x_{74,8}$.
- (2) from $\Delta^2 p$ to $\tau M h_1 d_0$.

See [5] for more details.

REMARK 7.105. If $M\Delta h_1^2 d_0$ is not hit by a differential, then $M\Delta h_1 d_0$ supports an h_1 extension, and there cannot be a hidden ν extension from $h_0 h_2 h_4 h_6$ to $M\Delta h_1 d_0$.

LEMMA 7.106. *There is a hidden ν extension from $\Delta e_1 + C_0$ to τMh_0g .*

PROOF. Table 11 shows that $2\bar{\kappa}$ is contained in $\tau\langle\nu, \eta, \eta\kappa\rangle$. Shuffle to obtain that

$$\nu\langle\eta, \eta\kappa, \tau\theta_{4.5}\rangle = \langle\nu, \eta, \eta\kappa\rangle\tau\theta_{4.5},$$

so $2\bar{\kappa}\theta_{4.5}$ is divisible by ν .

Table 21 shows that τMg detects $\bar{\kappa}\theta_{4.5}$, so τMh_0g detects $2\bar{\kappa}\theta_{4.5}$. Now we know that there is a hidden ν extension whose target is τMh_0g , and the only possible source is $\Delta e_1 + C_0$. \square

REMARK 7.107. One consequence of the proof of Lemma 7.106 is that $\Delta e_1 + C_0$ detects the Toda bracket $\langle\eta, \eta\kappa, \tau\theta_{4.5}\rangle$.

LEMMA 7.108. *There is a hidden ν extension from τh_1H_1 to $\tau^2 Mh_1g$.*

PROOF. Lemma 6.4 shows that the bracket $\langle\kappa, 2, \eta\rangle$ contains zero with indeterminacy generated by $\eta\rho_{15}$. The bracket $\langle\tau\eta\theta_{4.5}, \kappa, 2\rangle$ equals zero since $\pi_{61,32}$ is zero. Therefore, the Toda bracket $\langle\tau\eta\theta_{4.5}, \kappa, 2, \eta\rangle$ is well-defined.

Table 11 shows that τg detects $\langle\kappa, 2, \eta, \nu\rangle$. Therefore, $\tau^2 Mh_1g$ detects

$$\tau\eta\theta_{4.5}\langle\kappa, 2, \eta, \nu\rangle = \langle\tau\eta\theta_{4.5}, \kappa, 2, \eta\rangle\nu.$$

This shows that $\tau^2 Mh_1g$ is the target of a ν extension, and the only possible source is τh_1H_1 . \square

REMARK 7.109. The proof of Lemma 7.108 shows that τh_1H_1 detects the Toda bracket $\langle\tau\eta\theta_{4.5}, \kappa, 2, \eta\rangle$.

LEMMA 7.110. *There is no hidden ν extension on h_1h_6 .*

PROOF. Table 11 shows that h_1h_6 detects the Toda bracket $\langle\eta, 2, \theta_5\rangle$. Shuffle to obtain

$$\nu\langle\eta, 2, \theta_5\rangle = \langle\nu, \eta, 2\rangle\theta_5 = 0,$$

since $\langle\nu, \eta, 2\rangle$ is contained in $\pi_{5,3} = 0$. \square

LEMMA 7.111. *There is no hidden ν extension on h_3Q_2 .*

PROOF. Table 17 shows that $\tau^2\Delta h_2^2e_0g$ supports a hidden η extension. Therefore, it cannot be the target of a ν extension. \square

LEMMA 7.112. *There is a hidden ν extension from h_2^2A' to $h_1h_3(\Delta e_1 + C_0)$.*

PROOF. By comparison to $C\tau$, There cannot be a hidden ν extension from h_2^2A' to $\tau h_2C'' + h_1h_3(\Delta e_1 + C_0)$

Table 11 shows that $\Delta e_1 + C_0$ detects the Toda bracket $\langle\eta, \eta\kappa, \tau\theta_{4.5}\rangle$, and h_2A' detects the Toda bracket $\langle\nu, \eta, \tau\kappa\theta_{4.5}\rangle$. Note that h_2A' also detects $\langle\nu, \eta\kappa, \tau\theta_{4.5}\rangle$.

Now shuffle to obtain

$$(\eta\sigma + \epsilon)\langle\eta, \eta\kappa, \tau\theta_{4.5}\rangle + \nu^2\langle\nu, \eta\kappa, \tau\theta_{4.5}\rangle = \left\langle \begin{bmatrix} \eta\sigma + \epsilon & \nu^2 \end{bmatrix}, \begin{bmatrix} \eta \\ \nu \end{bmatrix}, \eta\kappa \right\rangle \tau\theta_{4.5}.$$

The matrix Toda bracket $\left\langle \begin{bmatrix} \eta\sigma + \epsilon & \nu^2 \end{bmatrix}, \begin{bmatrix} \eta \\ \nu \end{bmatrix}, \eta\kappa \right\rangle$ must equal $\{0, \nu^2\bar{\sigma}\}$, since $\nu^2\bar{\sigma} = \{h_1^2h_4c_0\}$ is the only non-zero element of $\pi_{25,15}$, and that element belongs to the indeterminacy because it is a multiple of ν^2 .

Next observe that $\tau\nu^2\bar{\sigma}\theta_{4.5}$ is zero because all possible values of $\bar{\sigma}\theta_{4.5}$ are multiples of η . This shows that

$$(\eta\sigma + \epsilon)\alpha + \nu^2\beta = 0,$$

for some α and β detected by $\Delta e_1 + C_0$ and h_2A' respectively. The product $(\eta\sigma + \epsilon)\alpha$ is detected by $h_1h_3(\Delta e_1 + C_0)$, so there must be a hidden ν extension from h_2^2A' to $h_1h_3(\Delta e_1 + C_0)$. \square

LEMMA 7.113. *There is no hidden ν extension on h_3A' .*

PROOF. Table 11 shows that h_3A' detects the Toda bracket $\langle \sigma, \kappa, \tau\eta\theta_{4.5} \rangle$. Now shuffle to obtain

$$\nu\langle \sigma, \kappa, \tau\eta\theta_{4.5} \rangle = \langle \nu, \sigma, \kappa \rangle \tau\eta\theta_{4.5} = \langle \eta, \nu, \sigma \rangle \tau\kappa\theta_{4.5}.$$

The Toda bracket $\langle \eta, \nu, \sigma \rangle$ is zero because it is contained in $\pi_{12,7} = 0$. \square

LEMMA 7.114. *There is no hidden ν extension on p' .*

PROOF. Table 21 shows that p' detects the product $\sigma\theta_5$. Therefore, it cannot support a hidden ν extension. \square

LEMMA 7.115. *There is a hidden ν extension from h_2^2C' to $\tau^2d_1g^2$.*

PROOF. Let α be an element of $\pi_{63,33}$ that is detected by $\tau X_2 + \tau C'$. Table 21 shows that $\epsilon\alpha$ is detected by d_0Q_2 , so $\eta\epsilon\alpha$ is detected by $\tau^3d_1g^2$. On the other hand, $\eta\sigma\alpha$ is zero by comparison to $C\tau$.

Now consider the relation $\eta^2\sigma + \nu^3 = \eta\epsilon$. This shows that $\nu^3\alpha$ is detected by $\tau^3d_1g^2$. Since $\nu^2\alpha$ is detected by $\tau h_2^2C'$, there must be a hidden ν extension from h_2^2C' to $\tau^2d_1g^2$. \square

LEMMA 7.116. *There is a hidden ν extension from $\tau h_1D_3'$ to τMd_0^2 .*

PROOF. Table 11 shows that $\tau h_1D_3'$ detects the Toda bracket $\langle \eta, \nu, \tau\bar{\kappa}\theta_{4.5} \rangle$. Now shuffle to obtain

$$\nu\langle \eta, \nu, \tau\bar{\kappa}\theta_{4.5} \rangle = \langle \nu, \eta, \nu \rangle \tau\bar{\kappa}\theta_{4.5}.$$

The bracket $\langle \nu, \eta, \nu \rangle$ equals $\eta\sigma + \epsilon$ [38].

Now we must compute $(\eta\sigma + \epsilon)\tau\bar{\kappa}\theta_{4.5}$. The product $\sigma\bar{\kappa}$ is zero, and Table 21 shows that $\epsilon\bar{\kappa}\theta_{4.5}$ is detected by Md_0^2 . These two observations imply that $(\eta\sigma + \epsilon)\tau\bar{\kappa}\theta_{4.5}$ is detected by τMd_0^2 . \square

LEMMA 7.117. *There is no hidden ν extension on h_6c_0 .*

PROOF. Table 11 shows that h_6c_0 detects the Toda bracket $\langle \epsilon, 2, \theta_5 \rangle$. Now shuffle to obtain

$$\nu\langle \epsilon, 2, \theta_5 \rangle = \langle \nu, \epsilon, 2 \rangle \theta_5.$$

Finally, the Toda bracket $\langle \nu, \epsilon, 2 \rangle$ is zero because it is contained in $\pi_{12,7} = 0$. \square

LEMMA 7.118. *There is a hidden ν extension from h_2^2C'' to τg^2t .*

PROOF. Let α be an element of $\pi_{53,30}$ that is detected by i_1 . Table 19 shows gt detects $\nu\alpha$. Therefore τg^2t detects $\nu\kappa\alpha$, so τg^2t must be the target of a hidden ν extension. The element h_2^2C'' is the only possible source for this extension. \square

LEMMA 7.119. *There is no hidden ν extension on Mh_1h_3g .*

PROOF. If there were a hidden ν extension from Mh_1h_3g to τg^2t , then there would also be a hidden ν extension with target τ^2g^2t . But there is no possible source for such an extension. \square

LEMMA 7.120. *If there is a hidden ν extension on $h_0h_3d_2$, then its target is $M\Delta h_1^2h_3$.*

PROOF. The only other possible target is e_0A' . However, Table 15 shows that e_0A' supports a hidden 2 extension, while $h_0h_3d_2$ does not. \square

LEMMA 7.121. *There is no hidden ν extension on τd_1g_2 .*

PROOF. Table 17 shows that τd_1g_2 is the target of a hidden η extension. Therefore, it cannot be the source of a hidden ν extension. \square

LEMMA 7.122. *There is no hidden ν extension on h_0h_4A .*

PROOF. Table 21 shows that h_0h_4A detects $\sigma^2\theta_5$, so it cannot support a hidden ν extension. \square

LEMMA 7.123. *There is a hidden ν extension from $h_3^2h_6$ to τh_1x_1 .*

PROOF. Table 11 shows that $h_3^2h_6$ detects $\langle\theta_5, 2, \sigma^2\rangle$. Let α be an element of $\pi_{77,40}$ that is contained in this Toda bracket. Then $\nu\alpha$ is an element of

$$\langle\theta_5, 2, \sigma^2\rangle\nu = \langle\theta_5, 2\sigma, \sigma\rangle\nu = \theta_5\langle 2\sigma, \sigma, \nu\rangle \subseteq \langle 2\sigma, \sigma\theta_5, \nu\rangle.$$

Table 21 shows that p' detects $\sigma\theta_5$. Therefore, the Toda bracket $\langle 2\sigma, \sigma\theta_5, \nu\rangle$ is detected by an element of the Massey product $\langle h_0h_3, p', h_2\rangle$. Table 3 shows that h_0e_2 equals the Massey product $\langle h_3, p', h_2\rangle$. By inspection of indeterminacy, the Massey product $\langle h_0h_3, p', h_2\rangle$ contains $h_0^2e_2 = \tau h_1x_1$ with indeterminacy generated by $h_0h_6e_0$.

We have now shown that $\nu\alpha$ is detected by either τh_1x_1 or $\tau h_1x_1 + h_0h_6e_0$. But $h_0h_6e_0 = h_2h_6d_0$ is a multiple of h_2 , so we may add an element in higher Adams filtration to α , if necessary, to conclude that $\nu\alpha$ is detected by τh_1x_1 . \square

LEMMA 7.124. *There is a hidden ν extension from $h_0^7h_4h_6$ to $\tau\Delta^2h_1d_1$.*

PROOF. Table 19 shows that there is a hidden ν extension from $h_0^6h_4h_6$ to Δ^2p . Therefore, there is also a hidden ν extension from $h_0^7h_4h_6$ to $h_0\cdot\Delta^2p = \tau\Delta^2h_1d_1$. \square

LEMMA 7.125. *There is no hidden ν extension on e_0A' .*

PROOF. A possible hidden ν extension from e_0A' to $\Delta h_1^2B_6$ would be detected by $C\tau$, but we have to be careful with the analysis of the homotopy of $C\tau$ because of the h_2 extension from $\overline{\Delta h_1d_1g}$ to $\Delta h_1^2B_6$ in the Adams E_∞ -page for $C\tau$.

Let α be an element of $\pi_{75,40}C\tau$ that is detected by $\overline{h_3C'}$. Then $\nu\alpha$ is detected by e_0A' , and $\nu\alpha$ maps to zero under projection to the top cell because h_3C' does not support a ν extension in the homotopy of the sphere.

Therefore, $\nu\alpha$ lies in the image of e_0A' under inclusion of the bottom cell. Since $\nu^2\alpha$ is zero, e_0A' cannot support a hidden ν extension to $\Delta h_1^2B_6$. \square

LEMMA 7.126. *There is no hidden ν extension on $h_1h_4h_6$.*

PROOF. Table 11 shows that $h_1h_4h_6$ detects the Toda bracket $\langle \eta_4, 2, \theta_5 \rangle$. Shuffle to obtain

$$\nu\langle \eta_4, 2, \theta_5 \rangle = \langle \nu, \eta_4, 2 \rangle \theta_5.$$

Finally, $\langle \nu, \eta_4, 2 \rangle$ must contain zero in $\pi_{20,11}$ because tmf detects every element of $\pi_{20,11}$. \square

LEMMA 7.127. *There is no hidden ν extension on $h_3^2n_1$.*

PROOF. The element h_2gD_3 cannot be the target of a hidden ν extension by comparison to $C\tau$.

The element $\tau h_3^2n_1 = h_3 \cdot h_3(\tau Q_3 + \tau n_1)$ detects a multiple of σ , so it cannot support a hidden ν extension. This rules out h_2gA' as a possible target. \square

LEMMA 7.128. *There is no hidden ν extension on τe_1g_2 .*

PROOF. After eliminating other possibilities by comparison to tmf , comparison to mmf , and by inspection of h_1 multiplications, the only possible target for a hidden ν extension is $Ph_1x_{76,6}$.

Let α be an element of $\pi_{82,45}$ that is detected by e_1g_2 . Then $\nu\alpha$ is detected by $h_2e_1g_2 = h_1^3h_4Q_3$. Choose an element β of $\pi_{83,45}$ that is detected by $h_1h_4Q_3$ such that $\tau\beta$ is zero. Then $\eta^2\beta$ is also detected by $h_1^3h_4Q_3$. However, $\nu\alpha + \eta^2\beta$ is not necessarily zero; it could be detected in Adams filtration at least 13. In any case, $\tau\nu\alpha$ equals $\tau\eta^2\beta = 0$ modulo filtration 13. In particular, $\tau\nu\alpha$ cannot be detected by $Ph_1x_{76,6}$ in filtration 11. \square

LEMMA 7.129. *There is no hidden ν extension on $P^2h_0h_2h_6$.*

PROOF. Table 19 shows that there is a hidden ν extension from $P^2h_2h_6$ to Δ^2h_0x . The target of a hidden ν extension on $P^2h_0h_2h_6$ must have Adams filtration greater than the filtration of Δ^2h_0x . The only possibilities are ruled out by comparison to tmf . \square

LEMMA 7.130. *There is a hidden ν extension from $(\Delta e_1 + C_0)g$ to τMh_0g^2 .*

PROOF. Let α be an element of $\pi_{62,33}$ that is detected by $\Delta e_1 + C_0$. Table 11 shows that $(\Delta e_1 + C_0)g$ detects $\langle \alpha, \eta^3, \eta_4 \rangle$. Then

$$\nu\langle \alpha, \eta^3, \eta_4 \rangle = \langle \nu\alpha, \eta^3, \eta_4 \rangle$$

by inspection of indeterminacies. Table 19 shows that τMh_0g detects $\nu\alpha$. The Toda bracket $\langle \nu\alpha, \eta^3, \eta_4 \rangle$ is detected by the Massey product

$$\langle \tau Mh_0g, h_1^3, h_1h_4 \rangle = \langle \tau Mh_0g, h_1^4, h_4 \rangle = Mh_0g\langle \tau, h_1^4, h_4 \rangle = \tau Mh_0g^2.$$

\square

LEMMA 7.131. *There is no hidden ν extension on h_2c_1A' .*

PROOF. Table 11 shows that $\tau h_2c_1A'$ detects $\langle \tau\theta_{4,5}\kappa, \eta, \nu \rangle \tau\bar{\sigma}$. Shuffle to obtain

$$\langle \tau\theta_{4,5}\kappa, \eta, \nu \rangle \tau\bar{\sigma}\nu = \tau\theta_{4,5}\kappa\langle \eta, \nu, \tau\nu\bar{\sigma} \rangle.$$

The Toda bracket $\langle \eta, \nu, \tau\nu\bar{\sigma} \rangle$ is zero because $\pi_{27,15}$ contains only a v_1 -periodic element detected by $P^3h_1^3$.

We now know that $\tau h_2c_1A'$ does not support a hidden ν extension. In particular, there cannot be a hidden ν extension from $\tau h_2c_1A'$ to $M\Delta h_0^2e_0$. The hidden

τ extension from $\tau^2 Mh_1g^2$ to $M\Delta h_0^2e_0$ implies that there cannot be a hidden ν extension from h_2c_1A' to $\tau^2 Mh_1g^2$.

Additional cases are ruled out by comparison to $C\tau$ and to mmf . \square

LEMMA 7.132.

- (1) *There is a hidden ν extension from $\Delta j_1 + \tau gC'$ to $\tau^2 Mh_1g^2$.*
- (2) *There is a hidden ν extension from $\tau^2 gC'$ to $M\Delta h_0^2e_0$.*

PROOF. Table 19 shows that there exists an element α in $\pi_{63,33}$ detected by τh_1H_1 such that ν is detected by $\tau^2 Mh_1g$. (Beware that there is a crossing extension here, so not every element detected by τh_1H_1 has the desired property.) Table 21 shows that $\tau^2 Mh_1g$ also detects $\tau\theta_{4.5}\eta\bar{\kappa}$. However, $\nu\alpha$ does not necessarily equal $\tau\theta_{4.5}\eta\bar{\kappa}$ because the difference could be detected in higher filtration by $\Delta^2 h_1^3 h_4$. In any case, $\nu\bar{\kappa}\alpha$ equals $\tau\theta_{4.5}\eta\bar{\kappa}^2$.

The product $\theta_{4.5}\eta\bar{\kappa}^2$ is detected by $\tau^2 Mh_1g^2$. The hidden τ extension from $\tau^2 Mh_1g^2$ to $M\Delta h_0^2e_0$ then implies that $\nu\bar{\kappa}\alpha = \tau\theta_{4.5}\eta\bar{\kappa}^2$ is detected by $M\Delta h_0^2e_0$.

We now know that $M\Delta h_0^2e_0$ is the target of a hidden ν extension. The only possible source is $\tau^2 gC'$. (Lemma 7.131 eliminates another possible source.) This establishes the second extension. The first extension follows from consideration of τ extensions. \square

REMARK 7.133. The proof of Lemma 7.132 shows that $\nu\bar{\kappa}\alpha$ is detected by $M\Delta h_0^2e_0$, where α is detected by τh_1H_1 . Note that $\bar{\kappa}\alpha$ is detected by $\tau^2 h_1H_1g = \tau h_2c_1A'$. But this does not show that $\tau h_2c_1A'$ supports a hidden ν extension. Rather, it shows that the source of the hidden ν extension is either $\tau h_2c_1A'$, or a non-zero element in higher filtration.

LEMMA 7.134. *There is no hidden ν extension on $h_2^2 h_4 h_6$.*

PROOF. Table 11 shows that $h_2^2 h_4 h_6$ detects the Toda bracket $\langle \nu\nu_4, 2, \theta_5 \rangle$. Shuffle to obtain

$$\nu\langle \nu\nu_4, 2, \theta_5 \rangle = \langle \nu, \nu\nu_4, 2 \rangle \theta_5.$$

The Toda bracket $\langle \nu, \nu\nu_4, 2 \rangle$ is zero because $\pi_{25,14}$ consists only of a ν_1 -periodic element detected by $P^2 h_1 c_0$. \square

LEMMA 7.135. *There is a hidden ν extension from $\tau h_1 f_2$ to $h_1 x_{87,7} + \tau^2 g_2^2$.*

PROOF. By comparison to $C\tau$, there must be a hidden ν extension whose target is either $h_1 x_{87,7}$ or $h_1 x_{87,7} + \tau^2 g_2^2$.

Table 15 shows that there is a hidden 2 extension from $\tau h_1 f_2$ to $\tau^2 h_2 h_4 Q_3$, and Table 19 shows that there is a hidden ν extension from $\tau^2 h_2 h_4 Q_3$ to $\tau^2 h_0 g_2^2$. This implies that the target of the ν extension on $\tau h_1 f_2$ must be $h_1 x_{87,7} + \tau^2 g_2^2$. \square

LEMMA 7.136. *There is no hidden ν extension on $x_{85,6}$.*

PROOF. Let α be an element of $\pi_{85,45}$ that is detected by $x_{85,6}$. Table 16 shows that if 2α is non-zero, then it is detected in filtration 11. Then $2\nu\alpha$ must be detected in filtration at least 13.

The product $\nu\alpha$ cannot be detected by τg_2^2 , for then $2\nu\alpha$ would be detected by $\tau h_0 g_2^2$ in filtration 9. This rules out τg_2^2 as a possible target for a hidden ν extension on $x_{85,6}$. \square

LEMMA 7.137. *There is no hidden ν extension on $P^2 h_6 c_0$.*

PROOF. Table 21 shows that $P^2h_6c_0$ detects the product $\rho_{23}\eta_6$, and $\nu\rho_{23}\eta_6$ is zero. \square

LEMMA 7.138. *There is a hidden ν extension from h_2^2gA' to $\Delta h_1^2g_2g$.*

PROOF. Comparison to $C\tau$ shows that h_2^2gA' supports a hidden ν extension whose target is either $\Delta h_1^2g_2g$ or $\Delta h_1^2g_2g + \tau h_2gC''$.

Let α be an element of $\pi_{84,46}$ that is detected by h_2gA' . Since h_2gA' does not support a hidden η extension, we may choose α such that $\eta\alpha$ is zero. Note that h_2^2gA' detects $\nu\alpha$.

Shuffle to obtain

$$\nu^2\alpha = \langle \eta, \nu, \eta \rangle \alpha = \eta \langle \nu, \eta, \alpha \rangle.$$

This shows that $\nu^2\alpha$ must be divisible by η . Consequently, the hidden ν extension on h_2^2gA' must have target $\Delta h_1^2g_2g$. \square

REMARK 7.139. The proof of Lemma 7.138 shows that Δh_1g_2g detects the Toda bracket $\langle \nu, \eta, \{h_2gA'\} \rangle$.

7.5. Miscellaneous hidden extensions

LEMMA 7.140.

- (1) *There is a hidden ϵ extension from $h_3^2h_5$ to Mc_0 .*
- (2) *There is a hidden ϵ extension from $\tau h_3^2h_5$ to MP .*

PROOF. Table 17 shows that Mh_1 detects the product $\eta\theta_{4.5}$. Then Mh_1c_0 detects $\eta\epsilon\theta_{4.5}$. This implies that Mc_0 detects $\epsilon\theta_{4.5}$.

This only shows that Mc_0 is the target of a hidden ϵ extension, whose source could be $h_3^2h_5$ or h_5d_0 . However, Lemma 7.147 rules out the latter case. This establishes the first hidden extension.

Table 14 shows that there is a hidden τ extension from Mc_0 to MP . Then the first hidden extension implies the second one. \square

REMARK 7.141. We claimed in [16, Table 33] that there is a hidden ϵ extension from $h_3^2h_5$ to Mc_0 . However, the argument given in [16, Lemma 4.108] only implies that Mc_0 is the target of a hidden extension from either $h_3^2h_5$ or h_5d_0 .

LEMMA 7.142. *There is a hidden κ extension from $h_3^2h_5$ to Md_0 .*

PROOF. Table 17 shows that Mh_1 detects the product $\eta\theta_{4.5}$. Then Mh_1d_0 detects the product $\eta\kappa\theta_{4.5}$, so Md_0 must detect the product $\kappa\theta_{4.5}$. This shows that Md_0 is the target of a hidden κ extension whose source is either $h_3^2h_5$ or h_5d_0 .

We showed in Lemma 7.147 that $\epsilon\alpha$ is zero for some element α of $\pi_{45,24}$ that is detected by h_5d_0 . Then $\epsilon\bar{\kappa}\alpha$ is also zero. Table 21 shows that $\epsilon\bar{\kappa}$ equals κ^2 . Therefore, $\kappa^2\alpha$ is zero. If $\kappa\alpha$ were detected by Md_0 , then $\kappa^2\alpha$ would be detected by Md_0^2 . It follows that there is no hidden κ extension from h_5d_0 to Md_0 . \square

REMARK 7.143. We showed in [16, Table 33] that there is a hidden κ extension from either $h_3^2h_5$ or h_5d_0 to Md_0 . Lemma 7.142 settles this uncertainty.

LEMMA 7.144. *There is a hidden $\bar{\kappa}$ extension from $h_3^2h_5$ to τMg .*

PROOF. Table 17 shows that Mh_1 detects the product $\eta\theta_{4.5}$. Then τMh_1g detects the product $\eta\bar{\kappa}\theta_{4.5}$, so τMg must detect the product $\bar{\kappa}\theta_{4.5}$. This shows that τMg is the target of a hidden $\bar{\kappa}$ extension whose source is either $h_3^2h_5$ or h_5d_0 .

We showed in Lemma 7.147 that $\epsilon\alpha$ is zero for some element α of $\pi_{45,24}$ that is detected by h_5d_0 . If $\bar{\kappa}\alpha$ were detected by τMg , then $\epsilon\bar{\kappa}\alpha$ would be detected by Md_0^2 because Table 21 shows that there is a hidden ϵ extension from τMg to Md_0^2 . Therefore, there is no hidden $\bar{\kappa}$ extension from h_5d_0 to τMg . \square

LEMMA 7.145. *There is a hidden $\{\Delta h_1h_3\}$ extension from $h_3^2h_5$ to $M\Delta h_1h_3$.*

PROOF. Table 17 shows that Mh_1 detects the product $\eta\theta_{4.5}$. Therefore, the element $M\Delta h_1^2h_3$ detects $\eta\{\Delta h_1h_3\}\theta_{4.5}$. This shows that $M\Delta h_1h_3$ is the target of a hidden $\{\Delta h_1h_3\}$ extension. Lemma 7.148 rules out h_5d_0 as a possible source. The only remaining possible source is $h_3^2h_5$. \square

LEMMA 7.146. *There is a hidden $\theta_{4.5}$ extension from $h_3^2h_5$ to M^2 .*

PROOF. The proof of Lemma 5.28 shows that M^2h_1 detects a multiple of $\eta\theta_{4.5}$. Therefore, it detects either $\eta\theta_{4.5}^2$ or $\eta\theta_{4.5}\{h_5d_0\}$.

Now $h_1h_5d_0$ detects $\eta\{h_5d_0\}$, which also detects $\eta_4\theta_4$ by Table 21. In fact, the proof of [16, Lemma 4.112] shows that these two products are equal. Then $\eta\theta_{4.5}\{h_5d_0\}$ equals $\eta_4\theta_4\theta_{4.5}$. Next, $\eta_4\theta_{4.5}$ lies in $\pi_{61,33}$. The only non-zero element of $\pi_{61,33}$ is detected by mmf , so the product $\eta_4\theta_{4.5}$ must be zero.

We have now shown that M^2h_1 detects $\eta\theta_{4.5}^2$. This implies that M^2 detects $\theta_{4.5}^2$. \square

LEMMA 7.147. *There is no hidden ϵ extension on h_5d_0 .*

PROOF. Table 11 shows that h_5d_0 detects the Toda bracket $\langle 2, \theta_4, \kappa \rangle$. Now shuffle to obtain

$$\epsilon\langle 2, \theta_4, \kappa \rangle = \langle \epsilon, 2, \theta_4 \rangle \kappa.$$

Table 11 shows that h_5c_0 detects the Toda bracket $\langle \epsilon, 2, \theta_4 \rangle$, and there is no indeterminacy. Let α in $\pi_{39,21}$ be the unique element of this Toda bracket. We wish to compute $\alpha\kappa$.

Table 11 shows that h_5c_0 also detects the Toda bracket $\langle \eta_5, \nu, 2\nu \rangle$, with indeterminacy generated by $\sigma\eta_5$. Let β in $\pi_{39,21}$ be an element of this Toda bracket. Then α and β are equal, modulo $\sigma\eta_5$ and modulo elements in higher filtration. Both τh_3d_1 and τ^2c_1g detect multiples of σ . Also, the difference between α and β cannot be detected by Δh_1d_0 by comparison to tmf .

This implies that α equals $\beta + \sigma\gamma$ for some element γ in $\pi_{31,17}$. Then

$$\alpha\kappa = (\beta + \sigma\gamma)\kappa = \beta\kappa$$

because $\sigma\kappa$ is zero.

Now shuffle to obtain

$$\beta\kappa = \langle \eta_5, \nu, 2\nu \rangle \kappa = \eta_5\langle \nu, 2\nu, \kappa \rangle.$$

Table 11 shows that $\langle \nu, 2\nu, \kappa \rangle$ contains $\eta\bar{\kappa}$, and its indeterminacy is generated by $\nu\nu_4$. We now need to compute $\eta_5\eta\bar{\kappa}$.

The product $\eta_5\bar{\kappa}$ is detected by $\tau h_1h_3g_2 = \tau h_1h_5g$, so $\eta_5\bar{\kappa}$ equals $\tau\eta\sigma\bar{\kappa}_2$, modulo elements of higher filtration. But these elements of higher filtration are either annihilated by η or detected by tmf , so $\eta_5\eta\bar{\kappa}$ equals $\tau\eta^2\sigma\bar{\kappa}_2$. By comparison to tmf , this latter expression must be zero. \square

LEMMA 7.148. *There is no hidden $\{\Delta h_1 h_3\}$ extension on $h_5 d_0$.*

PROOF. Table 11 shows that $h_5 d_0$ detects the Toda bracket $\langle \kappa, \theta_4, 2 \rangle$. By inspection of indeterminacies, we have

$$\{\Delta h_1 h_3\} \langle \kappa, \theta_4, 2 \rangle = \langle \{\Delta h_1 h_3\} \kappa, \theta_4, 2 \rangle.$$

Table 21 shows that $\tau d_0 l + \Delta c_0 d_0$ detects the product $\{\Delta h_1 h_3\} \kappa$. Now apply the Moss Convergence Theorem 2.16 with the Adams differential $d_2(h_5) = h_0 h_4^2$ to determine that the Toda bracket $\langle \{\Delta h_1 h_3\} \kappa, \theta_4, 2 \rangle$ is detected in Adams filtration at least 13.

The only element in sufficiently high filtration is $\tau^5 e_0 g^3$, but comparison to *mmf* rules this out. Thus the Toda bracket $\langle \{\Delta h_1 h_3\} \kappa, \theta_4, 2 \rangle$ contains zero. \square

LEMMA 7.149. *There is a hidden κ extension from h_5^2 to $h_0 h_4 A$.*

PROOF. Table 21 shows that there is a hidden σ^2 extension from h_5^2 to $h_0 h_4 A$. Lemma 7.156 implies that there also must be a hidden κ extension from h_5^2 to $h_0 h_4 A$. \square

LEMMA 7.150. *There is a hidden ρ_{15} extension from h_5^2 to either $h_0 x_{77,7}$ or $\tau^2 m_1$.*

PROOF. Table 11 shows that the Toda bracket $\langle 8, 2\sigma, \sigma \rangle$ contains ρ_{15} . Then ρ_{15} is also contained in $\langle 2, 8\sigma, \sigma \rangle$, although the indeterminacy increases.

Now shuffle to obtain

$$\rho_{15} \theta_5 = \theta_5 \langle 2, 8\sigma, \sigma \rangle = \langle \theta_5, 2, 8\sigma \rangle \sigma.$$

Table 11 shows that $h_0^3 h_3 h_6$ detects $\langle \theta_5, 2, 8\sigma \rangle$. Also, there is a σ extension from $h_0^3 h_3 h_6$ to $h_0 x_{77,7}$ in the homotopy of $C\tau$.

This implies that $\rho_{15} \theta_5$ is non-zero in $\pi_{77,40}$, and that it is detected in filtration at most 8. Moreover, it is detected in filtration at least 7, since ρ_{15} and θ_5 are detected in filtrations 4 and 2 respectively.

There are several elements in filtration 7 that could detect $\rho_{15} \theta_5$. The element $x_{77,7}$ (if it survives to the E_∞ -page) is ruled out by comparison to $C\tau$. The element $\tau h_1 x_{76,6}$ is ruled out because $\eta \rho_{15} \theta_5$ is detected in filtration at least 10, since $\eta \rho_{15}$ is detected in filtration 7.

The only remaining possibilities are $h_0 x_{77,7}$ and $\tau^2 m_1$. \square

LEMMA 7.151.

- (1) *There is a hidden ρ_{15} extension from $h_1 h_6$ to $Ph_6 c_0$.*
- (2) *There is a hidden ρ_{23} extension from $h_1 h_6$ to $P^2 h_6 c_0$.*

PROOF. Table 11 shows that $h_1 h_6$ detects $\langle \eta, 2, \theta_5 \rangle$. Then

$$\rho_{15} \langle \eta, 2, \theta_5 \rangle \subseteq \langle \eta \rho_{15}, 2, \theta_5 \rangle.$$

The last bracket is detected by $Ph_6 c_0$ because $d_2(h_6) = h_0 h_5^2$ and because Pc_0 detects $\eta \rho_{15}$. Also, its indeterminacy is in Adams filtration greater than 8. This establishes the first hidden extension.

The proof for the second extension is essentially the same, using that $P^2 c_0$ detects $\eta \rho_{23}$ and that the indeterminacy of $\langle \eta \rho_{23}, 2, \theta_5 \rangle$ is in Adams filtration greater than 12. \square

LEMMA 7.152. *There is a hidden ϵ extension from τMg to Md_0^2 .*

PROOF. First, we have the relation $c_0 \cdot h_1^2 X_2 = Mh_1 h_3 g$ in the Adams E_2 -page, which is detected in the homotopy of $C\tau$. Table 14 shows that there are hidden τ extensions from $h_1^2 X_2$ and $Mh_1 h_3 g$ to $\tau M g$ and $M d_0^2$ respectively. \square

LEMMA 7.153. *If $M\Delta h_1^2 d_0$ is non-zero in the E_∞ -page, then there is a hidden ϵ extension from $M\Delta h_1 h_3$ to $M\Delta h_1^2 d_0$.*

PROOF. Table 17 shows that Mh_1 detects the product $\eta\theta_{4.5}$. Since $M\Delta h_1^2 d_0$ equals $\Delta h_1 d_0 \cdot Mh_1$, it detects $\eta\{\Delta h_1 d_0\}\theta_{4.5}$.

Table 21 shows that $\eta\{\Delta h_1 d_0\}$ equals $\epsilon\{\Delta h_1 h_3\}$, since they are both detected by $\Delta h_1^2 d_0$ and there are no elements in higher Adams filtration. Therefore, the product $\epsilon\{\Delta h_1 h_3\}\theta_{4.5}$ is detected by $M\Delta h_1^2 d_0$. In particular, $\{\Delta h_1 h_3\}\theta_{4.5}$ is non-zero, and it can only be detected by $M\Delta h_1 h_3$. \square

7.6. Additional relations

LEMMA 7.154. *The product $(\eta\sigma + \epsilon)\theta_5$ is detected by $\tau h_0 h_2 Q_3$.*

PROOF. Table 11 indicates a hidden η extension from $h_2^2 h_6$ to $\tau h_0 h_2 Q_3$. Therefore, there exists an element α in $\pi_{69,36}$ such that $\tau h_0 h_2 Q_3$ detects $\eta\alpha$. (Beware of the crossing extension from p' to $h_1 p'$. This means that it is possible to choose such an α , but not any element detected by $h_2^2 h_6$ will suffice.)

Table 11 shows that $h_2^2 h_6$ detects the Toda bracket $\langle \nu^2, 2, \theta_5 \rangle$. Let β be an element of this Toda bracket. Since α and β are both detected by $h_2^2 h_6$, the difference $\alpha - \beta$ is detected in Adams filtration at least 4.

Table 21 shows that p' detects $\sigma\theta_5$, which belongs to the indeterminacy of $\langle \nu^2, 2, \theta_5 \rangle$. Therefore, we may choose β such that the difference $\alpha - \beta$ is detected in filtration at least 9. Since $\eta\alpha$ is detected by $\tau h_0 h_2 Q_3$ in filtration 7, it follows that $\eta\beta$ is also detected by $\tau h_0 h_2 Q_3$.

We now have an element β contained in $\langle \nu^2, 2, \theta_5 \rangle$ such that $\eta\beta$ is detected by $\tau h_0 h_2 Q_3$. Now consider the shuffle

$$\eta\langle \nu^2, 2, \theta_5 \rangle = \langle \eta, \nu^2, 2 \rangle \theta_5.$$

Table 11 shows that the last bracket equals $\{\epsilon, \epsilon + \eta\sigma\}$. Therefore, either $\epsilon\theta_5$ or $(\epsilon + \eta\sigma)\theta_5$ is detected by $\tau h_0 h_2 Q_3$. But $\epsilon\theta_5$ is detected by $h_1 p' = h_5^2 c_0$. \square

LEMMA 7.155. *There exists an element α of $\pi_{67,36}$ that is detected by $h_0 Q_3 + h_0 n_1$ such that $\tau\nu\alpha$ equals $(\eta\sigma + \epsilon)\theta_5$.*

PROOF. Lemma 7.154 shows that $\tau h_0 h_2 Q_3$ detects $(\epsilon + \eta\sigma)\theta_5$. The element $\tau h_0 h_2 Q_3$ also detects $\tau\nu\alpha$. Let β be the difference $\tau\nu\alpha - (\epsilon + \eta\sigma)\theta_5$, which is detected in higher Adams filtration. We will show that β must equal zero.

First, $\tau h_1 D_3'$ cannot detect β because $\eta^2\beta$ is zero, while Table 17 shows that $\tau^3 d_1 g^2$ detects $\eta^2\{\tau h_1 D_3'\}$. Second, Table 19 shows that $\tau h_1 h_3(\Delta e_1 + C_0)$ is the target of a hidden ν extension. Therefore, we may alter the choice of α to ensure that β is not detected by $\tau h_1 h_3(\Delta e_1 + C_0)$. Third, $\Delta^2 h_2 c_1$ is also the target of a ν extension. Therefore, we may alter the choice of α to ensure that β is not detected by $\Delta^2 h_2 c_1$. Finally, comparison to tmf implies that β is not detected by $\tau\Delta^2 h_1^2 g + \tau^3 \Delta h_2^2 g^2$. \square

LEMMA 7.156. *The product $(\sigma^2 + \kappa)\theta_5$ is zero.*

PROOF. Table 11 shows that h_5^2 detects the Toda bracket $\langle 2, \theta_4, \theta_4, 2 \rangle$. Shuffle to obtain

$$\langle 2, \theta_4, \theta_4, 2 \rangle (\sigma^2 + \kappa) = 2 \langle \theta_4, \theta_4, 2, \sigma^2 + \kappa \rangle.$$

The second bracket is well-defined because $\langle \theta_4, 2, \sigma^2 + \kappa \rangle$ contains zero [46, ?]. Next, $\langle \theta_4, \theta_4, 2, \sigma^2 + \kappa \rangle 2$ intersects $\langle \theta_4, \theta_4, \alpha \rangle$ for some α in $\langle 2, \sigma^2 + \kappa, 2 \rangle$.

The Toda bracket $\langle 2, \sigma^2 + \kappa, 2 \rangle$ equals

$$\langle 2, \sigma^2, 2 \rangle + \langle 2, \kappa, 2 \rangle,$$

which equals $\tau\eta\kappa$ by Corollary 6.2 since $\eta\sigma^2$ equals zero. The indeterminacy is generated by $2\rho_{15}$. Thus, α equals $\tau\eta\kappa + 2k\rho_{15}$ for some integer k . In every case, the bracket $\langle \theta_4, \theta_4, \alpha \rangle$ has no indeterminacy, so it equals

$$\langle \theta_4, \theta_4, \kappa \rangle \tau\eta + \langle \theta_4, \theta_4, 2 \rangle k\rho_{15}.$$

The second term is zero since $\langle \theta_4, \theta_4, 2 \rangle$ is contained in $\pi_{61,32} = 0$. The first bracket $\langle \theta_4, \theta_4, \kappa \rangle$ is zero by Lemma 6.20. \square

LEMMA 7.157. $\eta\sigma\{k_1\} + \nu\{d_1e_1\} = 0$ in $\pi_{73,41}$.

PROOF. We have the relation $h_1h_3k_1 + h_2d_1e_1 = 0$ in the Adams E_∞ -page, but $\eta\sigma\{k_1\} + \nu\{d_1e_1\}$ could possibly be detected in higher Adams filtration. However, it cannot be detected by h_2^2C'' or Mh_1h_3g by comparison to $C\tau$. Also, it cannot be detected by $\Delta h_1d_0e_0^2$ by comparison to mmf . \square

CHAPTER 8

Tables

Table 1 gives some notation for elements in $\pi_{*,*}$. The fourth column gives partial information that reduces the indeterminacies in the definitions, but does not completely specify a unique element in all cases. See Section 1.5 for further discussion.

Table 2 gives hidden values of the unit map $\pi_{*,*} \rightarrow \pi_{*,*}mmf$. The elements in the third column belong to the Adams E_∞ -page for mmf [14] [17]. See Section 2.2 for further discussion.

Table 3 lists information about some Massey products. The fifth column indicates the proof. When a differential appears in this column, it indicates the May differential that can be used with the May Convergence Theorem (see Remark 2.26) to compute the bracket. The sixth column shows where each specific Massey product is used in the manuscript. See Chapter 4 for more discussion.

Table 4 lists all of the multiplicative generators of the Adams E_2 -page through the 95-stem. The third column indicates the value of the d_2 differential, if it is non-zero. A blank entry in the third column indicates that the d_2 differential is zero. The fourth column indicates the proof. A blank entry in the fourth column indicates that there are no possible values for the differential. The fifth column gives alternative names for the element, as used in [7], [16], or [37]. See Sections 1.5 and 5.1 for further discussion.

Table 5 lists some elements in the Adams spectral sequence that are known to be permanent cycles. The third column indicates the proof. When a Toda bracket appears in the third column, the Moss Convergence Theorem 2.16 applied to that Toda bracket implies that the element is a permanent cycle (see Table 11 for more information). When a product appears in the third column, the element must survive to detect that product.

Table 6 lists the multiplicative generators of the Adams E_3 -page through the 95-stem whose d_3 differentials are non-zero, or whose d_3 differentials are zero for non-obvious reasons. See Section 5.2 for further discussion.

Table 7 lists the multiplicative generators of the Adams E_4 -page through the 95-stem whose d_4 differentials are non-zero, or whose d_4 differentials are zero for non-obvious reasons. See Section 5.3 for further discussion.

Table 8 lists the multiplicative generators of the Adams E_5 -page through the 95-stem whose d_5 differentials are non-zero, or whose d_5 differentials are zero for non-obvious reasons. See Section 5.4 for further discussion.

Table 9 lists the multiplicative generators of the Adams E_r -page, for $r \geq 6$, through the 90-stem whose d_r differentials are non-zero, or whose d_r differentials are zero for non-obvious reasons. See Section 5.5 for further discussion.

Table 10 lists the possible Adams differentials that remain unresolved. These possibilities also appear in Tables 4–9.

Table 11 lists information about some Toda brackets. The third column of Table 11 gives an element of the Adams E_∞ -page that detects an element of the Toda bracket. The fourth column of Table 11 gives partial information about indeterminacies, again by giving detecting elements of the Adams E_∞ -page. We have not completely analyzed the indeterminacies of some brackets when the details are inconsequential for our purposes; this is indicated by a blank entry in the fourth column. The fifth column indicates the proof of the Toda bracket, and the sixth column shows where each specific Toda bracket is used in the manuscript. See Chapter 6 for further discussion.

Tables 12 and 13 gives hidden values of the inclusion $\pi_{*,*} \rightarrow \pi_{*,*}C\tau$ of the bottom cell, and of the projection $\pi_{*,*}C\tau \rightarrow \pi_{*-1,*,+1}$ to the top cell. See Section 7.1 for further discussion.

Table 14 lists hidden τ extensions in the E_∞ -page of the \mathbb{C} -motivic Adams spectral sequence. See Section 7.1 for further discussion.

Tables 15, 17, and 19 list hidden extensions by 2 , η , and ν . The fourth column indicates the proof of each extension. The fifth column gives additional information about each extension, including whether it is a crossing extension and whether it has indeterminacy in the sense of Section 2.1.1. See Sections 7.2, 7.3, and 7.4 for further discussion.

Tables 16, 18, and 20 list possible hidden extensions by 2 , η , and ν that we have not yet resolved.

Finally, Table 21 gives some various hidden extensions by elements other than 2 , η , and ν . See Section 7.5 for further discussion.

Table 1: Notation for $\pi_{*,*}$

(s, w)	element	detected by	definition
$(0, -1)$	τ	τ	
$(0, 0)$	2	h_0	
$(1, 1)$	η	h_1	
$(3, 2)$	ν	h_2	
$(7, 4)$	σ	h_3	
$(8, 5)$	ϵ	c_0	
$(9, 5)$	μ_9	Ph_1	
$(14, 8)$	κ	d_0	
$(15, 8)$	ρ_{15}	$h_0^3 h_4$	
$(16, 9)$	η_4	$h_1 h_4$	
$(17, 9)$	μ_{17}	$P^2 h_1$	
$(19, 11)$	$\bar{\sigma}$	c_1	
$(20, 11)$	$\bar{\kappa}$	τg	$\langle \kappa, 2, \eta, \nu \rangle$
$(23, 12)$	ρ_{23}	$h_0^2 i + \tau P h_1 d_0$	
$(25, 13)$	μ_{25}	$P^3 h_1$	
$(30, 16)$	θ_4	h_4^2	
$(32, 17)$	η_5	$h_1 h_5$	in $\langle \eta, 2, \theta_4 \rangle$
$(32, 18)$	κ_1	d_1	
$(44, 24)$	$\bar{\kappa}_2$	g_2	
$(45, 24)$	$\theta_{4.5}$	h_4^3	$\eta \theta_{4.5} \in \{Mh_1\}$
$(62, 32)$	θ_5	h_5^2	
$(63, 32)$	η_6	$h_1 h_6$	in $\langle \eta, 2, \theta_5 \rangle$

Table 2: Some hidden values of the unit map of mmf

(s, f, w)	element	image
(28, 6, 17)	$h_1 h_3 g$	cg
(29, 7, 18)	$h_1^2 h_3 g$	$h_1 cg$
(32, 6, 17)	$\Delta h_1 h_3$	$\Delta c + \tau ag$
(33, 7, 18)	$\Delta h_1^2 h_3$	$\Delta h_1 c$
(35, 8, 18)	$\tau^3 h_1 e_0^2$	Pan
(40, 10, 21)	$\tau \Delta h_1^2 d_0$	$P(\Delta c + \tau ag)$
(48, 10, 29)	$h_1 h_3 g^2$	cg^2
(49, 11, 30)	$h_1^2 h_3 g^2$	$h_1 cg^2$
(52, 10, 29)	$\Delta h_1 h_3 g$	$(\Delta c + \tau ag)g$
(53, 11, 30)	$\Delta h_1^2 h_3 g$	$\Delta h_1 cg$
(54, 9, 28)	$h_0 h_5 i$	$\Delta^2 h_2^2$
(54, 11, 32)	$h_1^6 h_5 e_0$	dg^2
(55, 12, 33)	$h_1^7 h_5 e_0$	$h_1 dg^2$
(57, 10, 30)	$h_0 h_2 h_5 i$	$\Delta h_1 (\Delta c + \tau ag)$
(59, 12, 33)	$Ph_1^3 h_5 e_0$	$\Delta h_1 dg$
(60, 13, 34)	$\tau^2 h_0 g^3$	$\Delta h_2^2 dg$
(62, 14, 37)	$h_1^6 h_5 c_0 e_0$	cdg^2
(63, 15, 38)	$h_1^7 h_5 c_0 e_0$	$h_1 cdg^2$
(65, 12, 34)	$Ph_5 j$	$\Delta^2 h_2 d$
(66, 14, 37)	$Ph_1^2 h_5 c_0 e_0$	$(\Delta c + \tau ag)dg$
(67, 15, 38)	$Ph_1^3 h_5 c_0 e_0$	$\Delta h_1 cdg$
(68, 13, 37)	$Ph_2 h_5 j$	$\Delta^2 h_2^2 d$
(68, 14, 41)	$h_1 h_3 g^3$	cg^3
(69, 15, 42)	$h_1^2 h_3 g^3$	$h_1 cg^3$
(71, 15, 38)	$\Delta^2 h_0^2 h_2 g$	$\Delta h_1 (\Delta c + \tau ag)d$
(72, 14, 41)	$\Delta h_1 h_3 g^2$	$(\Delta c + \tau ag)g^2$
(73, 15, 42)	$\Delta h_1^2 h_3 g^2$	$\Delta h_1 cg^2$
(77, 15, 42)	$\Delta^2 h_2^3 g$	$\Delta h_1 (\Delta c + \tau ag)g$
(88, 18, 53)	$h_1 h_3 g^4$	cg^4
(89, 19, 54)	$h_1^2 h_3 g^4$	$h_1 cg^4$

Table 3: Some Massey products in $\text{Ext}_{\mathbb{C}}$

(s, f, w)	bracket	contains	indeterminacy	proof	used for
(2, 2, 1)	$\langle h_0, h_1, h_0 \rangle$	τh_1^2	0	Theorem 4.1	$\langle 2, \eta, 2 \rangle$
(3, 2, 2)	$\langle h_1, h_0, h_1 \rangle$	$h_0 h_2$	0	Theorem 4.1	$\langle \{h_1 x_{76,6}\}, 2, \eta \rangle, 7.71$
(6, 2, 4)	$\langle h_1, h_2, h_1 \rangle$	h_2^2	0	Theorem 4.1	$\langle \eta, \nu, \eta \rangle$
(8, 2, 5)	$\langle h_2, h_1, h_2 \rangle$	$h_1 h_3$	0	Theorem 4.1	$\langle \nu, \eta, \nu \rangle$
(8, 3, 5)	$\langle h_1^2, h_0, h_1, h_2 \rangle$	c_0	0	$d_1(h_{20}) = h_0 h_1$ $d_1(h_{21}) = h_1 h_2$	$\langle \eta^2, 2, \eta, \nu \rangle$
(8, 3, 5)	$\langle h_1, h_2, h_0 h_2 \rangle$	c_0	0	$d_2(h_0(1)) = h_0 h_2^2$	$\langle \eta, \nu, 2\nu \rangle, \langle \eta_5, \nu, 2\nu \rangle$
(9, 5, 5)	$\langle h_1, h_0, h_0^3 h_3 \rangle$	Ph_1	0	$d_4(b_{20}^2) = h_0^4 h_3$	$\langle \eta, 2, 8\sigma \rangle$
(11, 5, 6)	$\langle h_0, h_1, \tau h_1 c_0 \rangle$	Ph_2	0	$d_2(b_{20} h_0(1)) = \tau h_1^2 c_0$	$\langle 2, \eta, \tau \eta \epsilon \rangle$
(20, 4, 11)	$\langle \tau, h_1^4, h_4 \rangle$	τg	0	$d_4(g) = h_1^4 h_4$	7.130
(23, 9, 12)	$\langle h_3, h_0^4, h_0^4 h_4 \rangle$	$h_0^2 i + \tau P h_1 d_0$	0	$d_8(b_{20}^4) = h_0^8 h_4$	$\langle \sigma, 16, 2\rho_{15} \rangle$
(30, 6, 16)	$\langle h_0^2, h_3^2, h_0^2, h_3^2 \rangle$	Δh_2^2	0	$d_4(h_2 b_{30}) = h_0^2 h_3^2$	4.13
(32, 4, 18)	$\langle h_1, h_3^2, h_1, h_3^2 \rangle$	d_1	0	$d_2(h_1(1)) = h_1 h_3^2$	$\langle \eta, \sigma^2, \eta, \sigma^2 \rangle$
(33, 4, 18)	$\langle h_1 h_4^2, h_1, h_0 \rangle$	p	0	$d_4(h_3 b_{31}) = h_1^2 h_4^2$	$\langle \eta \theta_4, \eta, 2 \rangle$
(46, 7, 25)	$\langle h_1, h_0, h_0^2 g_2 \rangle$	$M h_1$	0	M	6.10
(54, 6, 29)	$\langle h_1, h_0, D_1 \rangle$	$\tau \Delta_1 h_1^2$	0	[22]	Example 2.31
(66, 6, 36)	$\langle h_1^2, h_4^2, h_1^2, h_4^2 \rangle$	$\Delta_1 h_3^2$	0	Lemma 4.13	6.15
(67, 14, 36)	$\langle P d_0, h_0^3, g_2 \rangle$	$M P d_0$	0	M	5.66
(68, 11, 38)	$\langle h_2 g, h_0^3, g_2 \rangle$	$M h_2 g$	0	M	5.39
(71, 13, 40)	$\langle h_1^3 h_4, h_1, \tau g n \rangle$	$\tau g^2 n$	$M h_0 h_2^2 g$	Lemma 4.8	5.40
(75, 18, 42)	$\langle \Delta h_0^2 d_0 e_0, h_1, h_1^3 h_4 \rangle$	$\Delta h_2^2 d_0^2 e_0$	0	Proposition 4.2	$\langle \tau \eta \kappa \bar{\kappa}^2, \eta, \eta^2 \eta_4 \rangle$
(80, 5, 42)	$\langle h_3, p', h_2 \rangle$	$h_0 e_2$	0	Lemma 4.11	7.123
(82, 12, 45)	$\langle \Delta e_1 + C_0, h_1^3, h_1 h_4 \rangle$	$(\Delta e_1 + C_0) g$	0	Lemma 4.10	$\langle \{\Delta e_1 + C_0\}, \eta^3, \eta_4 \rangle$
(86, 16, 46)	$\langle \Delta h_0^2 e_0, h_0^2, h_0 g_2 \rangle$	$M \Delta h_0^2 e_0$	0	M	6.25
(91, 13, 49)	$\langle M h_1, h_0, h_0^2 g_2 \rangle$	$M^2 h_1$	0	M	5.28
(93, 13, 49)	$\langle \tau^3 g G_0, h_0 h_2, h_2 \rangle$	$? \tau e_0 x_{76,9}$	$\tau M^2 h_2$	Lemma 4.12	5.60

Table 4: E_2 -page generators of the \mathbb{C} -motivic Adams spectral sequence

(s, f, w)	element	d_2	proof	other names
(0, 1, 0)	h_0			
(1, 1, 1)	h_1			
(3, 1, 2)	h_2			
(7, 1, 4)	h_3			
(8, 3, 5)	c_0			
(9, 5, 5)	Ph_1			
(11, 5, 6)	Ph_2			
(14, 4, 8)	d_0			
(15, 1, 8)	h_4	$h_0h_3^2$	$C\tau$	
(16, 7, 9)	Pc_0			
(17, 4, 10)	e_0	$h_1^2d_0$	$C\tau$	
(17, 9, 9)	P^2h_1			
(18, 4, 10)	f_0	$h_0^2e_0$	$C\tau$	
(19, 3, 11)	c_1		$C\tau$	
(19, 9, 10)	P^2h_2			
(20, 4, 11)	τg			
(22, 8, 12)	Pd_0			
(23, 5, 14)	h_2g			
(23, 7, 12)	i	Ph_0d_0	$C\tau$	
(24, 11, 13)	P^2c_0			
(25, 8, 14)	Pe_0	$Ph_1^2d_0$	$C\tau$	
(25, 13, 13)	P^3h_1			
(26, 7, 14)	j	Ph_0e_0	$C\tau$	
(27, 5, 16)	h_3g	$h_0h_2^2g$	$C\tau$	
(27, 13, 14)	P^3h_2			
(29, 7, 16)	k	$h_0d_0^2$	$C\tau$	
(30, 6, 16)	Δh_2^2		$C\tau$	r
(30, 12, 16)	P^2d_0			
(31, 1, 16)	h_5	$h_0h_4^2$	$C\tau$	
(31, 5, 17)	n			
(32, 4, 18)	d_1			
(32, 6, 17)	Δh_1h_3		h_0	q
(32, 7, 18)	l	$h_0d_0e_0$	$C\tau$	
(32, 15, 17)	P^3c_0			
(33, 4, 18)	p			
(33, 12, 18)	P^2e_0	$P^2h_1^2d_0$	$C\tau$	
(33, 17, 17)	P^4h_1			
(34, 11, 18)	Pj	$P^2h_0e_0$	$C\tau$	
(35, 7, 20)	m	$h_0e_0^2$	$C\tau$	
(35, 17, 18)	P^4h_2			
(36, 6, 20)	t		$C\tau$	
(37, 5, 20)	x			
(37, 8, 22)	e_0g	$h_1^2e_0^2$	$C\tau$	

Table 4: E_2 -page generators of the \mathbb{C} -motivic Adams spectral sequence

(s, f, w)	element	d_2	proof	other names
(38, 4, 21)	e_1			
(38, 6, 20)	Δh_3^2	$h_0^3 x$	$C\tau$	$y + ?h_2^2 d_1$
(38, 16, 20)	$P^3 d_0$			
(39, 7, 23)	$c_1 g$		$C\tau$	
(39, 9, 21)	$\Delta h_1 d_0$			u
(39, 15, 20)	$P^2 i$	$P^3 h_0 d_0$	$C\tau$	
(40, 4, 22)	f_1			
(40, 8, 23)	τg^2			
(40, 19, 21)	$P^4 c_0$			
(41, 3, 22)	c_2	$h_0 f_1$	$C\tau$	
(41, 10, 22)	$\Delta h_0^2 e_0$		$C\tau$	z
(41, 16, 22)	$P^3 e_0$	$P^3 h_1^2 d_0$	$C\tau$	
(41, 21, 21)	$P^5 h_1$			
(42, 9, 23)	$\Delta h_1 e_0$	$\Delta h_1^3 d_0$	$C\tau$	v
(42, 15, 22)	$P^2 j$	$P^3 h_0 e_0$	$C\tau$	
(43, 9, 26)	$h_2 g^2$			
(43, 21, 22)	$P^5 h_2$			
(44, 4, 24)	g_2			
(45, 9, 24)	$\tau \Delta h_1 g$			w
(46, 7, 25)	$M h_1$			B_1
(46, 8, 25)	$\Delta h_2 c_1$			N
(46, 11, 25)	$\Delta c_0 d_0$	$\tau h_0 d_0^2 e_0$	$C\tau$	u'
(46, 20, 24)	$P^4 d_0$			
(47, 9, 28)	$h_3 g^2$	$h_0 h_2^2 g^2$	$C\tau$	
(47, 13, 24)	$\Delta h_0^2 i$	$h_0 i^2$	$C\tau$	$Q', Q + Pu$
(47, 13, 25)	$P \Delta h_1 d_0$			
(48, 7, 26)	$M h_2$			$B_2 + ?h_0^2 h_5 e_0$
(48, 23, 25)	$P^5 c_0$			
(49, 11, 27)	$\Delta c_0 e_0$	$\Delta h_1^2 c_0 d_0 + \tau h_0 d_0 e_0^2$	$C\tau$	v'
(49, 20, 26)	$P^4 e_0$	$P^4 h_1^2 d_0$	$C\tau$	
(49, 25, 25)	$P^6 h_1$			
(50, 6, 27)	C			
(50, 10, 28)	$\Delta h_2^2 g$		$C\tau$	
(50, 13, 27)	$P \Delta h_1 e_0$	$P \Delta h_1^3 d_0$	$C\tau$	
(50, 19, 26)	$P^3 j$	$P^4 h_0 e_0$	$C\tau$	
(51, 9, 28)	$\Delta h_3 g$	$\Delta h_0 h_2^2 g$	$C\tau$	G_3
(51, 9, 29)	gn			
(51, 25, 26)	$P^6 h_2$			
(52, 5, 28)	D_1	$h_0^2 h_3 g_2$	$C\tau$	
(52, 8, 30)	$d_1 g$			
(53, 7, 30)	i_1		$C\tau$	
(53, 9, 29)	$M c_0$		h_0	$B_8, Ph_5 d_0$
(53, 10, 28)	MP			x'
(54, 6, 29)	$\tau \Delta_1 h_1^2$	$M h_1 h_3$	$C\tau$	G

Table 4: E_2 -page generators of the \mathbb{C} -motivic Adams spectral sequence

(s, f, w)	element	d_2	proof	other names
(54, 10, 28)	$\Delta^2 h_2^2$	MPh_0^2	$C\tau$	$R_1 + ?h_0^2 h_5 i$
(54, 15, 29)	$P\Delta c_0 d_0$	$\tau Ph_0 d_0^2 e_0$	$C\tau$	
(54, 24, 28)	$P^5 d_0$			
(55, 7, 30)	B_6		$C\tau$	
(55, 11, 32)	gm	$h_0 e_0^2 g$	$C\tau$	
(55, 17, 29)	$P^2 \Delta h_1 d_0$			
(55, 23, 28)	$P^4 i$	$P^5 h_0 d_0$	$C\tau$	
(56, 10, 29)	$\Delta^2 h_1 h_3$	τMPh_1^2	h_1	$Q_1 + ?gt$
(56, 10, 32)	gt		$C\tau$	
(56, 27, 29)	$P^6 c_0$			
(57, 6, 31)	D_4	$h_1 B_6$	$C\tau$	
(57, 7, 30)	Q_2		$C\tau$	
(57, 9, 31)	$\Delta h_1 d_1$			D_{11}
(57, 15, 31)	$P\Delta c_0 e_0$	$P\Delta h_1^2 c_0 d_0 + \tau h_0 d_0^4$	$C\tau$	
(57, 24, 30)	$P^5 e_0$	$P^5 h_1^2 d_0$	$C\tau$	
(57, 29, 29)	$P^7 h_1$			
(58, 6, 30)	D_2	$h_0 Q_2$	$C\tau$	
(58, 8, 33)	$e_1 g$			
(58, 17, 31)	$P^2 \Delta h_1 e_0$	$P^2 \Delta h_1^3 d_0$	$C\tau$	
(58, 23, 30)	$P^4 j$	$P^5 h_0 e_0$	$C\tau$	
(59, 7, 33)	j_1			
(59, 10, 32)	Md_0			B_{21}
(59, 11, 35)	$c_1 g^2$			
(59, 29, 30)	$P^7 h_2$			
(60, 7, 32)	Mh_4		$C\tau$	B_3
(60, 9, 32)	B_4	$Mh_0 d_0$	$C\tau$	
(60, 12, 35)	τg^3			
(60, 13, 36)	$h_0 g^3$			
(61, 4, 32)	D_3			
(61, 6, 32)	A'		$C\tau$	
(61, 6, 32)	$A + A'$	$Mh_0 h_4$	$C\tau$	
(61, 7, 33)	B_7			
(61, 9, 32)	Δx	$h_0^2 B_4 + \tau Mh_1 d_0$	Lemma 5.3 or $C\tau, h_1^2$	X_1
(62, 5, 33)	H_1	B_7	$C\tau$	
(62, 8, 33)	C_0			$x_{8,33} + h_0^6 h_5^2$
(62, 8, 33)	Δe_1			$E_1, x_{8,32} + x_{8,33}$
(62, 10, 32)	$\Delta^2 h_3^2$		$C\tau$	$x_{10,27} + x_{10,28} +$ $+ h_1 X_1, R$
(62, 10, 34)	Me_0	$Mh_1^2 d_0$	$C\tau$	$B_{22}, x_{10,28}$
(62, 19, 33)	$P^2 \Delta c_0 d_0$	$\tau P^2 h_0 d_0^2 e_0$	$C\tau$	
(62, 28, 32)	$P^6 d_0$			
(63, 1, 32)	h_6	$h_0 h_5^2$	$C\tau$	
(63, 7, 34)	C'		$C\tau$	$x_{7,33} + x_{7,44}$

Table 4: E_2 -page generators of the \mathbb{C} -motivic Adams spectral sequence

(s, f, w)	element	d_2	proof	other names
(63, 7, 34)	X_2	$Mh_1^2h_4$	$C\tau$	$x_{7,33}$
(63, 13, 38)	h_2g^3			
(63, 21, 33)	$P^3\Delta h_1d_0$			
(64, 6, 34)	A''	h_0X_2	$C\tau$	
(64, 10, 33)	$\Delta^2h_1h_4$			$x_{10,32}+?h_0^2h_3Q_2,$ q_1
(64, 14, 34)	$\Delta^2h_1^2d_0$	$MP^2h_1^2$	$C\tau$	$U, PQ_1+?km$
(64, 31, 33)	P^7c_0			
(65, 7, 36)	k_1			
(65, 10, 35)	τMg		$C\tau$	B_{23}
(65, 13, 34)	$\Delta^2h_0e_0$	$h_0 \cdot \Delta^2h_1^2d_0$	$C\tau$	$R_2+?gw$
(65, 19, 35)	$P^2\Delta c_0e_0$	$P^2\Delta h_1^2c_0d_0+$ $+\tau Ph_0d_0^4$	$C\tau$	
(65, 28, 34)	P^6e_0	$P^6h_1^2d_0$	$C\tau$	
(65, 33, 33)	P^8h_1			
(66, 6, 36)	$\Delta_1h_3^2$			r_1
(66, 7, 35)	τG_0	$h_2C_0 + h_1h_3Q_2$	$C\tau$	$x_{7,40}+?h_0r_1$
(66, 10, 34)	D'_2	$\tau^2Mh_0^2g$	$C\tau, i$	PD_2
(66, 10, 35)	τB_5	τMh_0^2g	i	
(66, 21, 35)	$P^3\Delta h_1e_0$	$P^3\Delta h_1^3d_0$	$C\tau$	
(66, 27, 34)	P^5j	$P^6h_0e_0$	$C\tau$	
(67, 5, 35)	τQ_3	$\tau\Delta_1h_0h_3^2$	g_2	
(67, 5, 36)	n_1	$\Delta_1h_0h_3^2$	$C\tau$	
(67, 6, 36)	$h_0Q_3+h_2^2D_3$		$C\tau$	
(67, 9, 36)	X_3		$C\tau$	$x_{9,40}$
(67, 9, 37)	C''			$x_{9,39}$
(67, 11, 35)	Δ^2c_1			$x_{11,35} + h_0^2x_{9,40}$
(67, 13, 40)	h_3g^3	$h_0h_2^2g^3$	$C\tau$	
(67, 33, 34)	P^8h_2			
(68, 4, 36)	d_2		$C\tau$	
(68, 8, 36)	Δg_2	h_0X_3	$C\tau$	$G_{21}+?h_0h_3A'$
(68, 11, 38)	Mh_2g			
(68, 13, 36)	Δ^2h_0g	MPh_0d_0	$C\tau$	$P^2D_1+?h_0^5G_{21},$ G_{11}
(69, 4, 36)	p'		$C\tau$	
(69, 8, 37)	D'_3	$h_1X_3+?th_1C''$	$C\tau$	PD_3
(69, 8, 38)	h_2G_0	h_1C''	$C\tau$	
(69, 10, 36)	$P(A + A')$	$\tau^2Mh_0h_2g$	$C\tau, i$	
(69, 11, 38)	h_2B_5		$C\tau$	
(69, 13, 36)	$\tau\Delta^2h_1g$		d_0	W_1
(69, 18, 36)	MP^3			$x_{18,20}+?d_0il$
(70, 4, 37)	p_1		$C\tau$	
(70, 6, 38)	h_2Q_3		$C\tau$	
(70, 17, 36)	$P\Delta^2h_0d_0$	MP^3h_0	$C\tau$	$R'_1, R_1+?d_0^2v$

Table 4: E_2 -page generators of the \mathbb{C} -motivic Adams spectral sequence

(s, f, w)	element	d_2	proof	other names
(70, 23, 37)	$P^3\Delta c_0d_0$	$\tau P^3h_0d_0^2e_0$	$C\tau$	
(70, 32, 36)	P^7d_0			
(71, 6, 38)	$x_{71,6}$	τd_1e_1	h_3	$x_{6,47}+?h_1^2p'$
(71, 7, 39)	l_1			
(71, 12, 37)	$\Delta^2h_4c_0$		$C\tau, tmf, h_0$	$x_{12,37}+?h_0d_0Q_2$
(71, 13, 38)	Δ^2h_2g		$C\tau$	$x_{13,34}$
(71, 13, 38)	Mj	MPh_0e_0	$C\tau$	$x_{13,35}$
(71, 13, 40)	Δh_3g^2	h_0m^2	$C\tau$	
(71, 13, 41)	g^2n			
(71, 25, 37)	$P^4\Delta h_1d_0$			
(71, 31, 36)	P^6i	$P^7h_0d_0$	$C\tau$	
(72, 12, 42)	d_1g^2			
(72, 17, 39)	$\Delta^2h_1^2c_0d_0$	$MP^2h_1^2c_0$	$C\tau$	
(72, 18, 38)	$P\Delta^2h_1^2d_0$	$MP^3h_1^2$	$C\tau$	
(72, 35, 37)	P^8c_0			
(73, 17, 38)	$P\Delta^2h_0e_0$	MP^3h_2	$C\tau$	
(73, 23, 39)	$P^3\Delta c_0e_0$	$P^3\Delta h_1^2c_0d_0+$ $+\tau P^2h_0d_0^4$	$C\tau$	
(73, 32, 38)	P^7e_0	$P^7h_1^2d_0$	$C\tau$	
(73, 37, 37)	P^9h_1			
(74, 8, 40)	$x_{74,8}$			$x_{8,51}+?Ph_0^2h_2h_6$
(74, 25, 39)	$P^4\Delta h_1e_0$	$P^4\Delta h_1^3d_0$	$C\tau$	
(74, 31, 38)	P^6j	$P^7h_0e_0$	$C\tau$	
(75, 7, 40)	$x_{75,7}$		$C\tau$	$x_{7,53}$
(75, 11, 42)	gB_6		$C\tau$	
(75, 13, 40)	Δ^2h_3g	$\Delta^2h_0h_2^2g$	$C\tau$	
(75, 15, 44)	g^2m	$h_0e_0^2g^2$	$C\tau$	
(75, 37, 38)	P^9h_2			
(76, 6, 40)	$x_{76,6}$	$h_0x_{75,7}$	$C\tau$	$x_{6,53}$
(76, 9, 40)	$x_{76,9}$		$C\tau$	$x_{9,51}+?h_1h_4B_3$
(76, 14, 44)	g^2t		$C\tau$	
(76, 16, 40)	$\Delta^2d_0^2$	τ^2d_0jm	$C\tau, tmf$	$x_{16,32}$
(77, 7, 40)	$x_{77,7}$	$\tau Mh_1h_4^2$	Lemma 5.5	$x_{7,57} + m_1$
(77, 7, 42)	m_1		$C\tau$	
(77, 8, 41)	$x_{77,8}$	$\Delta h_1h_3g_2$	$C\tau$	$x_{8,57}$
(77, 12, 41)	$M\Delta h_1h_3$		$C\tau, h_0, \tau g$	P^2D_3
(77, 13, 43)	Δh_1d_1g			
(77, 16, 40)	Δ^2h_0k	$\Delta^2h_0^2d_0^2$	$C\tau$	$x_{16,33}+?e_0g^3$
(77, 16, 46)	e_0g^3	$h_1^2e_0^2g^2$	$C\tau$	
(78, 6, 42)	t_1	h_0m_1	$C\tau$	
(78, 9, 41)	$x_{78,9}$		$C\tau$	$x_{9,55}+?h_0^7h_4h_6$
(78, 10, 40)	$x_{78,10}$	$h_0^5x_{77,7}$	$C\tau, h_1$	$P^2h_5^2$
(78, 12, 45)	e_1g^2			
(78, 27, 41)	$P^4\Delta c_0d_0$	$\tau P^4h_0d_0^2e_0$	$C\tau$	

Table 4: E_2 -page generators of the \mathbb{C} -motivic Adams spectral sequence

(s, f, w)	element	d_2	proof	other names
(78, 36, 40)	$P^8 d_0$			
(79, 5, 42)	x_1		$C\tau$	
(79, 11, 42)	ΔB_6		$C\tau$	
(79, 11, 45)	gj_1			
(79, 13, 41)	$\Delta^2 n$			$x_{13,42}$
(79, 15, 47)	$c_1 g^3$		$C\tau$	
(79, 16, 42)	$\Delta^2 d_0 e_0$	$\Delta^2 h_1^2 d_0^2 + \tau^2 d_0 k m$	$C\tau, tmf$	$x_{16,35}$
(79, 29, 40)	$P^4 \Delta h_0^2 i$	$P^4 h_0 i^2$	$C\tau$	
(79, 29, 41)	$P^5 \Delta h_1 d_0$		$C\tau$	
(80, 4, 42)	e_2	$h_0 x_1$	$C\tau$	
(80, 12, 42)	$\Delta^2 d_1$		h_0	$x_{12,44}$
(80, 13, 44)	gB_4	$Mh_0 e_0^2$	$C\tau$	
(80, 14, 41)	$\Delta^3 h_1 h_3$		h_0	$x_{14,42}$
(80, 16, 42)	$\Delta^2 h_0 l$	$\Delta^2 h_0^2 d_0 e_0$	$C\tau$	$x_{16,37} + ?g^4$
(80, 16, 47)	τg^4			
(80, 22, 42)	$P^2 \Delta^2 h_1^2 d_0$	$MP^4 h_1^2$	$C\tau$	
(80, 39, 41)	$P^9 c_0$			
(81, 10, 44)	gA'			
(81, 11, 45)	gB_7			
(81, 12, 42)	$\Delta^2 p$		h_0	$x_{12,45}$
(81, 21, 42)	$P^2 \Delta^2 h_0 e_0$	$h_0 \cdot P^2 \Delta^2 h_1^2 d_0$	$C\tau$	
(81, 27, 43)	$P^4 \Delta c_0 e_0$	$P^4 \Delta h_1^2 c_0 d_0 +$ $+\tau P^3 h_0 d_0^4$	$C\tau$	
(81, 36, 42)	$P^8 e_0$	$P^8 h_1^2 d_0$	$C\tau$	
(81, 41, 41)	$P^{10} h_1$			
(82, 6, 44)	$h_4 Q_3$			
(82, 9, 45)	gH_1	gB_7	$C\tau$	
(82, 12, 45)	$(\Delta e_1 + C_0)g$			
(82, 12, 45)	gC_0			
(82, 14, 46)	$Me_0 g$	$Mh_1^2 e_0^2$	$C\tau$	
(82, 16, 44)	$\Delta^2 e_0^2$	$\tau^2 \Delta h_2^2 e_0^3$	$C\tau, tmf$	$x_{16,38}$
(82, 17, 47)	$\Delta h_1 e_0 g^2$	$\Delta h_1^3 e_0^2 g + Mh_1^5 d_0 e_0$	$C\tau$	
(82, 29, 43)	$P^5 \Delta h_1 e_0$	$P^5 \Delta h_1^3 d_0$	$C\tau$	
(82, 35, 42)	$P^7 j$	$P^8 h_0 e_0$	$C\tau$	
(83, 11, 45)	Δj_1			
(83, 11, 46)	gC'			
(83, 15, 44)	$\Delta^2 m$	$\Delta^2 h_0 e_0^2 +$ $+\tau^3 \Delta h_1 e_0 g^2$	$C\tau, tmf$	$x_{15,41}$
(83, 17, 50)	$h_2 g^4$			
(83, 41, 42)	$P^{10} h_2$			
(84, 4, 44)	f_2		$C\tau$	
(84, 10, 45)	$Px_{76,6}$		$C\tau$	
(84, 14, 44)	$\Delta^2 t$		$C\tau$	$x_{14,46}$
(84, 15, 44)	$\Delta h_0^2 B_4$	$\Delta^2 h_0^2 m + \tau \Delta^2 h_1 e_0^2$	$C\tau, mmf$	$x_{15,42} + x_{15,43}$

Table 4: E_2 -page generators of the \mathbb{C} -motivic Adams spectral sequence

(s, f, w)	element	d_2	proof	other names
(84, 15, 45)	$M\Delta h_1 d_0$		$C\tau$	$x_{15,43}$
(85, 3, 44)	c_3	$h_0 f_2$	$C\tau$	
(85, 6, 45)	$x_{85,6}$		h_1	$x_{6,68} + h_0^3 c_3$
(85, 13, 44)	$\Delta^2 x$		$C\tau, h_1^2$	$x_{13,46}$
(85, 14, 48)	Mg^2		$C\tau$	
(85, 16, 46)	$\Delta^2 e_0 g$	$\Delta^2 h_1^2 e_0^2 + \tau^2 d_0 m^2$	tmf	$x_{16,42} + h_0^3 x_{13,46}$
(85, 26, 44)	MP^5			
(86, 11, 47)	$\tau g G_0$	$h_2(\Delta e_1 + C_0)g$	$C\tau$	
(86, 12, 45)	$\Delta^2 e_1$		d_0	$x_{12,48}$
(86, 12, 46)	$\Delta h_1 B_7$			
(86, 14, 44)	$\Delta^3 h_3^2$	$\Delta^2 h_0^3 x$	$C\tau, h_1$	$P^3 h_5^2 + ?gB_5$
(86, 14, 47)	$\tau B_5 g$	$\tau M h_0^2 g^2$	Lemma 5.6	
(86, 25, 44)	$\Delta^2 P^3 h_0 d_0$	$MP^5 h_0$	$C\tau$	$x_{25,24} + ?P^2 d_0^2 v$
(86, 31, 45)	$P^5 \Delta c_0 d_0$	$\tau P^5 h_0 d_0^2 e_0$	h_0	
(86, 40, 44)	$P^9 d_0$			
(87, 7, 45)	$x_{87,7}$			$x_{7,74}$
(87, 9, 48)	gQ_3		$C\tau$	
(87, 10, 46)	$\Delta h_1 H_1$	$\Delta h_1 B_7$	$C\tau$	$x_{10,60}$
(87, 13, 49)	gC''		$C\tau$	
(87, 15, 47)	$\Delta^2 c_1 g$		$C\tau$	
(87, 15, 47)	$M\Delta h_1 e_0$	$M\Delta h_1^3 d_0$	$C\tau$	$x_{15,47} + ?h_2 x_{14,46}$
(87, 17, 45)	$\Delta^3 h_1 d_0$	$\tau^5 e_0^3 m$	tmf	$x_{17,50}$
(87, 17, 52)	$h_3 g^4$	$h_0 h_2^2 g^4$	$C\tau$	
(87, 20, 46)	$P\Delta^2 d_0 e_0$	$P\Delta^2 h_1^2 d_0^2 + \tau^2 \Delta h_2^2 d_0^4$	$C\tau, tmf$	
(87, 33, 45)	$P^6 \Delta h_1 d_0$		$C\tau$	
(87, 39, 36)	$P^8 i$	$P^9 h_0 d_0$	$C\tau$	
(88, 10, 48)	$x_{88,10}$			
(88, 12, 46)	$\Delta^2 f_1$			$x_{12,51} + ?gG_{21} + ?Ph_0^3 h_6 e_0$
(88, 12, 48)	$\Delta g_2 g$			
(88, 16, 47)	$\tau \Delta^2 g^2$	$\tau^3 \Delta h_2^2 e_0 g^2$	tmf, h_0	$x_{16,48}$
(88, 26, 46)	$P^3 \Delta^2 h_1^2 d_0$	$MP^5 h_1^2$	$C\tau$	
(88, 43, 45)	$P^{10} c_0$			
(89, 11, 46)	$\Delta^2 c_2$	$\Delta^2 h_0 f_1$	$C\tau$	$x_{11,59}$
(89, 12, 50)	$h_2 g G_0$	$h_1 g C''$	$C\tau$	
(89, 15, 50)	$h_2 B_5 g$		$C\tau$	
(89, 18, 46)	$\Delta^3 h_0^2 e_0$	$\tau^6 e_0^4 g$	$C\tau, tmf$	$x_{18,50}$
(89, 19, 51)	$\Delta c_0 e_0 g^2$	$\Delta h_1^2 c_0 e_0^2 g + \tau h_0 e_0^4 g + M h_1^4 c_0 d_0 e_0$	$C\tau, h_0$	
(89, 25, 46)	$P^3 \Delta^2 h_0 e_0$	$MP^5 h_2$	$C\tau$	
(89, 31, 47)	$P^5 \Delta c_0 e_0$	$P^5 \Delta h_1^2 c_0 d_0 + \tau P^4 h_0 d_0^4$	$C\tau$	
(89, 40, 46)	$P^9 e_0$	$P^9 h_1^2 d_0$	$C\tau$	

Table 4: E_2 -page generators of the \mathbb{C} -motivic Adams spectral sequence

(s, f, w)	element	d_2	proof	other names
(89, 45, 45)	$P^{11}h_1$			
(90, 10, 48)	$x_{90,10}$			$x_{10,63}$
(90, 12, 48)	M^2			$x_{12,55}$
(90, 15, 49)	$M\Delta h_{1g}$		$C\tau$	rB_4
(90, 17, 47)	$\Delta^3 h_1 e_0$	$\Delta^3 h_1^3 d_0 + \tau^5 e_0^2 g m$	$C\tau, tmf$	$x_{17,52}$
(90, 33, 47)	$P^6 \Delta h_1 e_0$	$P^6 \Delta h_1^3 d_0$	$C\tau$	
(90, 39, 46)	$P^8 j$	$P^9 h_0 e_0$	$C\tau$	
(91, 8, 48)	$x_{91,8}$	$x_{90,10}$	$C\tau$	$x_{8,75}$
(91, 11, 48)	$x_{91,11}$	$M^2 h_0$	$C\tau$	$x_{11,61} + ?h_0^2 h_6 d_0^2$
(91, 17, 49)	$M\Delta c_0 d_0$	$\tau M h_0 d_0^2 e_0$	$C\tau, h_0$	
(91, 17, 50)	$\Delta^2 h_2 g^2$			
(91, 17, 52)	$\Delta h_3 g^3$	$h_0 g m^2$	$C\tau$	
(91, 17, 53)	$g^3 n$			
(91, 45, 46)	$P^{11}h_2$			
(92, 4, 48)	g_3		$C\tau$	
(92, 10, 48)	$x_{92,10}$		$C\tau$	$x_{10,65} + ?h_0^2 h_6 k$
(92, 10, 51)	$\Delta_1 h_1^2 e_1$		$C\tau$	
(92, 12, 48)	$\Delta^2 g_2$		$C\tau$	$x_{12,58} + ?h_0^2 x_{10,65}$
(92, 16, 54)	$d_1 g^3$			
(92, 18, 48)	$\Delta^3 h_2^2 d_0$	$\tau^6 d_0 e_0 g^3$	$C\tau, tmf$	$x_{18,55}$
(92, 24, 48)	$P^2 \Delta^2 d_0^2$	$\tau^2 d_0^3 i j$	$C\tau, tmf$	
(93, 8, 49)	$x_{93,8}$		$C\tau$	$x_{8,78} + h_0 h_6 r$
(93, 9, 51)	$\Delta_1 h_2 e_1$			
(93, 10, 49)	$\Delta h_3 H_1$	$h_1 x_{91,11}$	$C\tau$	$x_{10,67} + h_0^4 h_6 r$
(93, 12, 50)	$\Delta \Delta_1 e_0$	$M^2 h_1^2$	$C\tau$	$x_{12,60} + ?P h_6 c_0 d_0$
(93, 15, 54)	$g^2 i_1$		$C\tau$	
(93, 17, 48)	$\tau \Delta^3 h_{1g}$	$\tau^6 e_0^2 g m$	tmf, h_1	$x_{17,57}$
(94, 8, 49)	$x_{94,8}$	$? \tau h_1^2 x_{91,8}$		$x_{8,80}$
(94, 9, 50)	$x_{94,9}$		$C\tau$	
(94, 10, 50)	$y_{94,10}$		$C\tau$	$x_{10,70} + ?h_0^2 x_{8,80}$
(94, 10, 51)	$x_{94,10}$			
(94, 15, 49)	$\Delta^2 M h_1$			$x_{15,56}$
(94, 16, 49)	$\Delta^3 h_2 c_1$		d_0	$x_{16,54}$
(94, 17, 51)	$M\Delta c_0 e_0$	$M\Delta h_1^2 c_0 d_0 + \tau M h_0 d_0 e_0^2$	$C\tau, h_0$	
(94, 19, 49)	$\Delta^3 c_0 d_0$	$\tau \Delta^2 h_0 d_0^2 e_0$	h_0	
(94, 19, 50)	$\Delta^2 d_0 l$	$\Delta^2 h_0 d_0^2 e_0 + \tau^3 \Delta h_1 e_0^4$	$C\tau, tmf$	$x_{19,49}$
(94, 22, 48)	$\Delta^2 i^2$	$\tau^6 d_0^3 e_0^3$	tmf	$x_{22,39}$
(94, 35, 49)	$P^6 \Delta c_0 d_0$	$\tau P^6 h_0 d_0^2 e_0$	$C\tau$	
(94, 44, 48)	$P^{10} d_0$			
(95, 7, 50)	$x_{95,7}$	$x_{94,9} + ? \tau d_1 H_1$	$C\tau$	$x_{7,79}$
(95, 8, 51)	$x_{95,8}$	$x_{94,10}$	$C\tau$	
(95, 10, 50)	$\Delta x_{71,6}$	$\tau \Delta d_1 e_1$	$C\tau, h_3$	$x_{10,73} + ?h_0^2 h_6 l$

Table 4: E_2 -page generators of the \mathbb{C} -motivic Adams spectral sequence

(s, f, w)	element	d_2	proof	other names
(95, 11, 51)	$x_{95,11}$		$C\tau$	
(95, 15, 54)	g^2B_6			
(95, 16, 52)	$M\Delta h_2^2g$		$C\tau$	
(95, 17, 52)	$\Delta^2h_3g^2$	$P^2h_1^7h_6c_0$	$C\tau$	
(95, 18, 50)	$\Delta^3h_2^2e_0$	$\tau^6e_0^2g^3$	$C\tau, tmf$	$x_{18,57}$
(95, 19, 56)	g^3m	$h_0e_0^2g^3$	$C\tau$	
(95, 21, 48)	$\Delta^3h_0^2i$	$\Delta^2h_0i^2 + \tau^6d_0^3e_0m$	$C\tau, tmf$	$x_{21,43} + ?Px_{17,50}$
(95, 21, 49)	$P\Delta^3h_1d_0$	$\tau^5d_0^3e_0m$	tmf	
(95, 24, 50)	$P^2\Delta^2d_0e_0$	$P^2\Delta^2h_1^2d_0^2 + \tau^2d_0^3ik$	$C\tau, tmf$	
(95, 37, 49)	$P^7\Delta h_1d_0$		$C\tau$	

Table 5: Some permanent cycles in the \mathbb{C} -motivic Adams spectral sequence

(s, f, w)	element	proof
(36, 6, 20)	t	$\langle \tau, \eta^2 \kappa_1, \eta \rangle$
(64, 2, 33)	$h_1 h_6$	$\langle \eta, 2, \theta_5 \rangle$
(68, 7, 36)	$h_3 A'$	$\langle \sigma, \kappa, \tau \eta \theta_{4.5} \rangle$
(69, 3, 36)	$h_2^2 h_6$	$\langle \nu^2, 2, \theta_5 \rangle$
(69, 4, 36)	p'	$\sigma \theta_5$
(70, 5, 36)	$h_0^3 h_3 h_6$	$\langle 8\sigma, 2, \theta_5 \rangle$
(70, 14, 37)	$\tau \Delta^2 h_1^2 g + \tau^3 \Delta h_2^2 g^2$	$\langle \eta, \tau \kappa^2, \tau \bar{\kappa}^2 \rangle$
(71, 4, 37)	$h_6 c_0$	$\langle \epsilon, 2, \theta_5 \rangle$
(72, 6, 37)	$Ph_1 h_6$	$\langle \mu_9, 2, \theta_5 \rangle$
(74, 7, 38)	$Ph_0 h_2 h_6$	Lemma 5.68
(74, 8, 40)	$x_{74,8}$	$\theta_4 \bar{\kappa}_2$
(77, 3, 40)	$h_3^2 h_6$	$\langle \sigma^2, 2, \theta_5 \rangle$
(77, 5, 40)	$h_6 d_0$	$\langle \kappa, 2, \theta_5 \rangle$
(79, 3, 41)	$h_1 h_4 h_6$	$\langle \eta_4, 2, \theta_5 \rangle$
(79, 8, 41)	$Ph_6 c_0$	$\rho_{15} \eta_6$
(80, 6, 42)	$\tau h_1 x_1$	$\langle 2, \eta, \tau \eta \{h_1 x_{76,6}\} \rangle$
(80, 10, 41)	$P^2 h_1 h_6$	$\langle \mu_{17}, 2, \theta_5 \rangle$
(80, 12, 42)	$\Delta^2 d_1$	Lemma 5.19
(82, 4, 43)	$h_6 c_1$	$\langle \bar{\sigma}, 2, \theta_5 \rangle$
(83, 6, 44)	$h_0 h_6 g$	$\langle \kappa \eta_6, \eta, \nu \rangle$
(84, 4, 44)	$h_2^2 h_4 h_6$	$\langle \nu \nu_4, 2, \theta_5 \rangle$
(85, 5, 44)	$\tau h_1 f_2$	Lemma 5.46
(85, 9, 44)	$Ph_6 d_0$	$\langle \tau \eta^2 \bar{\kappa}, 2, \theta_5 \rangle$
(86, 5, 45)	$h_4 h_6 c_0$	$\sigma \{h_1 h_4 h_6\}$
(87, 5, 46)	$h_1^2 c_3$	$\langle \tau \{h_0 Q_3 + h_0 n_1\}, \nu_4, \eta \rangle$
(87, 12, 45)	$P^2 h_6 c_0$	$\rho_{23} \eta_6$
(88, 10, 48)	$x_{88,10}$	Lemma 5.73
(88, 14, 45)	$P^3 h_1 h_6$	$\langle \mu_{25}, 2, \theta_5 \rangle$
(90, 12, 48)	M^2	Lemma 5.28
(91, 7, 49)	$h_1 h_3 h_6 g$	$\langle \{h_1 h_3 g\}, 2, \theta_5 \rangle$
(92, 5, 48)	$h_0 g_3$	Lemma 5.76
(93, 6, 48)	$h_0^3 h_4^2 h_6$	$\langle \theta_5, 2, \theta_4 \rangle$
(93, 10, 50)	$h_1^2 x_{91,8}$	$\langle \kappa_1, \kappa, \tau \eta \theta_{4.5} \rangle$
(94, 6, 49)	$h_6 n$	$\langle \{n\}, 2, \theta_5 \rangle$
(95, 5, 50)	$h_6 d_1$	$\langle \kappa_1, 2, \theta_5 \rangle$
(95, 7, 49)	$\Delta h_1 h_3 h_6$	$\langle \{\Delta h_1 h_3\}, 2, \theta_5 \rangle$
(95, 16, 49)	$P^3 h_6 c_0$	Lemma 5.32

Table 6: Adams d_3 differentials in the \mathbb{C} -motivic Adams spectral sequence

(s, f, w)	element	d_3	proof
(15, 2, 8)	h_0h_4	h_0d_0	$C\tau$
(30, 6, 16)	Δh_2^2	$\tau h_1d_0^2$	tmf
(31, 4, 16)	$h_0^3h_5$	$h_0 \cdot \Delta h_2^2$	$C\tau$
(31, 8, 17)	τd_0e_0	Pc_0d_0	$d_4(\tau^2d_0e_0 + h_0^7h_5)$
(34, 2, 18)	h_2h_5	τh_1d_1	Lemma 5.12
(37, 8, 21)	τe_0g	$c_0d_0^2$	$d_4(\tau^2e_0g)$
(38, 4, 21)	e_1	h_1t	$C\tau$
(39, 12, 21)	τPd_0e_0	$P^2c_0d_0$	h_1
(40, 4, 22)	f_1	0	h_2
(46, 14, 24)	i^2	$\tau P^2h_1d_0^2$	tmf
(47, 16, 25)	$\tau P^2d_0e_0$	$P^3c_0d_0$	h_1
(47, 18, 24)	h_0^5Q'	$P^4h_0d_0$	$C\tau$
(49, 6, 27)	$h_1h_5e_0$	Mh_1^3	$C\tau$
(49, 11, 26)	τ^2d_0m	$P\Delta h_1^2d_0$	mmf
(50, 10, 28)	Δh_2^2g	$\tau h_1d_0e_0^2$	τ
(54, 8, 28)	h_5i	MPh_0	$C\tau$
(54, 6, 28)	$\tau^2\Delta_1h_1^2$	τMc_0	Lemma 5.13
(55, 11, 30)	τ^2gm	$\Delta h_1^2d_0^2$	$d_4(\tau^3gm)$
(55, 20, 29)	$\tau P^3d_0e_0$	$P^4c_0d_0$	h_1
(55, 7, 30)	B_6	τh_2gn	$C\tau$
(56, 8, 31)	$h_5c_0e_0$	$Mh_1^2c_0$	$C\tau$
(56, 9, 30)	Ph_5e_0	MPh_1^2	$C\tau$
(56, 10, 32)	gt	0	τ
(56, 13, 30)	$\tau\Delta h_1d_0e_0$	$P\Delta h_1c_0d_0$	$d_4(\tau^2\Delta h_1d_0e_0)$
(57, 12, 33)	τe_0g^2	$c_0d_0e_0^2$	$C\tau$
(57, 8, 30)	h_5j	MPh_2	$C\tau$
(57, 15, 30)	τ^2Pd_0m	$P^2\Delta h_1^2d_0$	mmf
(57, 7, 30)	Q_2	τ^2gt	$\tau\Delta h_1g$
(58, 8, 33)	e_1g	h_1gt	$C\tau$
(60, 12, 35)	τg^3	$Mh_1^6c_0$	$C\tau$
(61, 4, 32)	D_3	Mh_4	$C\tau$
(62, 13, 34)	$\tau\Delta h_1e_0g$	$\Delta h_1c_0d_0^2$	$C\tau$
(62, 8, 33)	Δe_1	Δh_2^2n	$C\tau$
(62, 8, 33)	C_0	Δh_2^2n	$C\tau$
(62, 10, 33)	$h_1 \cdot \Delta x + \tau Me_0$	MPc_0	$C\tau$
(62, 10, 32)	$\tau h_1 \cdot \Delta x$	0	Δh_2^2
(62, 10, 32)	$\Delta^2h_3^2$	0	Δh_2^2
(62, 22, 32)	P^2i^2	$\tau P^4h_1d_0^2$	tmf
(63, 8, 32)	$h_0^7h_6$	$\Delta^2h_0h_3^2 + ?\tau^2Mh_0e_0$	$C\tau$
(63, 24, 33)	$\tau P^4d_0e_0$	$P^5c_0d_0$	h_1
(64, 17, 34)	$\tau P\Delta h_1d_0e_0$	$P^2\Delta h_1c_0d_0$	$d_4(\tau^2P\Delta h_1d_0e_0)$
(65, 13, 36)	Δh_2^2m	$MPh_1^3c_0$	$C\tau$
(65, 19, 34)	$\tau^2P^2d_0m$	$P^3\Delta h_1^2d_0$	mmf

Table 6: Adams d_3 differentials in the \mathbb{C} -motivic Adams spectral sequence

(s, f, w)	element	d_3	proof
(67, 5, 35)	$\tau Q_3 + \tau n_1$	0	h_2
(67, 9, 37)	C''	nm	$C\tau$
(67, 9, 36)	X_3	0	τg
(68, 4, 36)	d_2	$h_0^2 Q_3$	$C\tau$
(68, 11, 35)	$\tau h_0^3 \cdot \Delta g_2$	$\tau^3 \Delta h_2^2 e_0 g$	Lemma 5.14
(68, 11, 38)	$Mh_2 g$	0	τ
(69, 8, 36)	$\tau D_3' + ?\tau^2 h_2 G_0$	$\tau^2 Mh_2 g$	Lemma 5.15
(69, 11, 38)	$h_2 B_5$	$Mh_1 c_0 d_0$	$C\tau$
(69, 13, 36)	$\tau \Delta^2 h_1 g$	$\tau^4 e_0^4$	tmf
(70, 2, 36)	$h_3 h_6$	$h_0 p'$	$C\tau$
(70, 4, 37)	p_1	$\tau h_1^2 Q_3$	$C\tau$
(70, 14, 37)	$\tau M P e_0$	$M P^2 c_0$	$d_4(\tau^2 M P e_0)$
(70, 14, 40)	m^2	$\tau h_1 e_0^4$	τ
(71, 28, 37)	$\tau P^5 d_0 e_0$	$P^6 c_0 d_0$	h_1
(72, 21, 39)	$\tau P^2 \Delta h_1 d_0 e_0$	$P^3 \Delta h_1 c_0 d_0$	$d_4(\tau^2 P^2 \Delta h_1 d_0 e_0)$
(73, 23, 38)	$\tau^2 P^3 d_0 m$	$P^4 \Delta h_1^2 d_0$	mmf
(74, 6, 38)	$Ph_2 h_6$	$\tau h_1 h_4 Q_2$	Lemma 5.12
(75, 7, 40)	$x_{75,7}$	$h_0^2 x_{74,8}$	$C\tau$
(75, 11, 42)	$g B_6$	$\tau h_2 g^2 n$	h_2
(75, 15, 42)	$\tau^2 g^2 m$	$\Delta h_1^2 d_0 e_0^2$	$d_4(\tau^3 g^2 m)$
(76, 5, 40)	$h_4 D_3$	$d_0 D_3$	$C\tau$
(76, 14, 41)	$\Delta^2 h_1 h_3 g$	0	h_1
(76, 14, 44)	$g^2 t$	0	τ
(77, 14, 40)	$\tau^2 M h_0 l$	$\Delta^2 h_0 d_0^2$	Lemma 5.16
(77, 16, 45)	$\tau e_0 g^3$	$c_0 e_0^4$	$d_4(\tau^2 e_0 g^3)$
(77, 17, 41)	$\Delta^2 h_1 d_0^2$	$\tau^3 d_0^3 e_0^2$	tmf
(78, 3, 40)	$h_0 h_4 h_6$	$h_0 h_6 d_0$	$C\tau$
(78, 12, 45)	$e_1 g^2$	$h_1 g^2 t$	$C\tau$
(78, 13, 40)	$h_0^3 x_{78,10}$	$\tau^6 e_0 g^3$	Lemma 5.17
(78, 18, 41)	$\tau M P^2 e_0$	$M P^3 c_0$	h_1
(78, 30, 40)	$P^4 i^2$	$\tau P^6 h_1 d_0^2$	tmf
(79, 5, 42)	x_1	$\tau h_1 m_1$	Lemma 5.18
(79, 32, 41)	$\tau P^6 d_0 e_0$	$P^7 c_0 d_0$	h_1
(79, 34, 40)	$P^4 h_0^5 Q'$	$P^8 h_0 d_0$	$C\tau$
(80, 14, 41)	$\Delta^3 h_1 h_3$	$\tau^4 \Delta h_1 e_0^2 g + ?\tau \Delta^2 h_0 d_0 e_0$	Lemma 5.20
(80, 14, 42)	$\tau^2 d_0 B_5$	$\Delta^2 h_0 d_0 e_0$	Lemma 5.21
(80, 25, 42)	$\tau P^3 \Delta h_1 d_0 e_0$	$P^4 \Delta h_1 c_0 d_0$	$d_4(\tau^2 P^3 \Delta h_1 d_0 e_0)$
(81, 3, 42)	$h_2 h_4 h_6$	0	Lemma 5.22
(81, 12, 42)	$\Delta^2 p$	0	Lemma 5.23
(81, 27, 42)	$\tau^2 P^4 d_0 m$	$P^5 \Delta h_1^2 d_0$	mmf
(82, 6, 44)	$h_4 Q_3$	$h_3 x_{74,8}$	$C\tau$
(82, 10, 42)	$P^2 h_2 h_6$	0	Lemma 5.22
(82, 12, 45)	$g C_0$	$\Delta h_2^2 g n$	$C\tau$
(82, 14, 45)	$\tau M e_0 g$	$M c_0 d_0^2$	h_1

Table 6: Adams d_3 differentials in the \mathbb{C} -motivic Adams spectral sequence

(s, f, w)	element	d_3	proof
(82, 17, 46)	$\tau\Delta h_1 e_0 g^2$	$\Delta h_1 c_0 d_0 e_0^2$	$d_4(\tau^2 \Delta h_1 e_0 g^2)$
(83, 5, 43)	$\tau h_6 g + \tau h_2 e_2$	0	Lemma 5.24
(83, 17, 45)	$\Delta^2 h_1 e_0^2$	$\tau^3 d_0 e_0^4$	τ
(84, 4, 44)	f_2	$\tau h_1 h_4 Q_3$	$C\tau, h_1^2$
(84, 19, 45)	$\Delta^2 c_0 d_0^2$	$\tau \Delta h_0^2 d_0^3 e_0$	mmf
(85, 6, 45)	$x_{85,6}$	0	Lemma 5.25
(85, 21, 45)	$P\Delta^2 h_1 d_0^2$	$\tau^3 d_0^6$	tmf
(86, 4, 45)	$h_1 c_3$	0	h_0
(86, 6, 46)	$h_2 h_6 g$	$?\tau h_1^3 h_4 Q_3$	
(86, 11, 45)	$\tau^3 g G_0$	0	d_0
(86, 12, 45)	$\Delta^2 e_1$	$\Delta^2 h_1 t$	$C\tau$
(86, 17, 46)	$\tau\Delta^2 h_1 e_0 g$	$\Delta^2 h_1 c_0 d_0^2 + \tau^4 e_0^5$	mmf
(86, 22, 45)	$\tau M P^3 e_0$	$M P^4 c_0$	h_1
(87, 7, 45)	$x_{87,7}$	0	d_0
(87, 13, 49)	$g C''$	$g n m$	$C\tau$
(87, 36, 45)	$\tau P^7 d_0 e_0$	$P^8 c_0 d_0$	h_1
(88, 12, 46)	$\Delta^2 f_1$	0	h_2
(88, 18, 46)	$\tau^2 M h_0 d_0 k$	$P\Delta^2 h_0 d_0 e_0$	Lemma 5.26
(88, 18, 46)	$\Delta^3 h_1^2 d_0$	$\tau^3 \Delta h_1 d_0^2 e_0^2 + ?P\Delta^2 h_0 d_0 e_0$	mmf
(88, 29, 46)	$\tau P^4 \Delta h_1 d_0 e_0$	$P^5 \Delta h_1 c_0 d_0$	$d_4(\tau^2 P^4 \Delta h_1 d_0 e_0)$
(89, 14, 51)	$h_1^9 h_6 e_0$	$M h_1^3 g^2$	$C\tau$
(89, 15, 50)	$h_2 B_5 g$	$M h_1 c_0 e_0^2$	Lemma 5.27
(89, 17, 48)	$\tau\Delta^2 h_1 g^2$	$\tau^4 e_0^4 g$	τ
(89, 31, 46)	$\tau^2 P^5 d_0 m$	$P^6 \Delta h_1^2 d_0$	mmf
(90, 18, 52)	$g m^2$	$\tau h_1 e_0^4 g$	τ
(90, 19, 49)	$\Delta^2 c_0 e_0^2$	$\tau \Delta h_2^2 d_0^3 e_0$	mmf
(92, 10, 48)	$x_{92,10}$	0	d_0
(92, 14, 50)	$m Q_2$	$\tau^3 g^3 n$	$C\tau, \tau$
(92, 23, 49)	$P\Delta^2 c_0 d_0^2$	$\tau P^2 \Delta h_2^2 d_0^2 e_0$	mmf
(93, 7, 48)	$\Delta h_2^2 h_6$	$\tau h_1 h_6 d_0^2$	Lemma 5.29
(93, 8, 49)	$x_{93,8}$	$\Delta h_2^2 H_1$	$C\tau$
(93, 13, 48)	$P^2 h_6 d_0$	0	Lemma 5.30
(93, 22, 48)	$\tau^2 M P h_0 d_0 j$	$P^2 \Delta^2 h_0 d_0^2$	Lemma 5.31
(93, 25, 49)	$P^2 \Delta^2 h_1 d_0^2$	$\tau^3 P d_0^6$	tmf
(94, 8, 49)	$?x_{94,8}$	$?h_1 x_{92,10}$	
(94, 9, 49)	$\tau h_6 d_0 e_0$	$P h_6 c_0 d_0$	h_1
(94, 15, 49)	$\Delta^2 M h_1$	$?\tau^3 M d_0 e_0^2$	
(94, 15, 52)	$P h_1^6 h_6 e_0$	$M \Delta h_1^4 g$	$C\tau$
(94, 17, 50)	$\tau^2 M d_0 m$	$M P \Delta h_1^2 d_0$	d_0
(94, 21, 50)	$\tau\Delta^2 h_1 d_0^2 e_0$	$P\Delta^2 h_1 c_0 d_0^2 + \tau^4 d_0^3 e_0^3$	mmf
(94, 26, 49)	$\tau M P^4 e_0 +$ $+ \tau P^2 \Delta^2 h_1^2 d_0^2$	$M P^5 c_0$	h_1
(94, 38, 48)	$P^6 i^2$	$\tau P^8 h_1 d_0^2$	tmf
(95, 15, 54)	$g^2 B_6$	$\tau h_2 g^3 n$	$C\tau$

Table 6: Adams d_3 differentials in the \mathbb{C} -motivic Adams spectral sequence

(s, f, w)	element	d_3	proof
(95, 16, 52)	$M\Delta h_2^2 g$	$\tau M h_1 d_0 e_0^2$	τ
(95, 19, 54)	$\tau^2 g^3 m$	$\Delta h_1^2 e_0^4 + M h_1^4 d_0^2 e_0$	mmf
(95, 20, 50)	$\Delta^3 h_1 c_0 d_0$	$\tau^4 d_0^3 e_0 m$	mmf
(95, 36, 48)	$h_0^{15} \cdot \Delta^3 h_0^2 i$	$P^6 h_0 i^2$	$C\tau$
(95, 40, 49)	$\tau P^8 d_0 e_0$	$P^9 c_0 d_0$	h_1

Table 7: Adams d_4 differentials in the \mathbb{C} -motivic Adams spectral sequence

(s, f, w)	element	d_4	proof
(31, 8, 16)	$\tau^2 d_0 e_0 + h_0^7 h_5$	$P^2 d_0$	$C\tau$
(37, 8, 20)	$\tau^2 e_0 g$	$P d_0^2$	tmf
(38, 2, 20)	$h_3 h_5$	$h_0 x$	$C\tau$
(39, 12, 20)	$\tau^2 P d_0 e_0$	$P^3 d_0$	tmf
(42, 6, 22)	$P h_2 h_5$	0	d_0
(47, 16, 24)	$\tau^2 P^2 d_0 e_0$	$P^4 d_0$	tmf
(50, 6, 27)	C	$P h_1^2 h_5 c_0$	$C\tau$
(50, 10, 26)	$\tau^2 \Delta h_2^2 g$	ij	tmf
(55, 11, 29)	$\tau^3 g m$	$P \Delta c_0 d_0 + \tau d_0^2 j$	tmf
(55, 20, 28)	$\tau^2 P^3 d_0 e_0$	$P^5 d_0$	tmf
(56, 13, 29)	$\tau^2 \Delta h_1 d_0 e_0$	$P^2 \Delta h_1 d_0$	tmf
(57, 12, 32)	$\tau^2 e_0 g^2$	d_0^4	d_0
(58, 14, 30)	$\tau^2 \Delta h_2^2 d_0^2$	$P ij$	tmf
(62, 10, 32)	$\tau h_1 \cdot \Delta x$	$\tau^2 \Delta h_2^2 d_0 e_0$	Lemma 5.35
(62, 10, 32)	$\Delta^2 h_3^2$	0	Lemma 5.36
(62, 13, 33)	$\tau^2 \Delta h_1 e_0 g$	$P \Delta h_1 d_0^2$	tmf
(63, 7, 34)	C'	$M h_2 d_0$	$C\tau$
(63, 7, 33)	τX_2	$\tau M h_2 d_0$	Lemma 5.37
(63, 11, 33)	$\tau^2 M h_1 e_0$	$M P^2 h_1$	d_0
(63, 15, 33)	$\tau^3 d_0^2 m$	$P^2 \Delta c_0 d_0 + \tau P d_0^2 j$	tmf
(63, 19, 32)	$h_0^{18} h_6$	$P^2 h_0 i^2$	$C\tau$
(64, 17, 33)	$\tau^2 \Delta h_2^2 d_0 g$	$P^3 \Delta h_1 d_0$	tmf
(66, 18, 34)	$\tau^2 P \Delta h_2^2 d_0^2$	$P^2 ij$	tmf
(68, 5, 36)	$h_0 d_2$	X_3	Lemma 5.38
(69, 11, 37)	$\tau h_2 B_5$	$M P h_1 d_0$	d_0
(70, 10, 39)	$h_2 C''$	$h_1^4 c_0 Q_2$	$C\tau$
(70, 14, 38)	$\tau^2 \Delta h_2^2 g^2$	$\Delta h_0^2 d_0^2 e_0$	τ
(70, 14, 36)	$\tau^2 M P e_0$	$M P^3$	d_0
(71, 19, 37)	$\tau^3 P d_0^2 m$	$P^3 \Delta c_0 d_0 + \tau P^2 d_0^2 j$	tmf
(71, 28, 36)	$\tau^2 P^5 d_0 e_0$	$P^7 d_0$	tmf
(72, 9, 40)	$h_2^2 G_0$	$\tau g^2 n$	Lemma 5.40
(72, 21, 37)	$\tau^2 P^2 \Delta h_1 d_0 e_0$	$P^4 \Delta h_1 d_0$	tmf

Table 7: Adams d_4 differentials in the \mathbb{C} -motivic Adams spectral sequence

(s, f, w)	element	d_4	proof
(74, 22, 38)	$\tau^2 P^2 \Delta h_2^2 d_0^2$	$P^3 i j$	<i>tmf</i>
(75, 5, 40)	$h_3 d_2$	$h_0 x_{74,8}$	$C\tau$
(75, 11, 40)	$\Delta h_0^2 h_3 g_2$	$\tau M h_1 d_0^2$	Lemma 5.41
(75, 15, 41)	$\tau^3 g^2 m$	$\Delta c_0 d_0^3 + \tau d_0^3 l$	<i>tmf</i>
(76, 14, 40)	$\tau^2 M d_0 e_0$	$M P^2 d_0$	h_1
(76, 14, 41)	$\Delta^2 h_1 h_3 g$	$\tau \Delta h_2^2 d_0^2 e_0$	Lemma 5.42
(77, 16, 44)	$\tau^2 e_0 g^3$	$d_0^3 e_0^2$	τ
(78, 5, 40)	$h_0^3 h_4 h_6$	$h_0^2 x_{77,7}$	$C\tau$
(79, 23, 41)	$\tau^3 P^2 d_0^2 m$	$P^4 \Delta c_0 d_0 + \tau P^3 d_0^2 j$	<i>tmf</i>
(79, 32, 40)	$\tau^2 P^6 d_0 e_0$	$P^8 d_0$	<i>tmf</i>
(80, 5, 42)	$h_0 e_2$	$\tau h_1^3 x_{76,6}$	Lemma 5.43
(80, 25, 41)	$\tau^2 P^3 \Delta h_1 d_0 e_0$	$P^5 \Delta h_1 d_0$	<i>tmf</i>
(81, 8, 43)	$\tau g D_3$	0	h_1
(81, 15, 42)	$\Delta^3 h_1^2 h_3$	$\tau^4 d_0 e_0^2 l$	Lemma 5.44
(82, 14, 44)	$\tau^2 M e_0 g$	$M P d_0^2$	h_1
(82, 17, 45)	$\tau^2 \Delta h_1 e_0 g^2$	$\Delta h_1 d_0^4$	τ
(82, 26, 42)	$\tau^2 P^3 \Delta h_2^2 d_0^2$	$P^4 i j$	<i>tmf</i>
(83, 11, 46)	$g C'$	$M h_0 e_0 g$	$C\tau$
(83, 11, 45)	Δj_1	$\tau M h_0 e_0 g$	Lemma 5.45
(86, 4, 45)	$h_1 c_3$	$\tau h_0 h_2 h_4 Q_3$	Lemma 5.48
(86, 10, 44)	$h_0^2 h_6 i$	0	$d_0, C\tau$
(86, 22, 44)	$\tau^2 M P^3 e_0$	$M P^5$	h_1
(87, 7, 45)	$x_{87,7}$	0	Lemma 5.49
(87, 10, 45)	$\tau \Delta h_1 H_1$	0	Lemma 5.50
(87, 15, 47)	$\Delta^2 c_1 g$	0	τ
(87, 27, 45)	$\tau^3 P^3 d_0^2 m$	$P^5 \Delta c_0 d_0 + \tau P^4 d_0^2 j$	<i>tmf</i>
(87, 36, 44)	$\tau^2 P^7 d_0 e_0$	$P^9 d_0$	<i>tmf</i>
(88, 17, 48)	$\Delta^2 h_0 g^2$	$\tau \Delta h_1 d_0^2 e_0^2$	<i>mmf</i>
(88, 29, 45)	$\tau^2 P^4 \Delta h_1 d_0 e_0$	$P^6 \Delta h_1 d_0$	<i>tmf</i>
(89, 15, 49)	$\tau h_2 B_5 g$	$M h_1 d_0^3$	Lemma 5.51
(90, 14, 51)	$h_2 g C''$	$P h_1^{10} h_6 c_0$	$C\tau$
(90, 18, 50)	$\tau^2 g m^2$	$\Delta h_2^2 d_0^3 e_0$	<i>mmf</i>
(90, 30, 46)	$\tau^2 \Delta P^4 h_2^2 d_0^2$	$P^5 i j$	<i>tmf</i>
(91, 12, 48)	$\Delta h_2^2 A'$	0	Lemma 5.52
(91, 20, 50)	$\Delta^2 h_1 c_0 e_0^2$	$\tau^2 d_0^4 e_0^2$	<i>mmf</i>
(92, 13, 52)	$h_2^2 g G_0$	$\tau g^3 n$	τ
(93, 3, 48)	$h_4^2 h_6$	$h_0^3 g_3$	Lemma 5.53
(95, 16, 50)	$\Delta^2 M h_1^2$	$? M P \Delta h_0^2 e_0$	
(95, 16, 50)	$\tau^2 M \Delta h_2^2 g$	$M P \Delta h_0^2 e_0$	d_0
(95, 19, 53)	$\tau^3 g^3 m$	$\Delta c_0 d_0^2 e_0^2 + \tau d_0^3 e_0 m$	τ
(95, 31, 49)	$\tau^3 P^4 d_0^2 m$	$P^6 \Delta c_0 d_0 + \tau P^5 d_0^2 j$	<i>tmf</i>
(95, 40, 48)	$P^6 \Delta h_0^5 i + \tau^2 P^8 d_0 e_0$	$P^{10} d_0$	<i>tmf</i>

Table 8: Adams d_5 differentials in the \mathbb{C} -motivic Adams spectral sequence

(s, f, w)	element	d_5	proof
(56, 9, 29)	$\tau P h_5 e_0$	$\tau \Delta h_0^2 d_0 e_0$	[16, Lemma 3.92]
(61, 6, 32)	A'	$\tau M h_1 d_0$	[44, Theorem 12.1]
(63, 11, 33)	$\tau h_1^2 \cdot \Delta x$	$\tau^3 d_0^2 e_0^2$	Lemma 5.55
(63, 23, 32)	$h_0^{22} h_6$	$P^6 d_0$	$C\tau$
(67, 6, 36)	$h_0 Q_3 + h_2^2 D_3$	0	$C\tau, \tau$
(68, 12, 36)	$h_5 d_0 i$	$\tau \Delta h_1 d_0^3$	Lemma 5.56
(70, 4, 36)	$\tau p_1 + h_0^2 h_3 h_6$	$\tau^2 h_2^2 C'$	Lemma 5.57, [5]
(72, 7, 39)	$h_1 x_{71,6}$	0	Lemma 5.58
(73, 7, 38)	$h_4 D_2$	$\tau^4 d_1 g^2$	Lemma 5.59
(81, 10, 44)	$g A'$	$\tau M h_1 e_0^2$	τ
(85, 6, 45)	$x_{85,6}$	0	h_2
(86, 10, 44)	$h_0^2 h_6 i$	$\Delta^2 h_0^2 x$	$C\tau$
(86, 11, 45)	$\tau^3 g G_0$	$\tau M \Delta h_1^2 d_0$	Lemma 5.60
(92, 4, 48)	g_3	$h_6 d_0^2$	Lemma 5.61
(92, 12, 48)	$\Delta^2 g_2$	$? \tau^2 \Delta^2 h_2 g^2$	
(93, 8, 48)	$h_0 \cdot \Delta h_2^2 h_6$	$\Delta^2 h_0 g_2$	$C\tau$
(93, 13, 50)	$e_0 x_{76,9}$	$M \Delta h_1 c_0 d_0$	Lemma 5.62

Table 9: Higher Adams differentials in the \mathbb{C} -motivic Adams spectral sequence

(s, f, w)	element	r	d_r	proof
(67, 5, 35)	$\tau Q_3 + \tau n_1$	6	0	Lemma 5.65
(68, 7, 37)	$h_2^2 H_1$	6	$M c_0 d_0$	Lemma 5.66
(68, 7, 36)	$\tau h_2^2 H_1$	7	$M P d_0$	Lemma 5.66
(77, 7, 42)	m_1	7	0	Lemma 5.69
(80, 6, 43)	$h_1 x_1$	8	0	Lemma 5.70
(81, 3, 42)	$h_2 h_4 h_6$	6	0	Lemma 5.71
(83, 5, 43)	$\tau h_6 g + \tau h_2 e_2$	9	$? \Delta^2 h_2 n$	
(85, 6, 45)	$x_{85,6}$	9	$? M \Delta h_1 d_0$	
(86, 6, 46)	$h_2 h_6 g + ? h_1^2 f_2$	10	$? M \Delta h_1^2 d_0$	
(87, 7, 45)	$x_{87,7}$	9	$? \tau M \Delta h_0^2 e_0$	
(87, 9, 48)	$g Q_3$	6	0	Lemma 5.65
(87, 10, 45)	$\tau \Delta h_1 H_1$	6	$? \tau M \Delta h_0^2 e_0$	
(88, 11, 49)	$h_2^2 g H_1$	6	$M c_0 e_0^2$	Lemma 5.74
(88, 11, 48)	$\tau h_2^2 g H_1$	7	0	Lemma 5.74
(88, 12, 46)	$\Delta^2 f_1$	6	$? \tau^2 M d_0^3$	
(88, 12, 48)	$\Delta g_2 g$	6	$M d_0^3$	Lemma 5.75
(92, 10, 51)	$\Delta_1 h_1^2 e_1$	6	0	Lemma 5.77
(93, 9, 51)	$\Delta_1 h_2 e_1$	8	0	Lemma 5.79
(93, 13, 49)	$\tau e_0 x_{76,9}$	6	$M P \Delta h_1 d_0$	[5]

Table 10: Unknown Adams differentials in the \mathbb{C} -motivic Adams spectral sequence

(s, f, w)	element	r	possible d_r
(63, 8, 32)	$h_0^7 h_6$	3	$\Delta^2 h_0 h_3^2 + ?\tau^2 M h_0 e_0$
(69, 8, 37)	D_3'	2	$h_1 X_3 + ?\tau h_1 C''$
(80, 14, 41)	$\Delta^3 h_1 h_3$	3	$\tau^4 \Delta h_1 e_0^2 g + ?\tau \Delta^2 h_0 d_0 e_0$
(83, 5, 43)	$\tau h_6 g + \tau h_2 e_2$	9	$? \Delta^2 h_2 n$
(85, 6, 45)	$x_{85,6}$	9	$? M \Delta h_1 d_0$
(86, 6, 46)	$h_2 h_6 g$	3	$? \tau h_1^3 h_4 Q_3$
(86, 6, 46)	$h_2 h_6 g + ? h_1^2 f_2$	10	$? M \Delta h_1^2 d_0$
(87, 7, 45)	$x_{87,7}$	9	$? \tau M \Delta h_0^2 e_0$
(87, 10, 45)	$\tau \Delta h_1 H_1$	6	$? \tau M \Delta h_0^2 e_0$
(88, 12, 46)	$\Delta^2 f_1$	6	$? \tau^2 M d_0^3$
(88, 18, 46)	$\Delta^3 h_1^2 d_0$	3	$\tau^3 \Delta h_1 d_0^2 e_0^2 + ? P \Delta^2 h_0 d_0 e_0$
(92, 12, 48)	$\Delta^2 g_2$	5	$? \tau^2 \Delta^2 h_2 g^2$
(94, 8, 49)	$x_{94,8}$	2	$? \tau h_1^2 x_{91,8}$
(94, 8, 49)	$x_{94,8}$	3	$? h_1 x_{92,10}$
(94, 15, 49)	$\Delta^2 M h_1$	3	$? \tau^3 M d_0 e_0^2$
(95, 7, 50)	$x_{95,7}$	2	$x_{94,9} + ? \tau d_1 H_1$
(95, 16, 50)	$\Delta^2 M h_1^2$	4	$? M P \Delta h_0^2 e_0$

Table 11: Some Toda brackets

(s, w)	bracket	contains	indeterminacy	proof	used in
(2, 1)	$\langle 2, \eta, 2 \rangle$	τh_1^2	0	$\langle h_0, h_1, h_0 \rangle$	6.22, 7.20, 7.21, 7.27, 7.38, 7.39 7.41, 7.44, 7.51, 7.52, 7.85
(6, 4)	$\langle \eta, \nu, \eta \rangle$	h_2^2	0	$\langle h_1, h_2, h_1 \rangle$	7.91, 7.98
(8, 5)	$\langle \nu, \eta, \nu \rangle$	$h_1 h_3$	0	$\langle h_2, h_1, h_2 \rangle$	7.116
(8, 5)	$\langle \eta, \nu, 2\nu \rangle$	c_0	$h_1 h_3$	$\langle h_1, h_2, h_0 h_2 \rangle$	7.80, 7.154
(8, 5)	$\langle \eta^2, \nu, \eta, 2 \rangle$	c_0	0	$\langle h_1^2, h_2, h_1, h_0 \rangle$	6.5
(9, 5)	$\langle 2, \epsilon, 2 \rangle$	$\tau h_1 c_0$	0	Corollary 6.2	7.28
(9, 5)	$\langle \eta, 2, 8\sigma \rangle$	Ph_1	$\tau h_1^2 h_3, \tau h_1 c_0$	$\langle h_1, h_0, h_0^3 h_3 \rangle$	7.31, 7.74
(10, 5)	$\langle 2, \mu_9, 2 \rangle$	τPh_1^2	0	Corollary 6.2	7.31
(11, 6)	$\langle 2, \eta, \tau \eta \epsilon \rangle$	Ph_2	$Ph_0 h_2, \tau Ph_1^3$	$\langle h_0, h_1, \tau h_1 c_0 \rangle$	7.78
(15, 8)	$\langle 8, 2\sigma, \sigma \rangle$	$h_0^3 h_4$	$h_0^6 h_4, h_0^7 h_4$	$d_2(h_4) = h_0 h_3^2$	7.150
(15, 8)	$\langle 2, \kappa, 2 \rangle$	$\tau h_1 d_0$	$h_0^4 h_4, h_0^5 h_4$ $h_0^6 h_4, h_0^7 h_4$	Corollary 6.2	7.37, 7.156
(16, 9)	$\langle \eta, \sigma^2, 2 \rangle$	$h_1 h_4$	Pc_0	$d_2(h_4) = h_0 h_3^2$	6.19
(16, 9)	$\langle \eta, 2, \sigma^2 \rangle$	$h_1 h_4$	Pc_0	$d_2(h_4) = h_0 h_3^2$	6.6
(17, 9)	$\langle 2, \eta_4, 2 \rangle$	$\tau h_1^2 h_4$	0	Corollary 6.2	7.39
(17, 10)	$\langle 2, \eta, \eta \kappa \rangle$	$h_2 d_0$	0	$d_2(e_0) = h_1^2 d_0$	7.84
(18, 10)	$\langle \nu, \sigma, 2\sigma \rangle$	$h_2 h_4$	0	$d_2(h_4) = h_0 h_3^2$	5.48
(20, 11)	$\langle \kappa, 2, \eta, \nu \rangle$	τg	$\tau h_0^2 g$	Lemma 6.5	7.108
(20, 12)	$\langle \nu, \eta, \eta \kappa \rangle$	$h_0 g$	$h_0^2 g$	$d_2(e_0) = h_1^2 d_0$	7.106
(21, 12)	$\langle \nu, 2\nu, \kappa \rangle$	$\tau h_1 g$	$h_2^2 h_4$	$d_2(f_0) = h_0^2 e_0$	7.147
(23, 12)	$\langle \sigma, 16, 2\rho_{15} \rangle$	$h_0^2 i + \tau Ph_1 d_0$	$\tau h_4 c_0$	$\langle h_3, h_0^4, h_0^4 h_4 \rangle$	6.16
(30, 16)	$\langle \sigma^2, 2, \sigma^2, 2 \rangle$	h_4^2	0	$d_2(h_4) = h_0 h_3^2$	7.36
(32, 17)	$\langle \eta, 2, \theta_4 \rangle$	$h_1 h_5$	0	$d_2(h_5) = h_0 h_4^2$	6.14
(32, 18)	$\langle \eta, \sigma^2, \eta, \sigma^2 \rangle$	d_1	0	$\langle h_1, h_3^2, h_1, h_3^2 \rangle$	7.81
(33, 18)	$\langle \eta \theta_4, \eta, 2 \rangle$	p	0	$\langle h_1 h_4^2, h_1, h_0 \rangle$	7.22

Table 11: Some Toda brackets

(s, w)	bracket	detected by	indeterminacy	proof	used in
(36, 20)	$\langle \tau, \eta^2 \kappa_1, \eta \rangle$	t	Ph_1^4	Lemma 6.9	Table 5
(36, 20)	$\langle \nu, \eta, \eta \theta_4 \rangle$	t	0	[6, ?]	7.22
(39, 21)	$\langle \eta_5, \nu, 2\nu \rangle$	$h_5 c_0$	$h_1 h_3 h_5$	$\langle h_1 h_5, h_2, h_0 h_2 \rangle =$ $= h_5 \langle h_1, h_2, h_0 h_2 \rangle$	7.147
(39, 21)	$\langle \epsilon, 2, \theta_4 \rangle$	$h_5 c_0$	0	$d_2(h_5) = h_0 h_4^2$	7.147
(45, 24)	$\langle 2, \theta_4, \kappa \rangle$	$h_5 d_0$	$h_0 h_3^2 h_5, h_0 h_5 d_0$ $h_0^2 h_5 d_0$	$d_2(h_5) = h_0 h_4^2$	7.147, 7.148
(57, 30)	$\langle \tau, \tau \eta \kappa \bar{\kappa}^2, \eta \rangle$	$h_0 h_2 h_5 i$	$P^6 h_1 c_0$	$d_5(\tau P h_5 e_0) = \tau \Delta h_0^2 d_0 e_0$	5.42
(62, 32)	$\langle 2, \theta_4, \theta_4, 2 \rangle$	h_5^2	$h_5 n$	$d_2(h_5) = h_0 h_4^2$	5.76, 7.156
(62, 33)	$\langle \eta, \eta \kappa, \tau \theta_{4.5} \rangle$	$\Delta e_1 + C_0$	0	Remark 7.107	7.106
(63, 33)	$\langle \tau \eta \theta_{4.5}, \kappa, 2, \eta \rangle$	$\tau h_1 H_1$	0	Remark 7.109	7.108
(64, 33)	$\langle \eta, 2, \theta_5 \rangle$	$h_1 h_6$	$\tau h_1^2 h_5^2, \tau^2 h_1 X_2$ $\tau h_3 Q_2, P^7 c_0$	$d_2(h_6) = h_0 h_5^2$	5.61, 7.21, 7.27, 7.41, Table 5 7.44, 7.52, 7.67, 7.78, 7.84 7.110, 7.151
(64, 34)	$\langle \nu, \eta, \tau \kappa \theta_{4.5} \rangle$	$h_2 A'$	0	$d_5(A') = \tau M h_1 d_0$	7.112, 7.131
(66, 36)	$\langle \eta^2, \theta_4, \eta^2, \theta_4 \rangle$	$\Delta_1 h_3^2$		Lemma 6.15	7.22
(68, 36)	$\langle \sigma, \kappa, \tau \eta \theta_{4.5} \rangle$	$h_3 A'$	0	$d_5(A') = \tau M h_1 d_0$	7.24, 7.113, Table 5
(69, 36)	$\langle \nu^2, 2, \theta_5 \rangle$	$h_2^2 h_6$		$d_2(h_6) = h_0 h_5^2$	7.154, Table 5
(70, 36)	$\langle 8\sigma, 2, \theta_5 \rangle$	$h_0^3 h_3 h_6$	$\tau h_1 p'$	$d_2(h_6) = h_0 h_5^2$	7.31, 7.74, 7.150, Table 5
(70, 37)	$\langle \eta, \nu, \tau \theta_{4.5} \bar{\kappa} \rangle$	$\tau h_1 D_3'$		Lemma 6.17	7.26, 7.116
(70, 37)	$\langle \eta, \tau \kappa^2, \tau \bar{\kappa}^2 \rangle$	$\tau \Delta^2 h_1^2 g +$ $+\tau^3 \Delta h_2^2 g^2$		$d_3(\Delta h_2^2) = \tau h_1 d_0^2$ $d_3(\tau \Delta^2 h_1 g) = \tau^4 e_0^4$	Table 5
(71, 37)	$\langle \epsilon, 2, \theta_5 \rangle$	$h_6 c_0$		$d_2(h_6) = h_0 h_5^2$	7.27, 7.28, 7.117, Table 5
(71, 39)	$\langle \nu, \epsilon, \kappa \theta_{4.5} \rangle$	$h_2^3 H_1$	$\tau M h_2^2 g$	$d_6(h_2^3 H_1) = M c_0 d_0$	7.77
(72, 37)	$\langle \mu_9, 2, \theta_5 \rangle$	$P h_1 h_6$		$d_2(h_6) = h_0 h_5^2$	7.31, Table 5
(72, 38)	$\langle \tau \bar{\kappa} \theta_{4.5}, 2\nu, \nu \rangle$	$h_0 d_0 D_2$	$h_1^2 h_3 h_6$	$d_2(P(A+A')) = \tau^2 M h_0 h_2 g$	7.80

Table 11: Some Toda brackets

(s, w)	bracket	detected by	indeterminacy	proof	used in
(72, 38)	$\langle \sigma^2, 2, \{t\}, \tau\bar{\kappa} \rangle$	$h_4 Q_2 + h_3^2 D_2$		Lemma 6.19	7.32
(75, 42)	$\langle \tau\eta\kappa\bar{\kappa}^2, \eta, \eta^2\eta_4 \rangle$	$\Delta h_2^2 d_0^2 e_0$	0	$\langle \Delta h_0^2 d_0 e_0, h_1, h_1^3 h_4 \rangle$	5.42
(77, 40)	$\langle \sigma^2, 2, \theta_5 \rangle$	$h_3^2 h_6$		$d_2(h_6) = h_0 h_5^2$ $d_2(h_4) = h_0 h_3^2$	7.123, Table 5
(77, 40)	$\langle \kappa, 2, \theta_5 \rangle$	$h_6 d_0$		Lemma 6.21	7.37, Table 5
(79, 41)	$\langle \eta_4, 2, \theta_5 \rangle$	$h_1 h_4 h_6$		$d_2(h_6) = h_0 h_5^2$	7.39, 7.126, Table 5
(79, 42)	$\langle \{h_1 x_{76,6}\}, 2, \eta \rangle$	$h_0 h_2 x_{76,6}$		$\langle h_1 x_{76,6}, h_0, h_1 \rangle =$ $= x_{76,6} \langle h_1, h_0, h_1 \rangle$	7.39
(80, 41)	$\langle \mu_{17}, 2, \theta_5 \rangle$	$P^2 h_1 h_6$		$d_2(h_6) = h_0 h_5^2$	Table 5
(80, 42)	$\langle 2, \eta, \tau\eta\{h_1 x_{76,6}\} \rangle$	$\tau h_1 x_1$		Lemma 6.22	Table 5
(82, 43)	$\langle \bar{\sigma}, 2, \theta_5 \rangle$	$h_6 c_1$		$d_2(h_6) = h_0 h_5^2$	7.90, Table 5
(82, 45)	$\langle \{\Delta e_1 + C_0\}, \eta^3, \eta_4 \rangle$	$(\Delta e_1 + C_0)g$	$h_1^3 x_1$	$\langle \Delta e_1 + C_0, h_1^3, h_1 h_4 \rangle$	7.130
(83, 44)	$\langle \eta_6 \kappa, \eta, \nu \rangle$	$h_0 h_6 g$		$d_2(h_6 e_0) = h_1^2 h_6 d_0$	Table 5, 7.91
(84, 44)	$\langle \nu\nu_4, 2, \theta_5 \rangle$	$h_2^2 h_4 h_6$		$d_2(h_6) = h_0 h_5^2$	7.134, Table 5
(85, 44)	$\langle \tau\eta^2\bar{\kappa}, 2, \theta_5 \rangle$	$P h_6 d_0$		$d_2(h_6) = h_0 h_5^2$	Table 5
(87, 46)	$\langle \theta_4, \tau\bar{\kappa}, \{t\} \rangle$	$\tau^2 g Q_3$	0	$d_3(Q_2) = \tau^2 g t$	5.65
(87, 46)	$\langle \tau\{h_0 Q_3 + h_0 n_1\}, \nu_4, \eta \rangle$	$h_1^2 c_3$		Lemma 6.26	7.50, Table 5
(88, 45)	$\langle \mu_{25}, 2, \theta_5 \rangle$	$P^3 h_1 h_6$		$d_2(h_6) = h_0 h_5^2$	Table 5
(89, 49)	$\langle \nu, \eta, \{h_2 g A'\} \rangle$	$\Delta h_1 g_2 g$	$\tau h_2^2 g C'$	Remark 7.139	7.138
(91, 49)	$\langle \{h_1 h_3 g\}, 2, \theta_5 \rangle$	$h_1 h_3 h_6 g$		$d_2(h_6) = h_0 h_5^2$	Table 5
(92, 48)	$\langle \theta_4, \theta_4, 2, \theta_4 \rangle$	$h_0 g_3$		Lemma 5.76	
(93, 48)	$\langle \theta_4, 2, \theta_5 \rangle$	$h_0^3 h_4^2 h_6$		$C\tau$	Table 5
(93, 50)	$\langle \kappa_1, \kappa, \tau\eta\theta_{4.5} \rangle$	$h_1^2 x_{91,8}$		$d_5(A') = \tau M h_1 d_0$	Table 5
(94, 49)	$\langle \{n\}, 2, \theta_5 \rangle$	$h_6 n$		$d_2(h_6) = h_0 h_5^2$	Table 5
(95, 49)	$\langle \{\Delta h_1 h_3\}, 2, \theta_5 \rangle$	$\Delta h_1 h_3 h_6$		$d_2(h_6) = h_0 h_5^2$	Table 5
(95, 50)	$\langle \kappa_1, 2, \theta_5 \rangle$	$h_6 d_1$		$d_2(h_6) = h_0 h_5^2$	Table 5

Table 12: Hidden values of inclusion of the bottom cell into $C\tau$

(s, f, w)	source	value	proof
(50, 6, 26)	τC	$\overline{Ph_1^2 h_5 c_0}$	
(57, 10, 30)	$h_0 h_2 h_5 i$	$\overline{\Delta^2 h_1^2 h_3}$	
(63, 6, 33)	$\tau h_1 H_1$	$\overline{h_1 B_7}$	
(63, 7, 33)	$\tau X_2 + \tau C'$	$\overline{h_5 d_0 e_0}$	
(64, 8, 34)	$\tau h_1 X_2$	$\overline{h_1 h_5 d_0 e_0}$	
(66, 8, 35)	$\tau h_2 C'$	τB_5	
(70, 7, 37)	$\tau h_1 h_3 H_1$	$\overline{h_1 h_3 B_7}$	
(70, 8, 37)	$\tau h_1 D'_3$	$\overline{h_1^2 X_3}$	
(70, 10, 38)	$h_1 h_3 (\Delta e_1 + C_0) + \tau h_2 C''$	$\overline{h_1^4 c_0 Q_2}$	Lemma 7.7
(71, 5, 37)	$\tau h_1 p_1$	$\overline{h_0^2 h_2 Q_3}$	
(74, 6, 39)	$h_3 (\tau Q_3 + \tau n_1)$	$\overline{h_1^4 p'}$	
(75, 11, 41)	$h_1^3 h_4 Q_2$	$\overline{\tau h_2 g^2 n}$	
(77, 6, 40)	$\tau h_1 h_4 D_3$	$\overline{x_{77,7}}$	
(81, 8, 43)	$\tau g D_3$	$\overline{h_1^4 x_{76,6}}$	
(83, 10, 45)	$h_2 c_1 A'$	$\overline{h_1 g B_7}$	
(85, 5, 44)	$\tau h_1 f_2$	$\overline{h_0^3 c_3}$	
(86, 6, 45)	$\tau h_1^2 f_2$	$\overline{\tau h_1^3 h_4 Q_3}$	
(86, 6, 45)	$\tau h_2 h_6 g + ? \tau h_1^2 f_2$	$\overline{? \Delta^2 e_1 + \tau \Delta h_2 e_1 g}$	
(86, 7, 45)	$\tau h_1 x_{85,6}$	$\overline{? \Delta^2 e_1 + \tau \Delta h_2 e_1 g}$	
(86, 12, 47)	$\tau h_2 g C'$	$\tau B_5 g$	
(87, 10, 45)	$\tau \Delta h_1 H_1$	$\overline{? \Delta h_1 B_7}$	
(88, 11, 48)	$\tau h_2^2 g H_1$	$\overline{\Delta g_2 g}$	
(90, 14, 50)	$\tau h_2 g C'''$	$\overline{Ph_1^{10} h_6 c_0}$	
(90, 19, 49)	$\tau^3 g m^2$	$\overline{\Delta^2 c_0 e_0^2}$	

Table 13: Hidden values of projection from $C\tau$ to the top cell

(s, f, w)	source	value	crossing source
(30, 6, 16)	Δh_2^2	$h_1 d_0^2$	
(34, 2, 18)	$h_2 h_5$	$h_1 d_1$	
(38, 7, 20)	$h_0 y$	$\tau h_0 e_0 g$	
(41, 4, 22)	$h_0 c_2$	$h_1 h_3 d_1$	
(44, 10, 24)	$\Delta h_2^2 d_0$	$h_1 d_0^3$	
(50, 10, 28)	$\Delta h_2^2 g$	$h_1 d_0 e_0^2$	
(55, 7, 30)	B_6	$h_2 g n$	
(56, 10, 29)	$\Delta^2 h_1 h_3$	$\Delta h_0^2 d_0 e_0$	
(57, 7, 30)	Q_2	$\tau g t$	
(58, 7, 30)	$h_0 D_2$	$\Delta h_1 d_1$	
(58, 11, 32)	$Ph_1^2 h_5 e_0$	$\tau h_2 e_0^2 g$	
(59, 8, 33)	$h_1^2 D_4$	$h_2^2 d_1 g$	
(61, 6, 32)	A'	$M h_1 d_0$	$\overline{h_1^3 Q_2}$
(62, 11, 32)	$\overline{Ph_5 c_0 d_0}$	$\tau \Delta h_2^2 d_0 e_0$	

Table 13: Hidden values of projection from $C\tau$ to the top cell

(s, f, w)	source	value	crossing source
(63, 12, 33)	$\overline{Ph_1h_5c_0d_0}$	$\tau^2d_0^2e_0^2$	
(64, 14, 36)	km	$h_1d_0^2e_0^2$	
(65, 6, 35)	h_2H_1	d_1^2	
(65, 8, 34)	$h_0h_3D_2$	$h_1^2(\Delta e_1 + C_0)$	
(66, 3, 34)	$h_0h_2h_6$	$h_1^3h_5^2$	
(68, 5, 36)	h_0d_2	$h_1 \cdot \Delta_1h_3^2$	
(68, 7, 37)	$\overline{h_2^2H_1}$	$h_1^4X_2$	$\overline{h_1h_3j_1}$
(68, 12, 35)	$\overline{\tau Mc_0d_0}$	$\tau^2\Delta h_2^2e_0g$	
(68, 12, 36)	$\overline{h_5d_0i}$	$\Delta h_1d_0^3$	
(69, 9, 36)	$\overline{h_1X_3}$	τMh_2g	
(69, 10, 36)	$P(A + A')$	τMh_0h_2g	
(69, 13, 36)	$\tau\Delta^2h_1g$	$\tau^3e_0^4$	
(70, 4, 36)	$h_0^2h_3h_6$	$\tau h_2^2C'$	
(70, 14, 40)	m^2	$h_1e_0^4$	
(72, 9, 40)	$h_2^2G_0$	g^2n	
(72, 10, 38)	d_0D_2	τMh_2^2g	
(73, 7, 38)	h_4D_2	$\tau^3d_1g^2$	
(74, 6, 38)	Ph_2h_6	$h_1h_4Q_2$	
(75, 11, 40)	$h_2d_0D_2$	$Mh_1d_0^2$	
(75, 11, 42)	gB_6	h_2g^2n	
(76, 14, 41)	$\Delta^2h_1h_3g$	$\Delta h_2^2d_0^2e_0$	
(77, 11, 41)	τgQ_2	τ^2g^2t	
(78, 7, 42)	$h_3x_{71,6}$	$h_1d_1g_2$	
(78, 13, 40)	$h_0^3x_{78,10}$	$\tau^5e_0g^3$	
(80, 5, 42)	h_0e_2	$h_1^3x_{76,6}$	
(80, 14, 41)	$\Delta^3h_1h_3$	$\tau^3\Delta h_1e_0^2g$	
(81, 10, 44)	gA'	$Mh_1e_0^2$	$\overline{Ph_1^9h_6}$
(81, 15, 42)	$\Delta^3h_1^2h_3$	$\tau^3d_0e_0^2l$	
(82, 16, 44)	$\Delta^2e_0^2$	$\tau\Delta h_2^2e_0^3$	
(83, 5, 43)	τh_6g	$?\tau(\Delta e_1 + C_0)g$	$\overline{Ph_1^3h_6c_0}$
(83, 17, 45)	$\Delta^2h_1e_0^2$	$\tau^2d_0e_0^4$	
(84, 4, 44)	f_2	$h_1h_4Q_3$	
(84, 18, 48)	d_0m^2	$h_1d_0e_0^4$	
(85, 5, 45)	h_1f_2	$h_1^2h_4Q_3$	
(85, 6, 45)	$x_{85,6}$	$?\Delta h_1j_1$	
(85, 10, 47)	h_2gH_1	d_1^2g	
(86, 4, 45)	h_1c_3	$h_0h_2h_4Q_3$	
(86, 6, 46)	$h_1^2f_2$	$h_1^3h_4Q_3$	
(86, 6, 46)	h_2h_6g	$?\tau Mh_0g^2$	
(86, 7, 46)	$h_1x_{85,6}$	$?\tau Mh_0g^2$	
(86, 12, 45)	$\Delta^2e_1 + \tau\overline{\Delta h_2e_1g}$	$?M\Delta h_1^2d_0$	$\overline{P^2h_1^6h_6}$
(86, 15, 44)	$\Delta^3h_0h_3^2$	$\tau\Delta^2h_0e_0g$	
(87, 7, 45)	$\overline{x_{87,7}}$	$?M\Delta h_0^2e_0$	
(87, 11, 45)	$\overline{\Delta h_1B_7}$	$?M\Delta h_0^2e_0$	

Table 13: Hidden values of projection from $C\tau$ to the top cell

(s, f, w)	source	value	crossing source
(87, 17, 45)	$\Delta^3 h_1 d_0$	$\tau^4 e_0^3 m$	
(88, 11, 49)	$h_2^2 g H_1$	$M h_1^2 g^2$	
(88, 16, 47)	$\tau \Delta^2 g^2$	$\tau^2 \Delta h_2^2 e_0 g^2$	
(88, 17, 48)	$\Delta^2 h_0 g^2$	$\Delta h_1 d_0^2 e_0^2$	
(89, 12, 46)	$\Delta^2 h_0 c_2$	$\Delta^2 h_1 h_3 d_1$	
(89, 17, 48)	$\tau \Delta^2 h_1 g^2$	$\tau^3 e_0^4 g$	
(90, 18, 52)	$g m^2$	$h_1 e_0^4 g$	

Table 14: Hidden τ extensions in the \mathbb{C} -motivic Adams spectral sequence

(s, f, w)	from	to	proof
(22, 7, 13)	$c_0 d_0$	$P d_0$	
(23, 8, 14)	$h_1 c_0 d_0$	$P h_1 d_0$	
(28, 6, 17)	$h_1 h_3 g$	d_0^2	
(29, 7, 18)	$h_1^2 h_3 g$	$h_1 d_0^2$	
(40, 9, 23)	$\tau h_0 g^2$	$\Delta h_1^2 d_0$	
(41, 9, 23)	$\tau^2 h_1 g^2$	$\Delta h_0^2 e_0$	
(42, 11, 25)	$c_0 e_0^2$	d_0^3	
(43, 12, 26)	$h_1 c_0 e_0^2$	$h_1 d_0^3$	
(46, 6, 26)	$h_1^2 g_2$	$\Delta h_2 c_1$	
(47, 12, 26)	$\Delta h_1 c_0 d_0$	$P \Delta h_1 d_0$	
(48, 10, 29)	$h_1 h_3 g^2$	$d_0 e_0^2$	
(49, 11, 30)	$h_1^2 h_3 g^2$	$h_1 d_0 e_0^2$	
(52, 10, 29)	$\Delta h_1 h_3 g$	$\tau^2 e_0 m$	
(53, 9, 29)	$M c_0$	$M P$	
(53, 11, 30)	$\Delta h_1^2 h_3 g$	$\Delta h_1 d_0^2$	
(54, 8, 31)	$h_1 i_1$	$M h_1 c_0$	
(54, 10, 30)	$M h_1 c_0$	$M P h_1$	
(54, 11, 32)	$h_1^6 h_5 e_0$	$\tau e_0^2 g$	
(55, 12, 33)	$h_1^7 h_5 e_0$	$\tau h_1 e_0^2 g$	
(55, 13, 31)	$\tau^2 h_1 e_0^2 g$	$\Delta h_0^2 d_0 e_0$	
(59, 7, 33)	j_1	$M d_0$	
(59, 12, 33)	$P h_1^3 h_5 e_0$	$\tau \Delta h_1 d_0 g$	
(60, 9, 34)	$h_1^3 D_4$	$M h_1 d_0$	
(60, 13, 34)	$\tau^2 h_0 g^3$	$\Delta c_0 d_0^2 + \tau d_0^2 l$	
(61, 13, 35)	$\tau^2 h_1 g^3$	$\Delta h_2^2 d_0 e_0$	
(62, 14, 37)	$h_1^6 h_5 c_0 e_0$	$d_0^2 e_0^2$	
(63, 15, 38)	$h_0^2 h_2 g^3$	$h_1 d_0^2 e_0^2$	
(65, 9, 36)	$h_1^2 X_2$	$\tau M g$	
(66, 10, 37)	$h_1^3 X_2$	$\tau M h_1 g$	
(66, 14, 37)	$P h_1^2 h_5 c_0 e_0$	$\tau^2 d_0 e_0 m$	
(67, 15, 38)	$P h_1^3 h_5 c_0 e_0$	$\Delta h_1 d_0^3$	

Table 14: Hidden τ extensions in the \mathbb{C} -motivic Adams spectral sequence

(s, f, w)	from	to	proof
(68, 14, 41)	$h_1 h_3 g^3$	e_0^4	
(69, 15, 42)	$h_1^2 h_3 g^3$	$h_1 e_0^4$	
(70, 8, 39)	$d_1 e_1$	$h_1 h_3 (\Delta e_1 + C_0)$	Lemma 7.7
(70, 10, 38)	$\tau h_2 C'' + h_1 h_3 (\Delta e_1 + C_0)$	$\Delta^2 h_2 c_1$	
(71, 8, 39)	$h_2^3 H_1$	$h_3^2 Q_2$	
(72, 7, 39)	$h_1 x_{71,6}$	$h_0 d_0 D_2$	
(72, 14, 41)	$\Delta h_1 h_3 g^2$	$\tau^2 e_0 g m$	
(73, 6, 39)	$h_1^2 h_6 c_0$	$h_0 h_4 D_2$	
(73, 11, 40)	$\tau h_2^2 C''$	$\Delta^2 h_1^2 h_4 c_0$	
(73, 12, 41)	$M h_1 h_3 g$	$M d_0^2$	
(73, 15, 42)	$\Delta h_1^2 h_3 g^2$	$\Delta h_1 d_0 e_0^2$	
(74, 13, 42)	$M h_1^2 h_3 g$	$M h_1 d_0^2$	
(75, 17, 43)	$\tau^2 h_1 e_0^2 g^2$	$\Delta h_2^2 d_0^2 e_0$	
(77, 15, 42)	$\Delta^2 h_3^2 g$	$\tau^5 e_0 g^3$	
(78, 8, 43)	$h_1 m_1$	$M \Delta h_1^2 h_3$	
(79, 11, 45)	$g j_1$	$M e_0^2$	
(80, 12, 46)	$h_1 g j_1$	$M h_1 e_0^2$	
(80, 17, 47)	$\tau h_0 g^4$	$\Delta h_1^2 e_0^2 g$	
(80, 18, 46)	$\Delta h_1^2 e_0^2 g$	$\Delta c_0 d_0 e_0^2 + \tau d_0 e_0^2 t$	
(81, 17, 47)	$\tau^2 h_1 g^4$	$\Delta h_2^2 e_0^3$	
(82, 12, 44)	$\tau (\Delta e_1 + C_0) g$	$? \Delta^2 h_2 n$	
(82, 19, 49)	$c_0 e_0^2 g^2$	$d_0 e_0^4$	
(83, 20, 50)	$h_1 c_0 e_0^2 g^2$	$h_1 d_0 e_0^4$	
(84, 12, 46)	$\Delta h_1 j_1$	$? M \Delta h_1 d_0$	
(85, 8, 45)	$h_6 c_0 d_0$	$P h_6 d_0$	
(85, 15, 47)	$\tau M h_0 g^2$	$? M \Delta h_1^2 d_0$	
(86, 9, 46)	$h_1 h_6 c_0 d_0$	$P h_1 h_6 d_0$	
(86, 12, 47)	$\tau h_2 g C'$	$\Delta^2 h_3^2 d_1$	
(86, 15, 47)	$\tau^2 M h_1 g^2$	$M \Delta h_0^2 e_0$	
(87, 20, 50)	$\Delta h_1 c_0 e_0^2 g$	$\Delta h_1 d_0^2 e_0^2$	
(88, 18, 53)	$h_1 h_3 g^4$	$e_0^4 g + M h_1^6 e_0 g$	
(89, 13, 49)	$\tau h_2^2 g C'$	$\Delta^2 h_1^2 h_3 d_1$	
(89, 19, 54)	$h_1^2 h_3 g^4$	$h_1 e_0^4 g + M h_1^7 e_0 g$	
(90, 14, 50)	$\tau h_2 g C''$	$\Delta^2 h_2 c_1 g$	

Table 15: Hidden 2 extensions

(s, f, w)	source	target	proof	notes
(23, 6, 13)	$\tau h_0 h_2 g$	$Ph_1 d_0$	τ	
(23, 6, 14)	$h_0 h_2 g$	$h_1 c_0 d_0$	$C\tau$	
(40, 8, 22)	$\tau^2 g^2$	$\Delta h_1^2 d_0$	τ	
(43, 10, 25)	$\tau h_0 h_2 g^2$	$h_1 d_0^3$	τ	
(43, 10, 26)	$h_0 h_2 g^2$	$h_1 c_0 e_0^2$	$C\tau$	
(47, 10, 25)	$\tau \Delta h_2^2 e_0$	$P \Delta h_1 d_0$	τ	
(47, 10, 26)	$\Delta h_2^2 e_0$	$\Delta h_1 c_0 d_0$	$C\tau$	
(51, 6, 28)	$h_0 h_3 g_2$	$\tau g n$	[45]	
(54, 9, 28)	$h_0 h_5 i$	$\tau^4 e_0^2 g$	Lemma 7.19, [5]	
(60, 12, 33)	$\tau^3 g^3$	$\Delta c_0 d_0^2 + \tau d_0^2 l$	τ	
(63, 6, 33)	$\tau h_1 H_1$	$\tau h_1 (\Delta e_1 + C_0)$	Lemma 7.20	
(63, 14, 37)	$\tau h_0 h_2 g^3$	$h_1 d_0^2 e_0^2$	τ	
(64, 2, 33)	$h_1 h_6$	$\tau h_1^2 h_5^2$	Lemma 7.21	
(65, 9, 36)	$h_1^2 X_2$	$M h_0 g$	τ	
(67, 14, 37)	$\tau \Delta h_2^2 e_0 g$	$\Delta h_1 d_0^3$	τ	
(70, 7, 37)	$\tau h_1 h_3 H_1$	$\tau h_1 h_3 (\Delta e_1 + C_0)$	Lemma 7.20	
(71, 4, 37)	$h_6 c_0$	$\tau h_1^2 p'$	Lemma 7.28	
(71, 8, 39)	$h_2^3 H_1$	$\tau M h_2^2 g$	Lemma 7.30	
(74, 6, 39)	$h_3 (\tau Q_3 + \tau n_1)$	$\tau x_{74,8}$	Lemma 7.36	indet
(74, 10, 41)	$h_3 C''$	$M h_1 d_0^2$	$C\tau$	
(74, 14, 40)	$\Delta^2 h_2^2 g$	$\tau^4 e_0^2 g^2$	mmf	
(77, 6, 41)	$\tau h_1 h_4 D_3$	$h_0 x_{77,7}$	$C\tau$	
(78, 10, 42)	$e_0 A'$	$M \Delta h_1^2 h_3$	Lemma 7.38	
(80, 16, 45)	$\tau^3 g^4$	$\Delta c_0 d_0 e_0^2 + \tau d_0 e_0^2 l$	τ	
(80, 16, 46)	$\tau^2 g^4$	$\Delta h_1^2 e_0^2 g + M h_1^3 d_0 e_0$	τ	
(83, 18, 49)	$\tau h_0 h_2 g^4$	$h_1 d_0 e_0^4$	τ	
(83, 18, 50)	$h_0 h_2 g^4$	$h_1 c_0 e_0^2 g^2$	$C\tau$	
(85, 5, 44)	$\tau h_1 f_2$	$\tau^2 h_2 h_4 Q_3$	Lemma 7.45	
(85, 14, 46)	$\tau^2 M g^2$	$? M \Delta h_1^2 d_0$	τ	
(86, 7, 45)	$\tau h_0 h_2 h_6 g$	$Ph_1 h_6 d_0$	τ	
(86, 7, 46)	$h_0 h_2 h_6 g$	$h_1 h_6 c_0 d_0$	$C\tau$	
(87, 9, 48)	$g Q_3$	$B_6 d_1$	$C\tau$	
(87, 10, 45)	$? \tau \Delta h_1 H_1$	$? \tau \Delta^2 h_3 d_1$ $? \tau^2 \Delta^2 c_1 g$	Lemma 7.51	
(87, 18, 49)	$\tau \Delta h_2^2 e_0 g^2$	$\Delta h_1 d_0^2 e_0^2$	$C\tau$	
(87, 18, 50)	$\Delta h_2^2 e_0 g^2$	$\Delta h_1 c_0 e_0^2 g + M h_1^3 c_0 d_0 e_0$	$C\tau$	
(90, 10, 50)	$h_2 g Q_3$	$\tau d_1 e_1 g$	$C\tau$	

Table 16: Possible hidden 2 extensions

(s, f, w)	source	target
(59, 7, 33)	j_1	$?\tau^2 c_1 g^2$
(72, 7, 39)	$h_1 x_{71,6}$	$?\tau^3 d_1 g^2$
(79, 11, 45)	gj_1	$?\tau^2 c_1 g^3$
(82, 6, 44)	$h_5^2 g$	$?\tau(\Delta e_1 + C_0)g$
(82, 8, 43)	$\tau h_2^2 x_{76,6}$	$?\Delta^2 h_2 n$
(82, 8, 44)	$h_2^2 x_{76,6}$	$?\tau(\Delta e_1 + C_0)g$
(85, 6, 45)	$x_{85,6}$	$?\tau P h_1 x_{76,6}$
(86, 7, 45)	$?\tau h_1 x_{85,6}$	$?P h_1 h_6 d_0$
		$?\tau \Delta^2 h_2^2 d_1$
		$?\tau M \Delta h_0^2 e_0$
(87, 7, 45)	$x_{87,7}$	$?\tau^3 g Q_3$
		$?\tau \Delta^2 h_3 d_1$
		$?\tau^2 \Delta^2 c_1 g$

Table 17: Hidden η extensions

(s, f, w)	source	target	proof	notes
(15, 4, 8)	$h_0^3 h_4$	Pc_0	$C\tau$	
(21, 5, 11)	$\tau^2 h_1 g$	Pd_0	τ	
(21, 5, 12)	$\tau h_1 g$	$c_0 d_0$	$C\tau$	
(23, 9, 12)	$h_0^2 i$	$P^2 c_0$	$C\tau$	
(31, 11, 16)	$h_0^{10} h_5$	$P^3 c_0$	$C\tau$	
(38, 4, 20)	$h_0^2 h_3 h_5$	$\tau^2 c_1 g$	[16, Table 29]	
(39, 17, 20)	$P^2 h_0^2 i$	$P^4 c_0$	$C\tau$	
(40, 8, 21)	$\tau^3 g^2$	$\Delta h_0^2 e_0$	τ	
(41, 5, 23)	$h_1 f_1$	$\tau h_2 c_1 g$	[16, Table 29]	crossing
(41, 9, 23)	$\tau^2 h_1 g^2$	d_0^3	τ	
(41, 9, 24)	$\tau h_1 g^2$	$c_0 e_0^2$	$C\tau$	
(41, 10, 22)	$\Delta h_0^2 e_0$	τd_0^3	τ	
(45, 3, 24)	$h_3^2 h_5$	Mh_1	[16, Table 29]	crossing
(45, 5, 24)	$\tau h_1 g_2$	$\Delta h_2 c_1$	τ	
(45, 9, 24)	$\tau \Delta h_1 g$	$\tau d_0 l + \Delta c_0 d_0$	$C\tau$	
(46, 11, 24)	$\tau^2 d_0 l$	$P \Delta h_1 d_0$	τ	
(47, 10, 26)	$\Delta h_2^2 e_0$	$\tau d_0 e_0^2$	mmf	
(47, 20, 24)	$h_0^7 Q'$	$P^5 c_0$	$C\tau$	
(50, 6, 26)	τC	$\tau^2 g n$	[45]	
(52, 11, 28)	$\tau^2 e_0 m$	$\Delta h_1 d_0^2$	τ	
(54, 12, 29)	$\tau^3 e_0^2 g$	$\Delta h_0^2 d_0 e_0$	τ	
(55, 25, 28)	$P^4 h_0^2 i$	$P^6 c_0$	$C\tau$	
(59, 13, 31)	$\tau \Delta h_1 d_0 g$	$\Delta c_0 d_0^2 + \tau d_0^2 l$	τ	
(60, 12, 33)	$\tau^3 g^3$	$\Delta h_2^2 d_0 e_0$	τ	
(61, 9, 35)	$h_1^2 j_1$	$\tau h_2 c_1 g^2$	$C\tau$	crossing
(61, 13, 35)	$\tau^2 h_1 g^3$	$d_0^2 e_0^2$	τ	
(61, 14, 34)	$\Delta h_2^2 d_0 e_0$	$\tau d_0^2 e_0^2$	τ	
(63, 6, 33)	$\tau h_1 H_1$	$h_3 Q_2$	$C\tau$	indet
(63, 26, 32)	$h_0^{25} h_6$	$P^7 c_0$	$C\tau$	
(64, 8, 34)	$\tau h_1 X_2$	$c_0 Q_2$	$C\tau$	
(64, 8, 33)	$\tau^2 h_1 X_2$	$\tau^2 M h_0 g$	Lemma 7.66	indet
(65, 13, 35)	$\tau^2 \Delta h_1 g^2$	$\tau^2 d_0 e_0 m$	τ	
(66, 15, 36)	$\tau^2 d_0 e_0 m$	$\Delta h_1 d_0^3$	τ	
(67, 14, 38)	$\Delta h_2^2 e_0 g$	τe_0^4	$C\tau$	
(68, 7, 36)	$h_3 A'$	$h_3 (\Delta e_1 + C_0)$	Lemma 7.70	
(69, 3, 36)	$h_2^2 h_6$	$\tau h_0 h_2 Q_3$	Lemma 7.71	crossing
(70, 7, 37)	$\tau h_1 h_3 H_1$	$h_3^2 Q_2$	Lemma 7.72	
(70, 9, 37)	$\tau h_1 D'_3$	$d_0 Q_2$	$C\tau$	
(71, 5, 37)	$\tau h_1 p_1$	$h_4 Q_2$	$C\tau$	
(71, 13, 38)	$\Delta^2 h_2 g$	$\tau^3 e_0 g m$	mmf	
(71, 33, 36)	$P^6 h_0^2 i$	$P^8 c_0$	$C\tau$	
(72, 5, 37)	$\tau h_1 h_6 c_0$	$\tau^2 h_2^2 Q_3$	Lemma 7.78	indet
(72, 11, 38)	$h_0 d_0 D_2$	$\tau M d_0^2$	Lemma 7.80	indet
(72, 15, 40)	$\tau^2 e_0 g m$	$\Delta h_1 d_0 e_0^2$	τ	

Table 17: Hidden η extensions

(s, f, w)	source	target	proof	notes
(74, 16, 41)	$\tau^3 e_0^2 g^2$	$\Delta h_2^2 d_0^2 e_0$	τ	
(75, 6, 40)	$h_0 h_3 d_2$	$\tau d_1 g_2$	Lemma 7.81	
(75, 10, 42)	$h_1^4 x_{71,6}$	$h_1 g B_6$	$C\tau$	
(75, 11, 41)	$h_1^3 h_4 Q_2$	$\tau^2 g^2 t$	$C\tau$	
(76, 9, 40)	$x_{76,9}$	$M\Delta h_1 h_3$	$C\tau$	
(77, 6, 40)	$\tau h_1 h_4 D_3$	$x_{78,9}$	$C\tau$	crossing
(78, 8, 40)	$h_0^6 h_4 h_6$	$\tau\Delta B_6$	$C\tau$, [5]	
(78, 10, 42)	$e_0 A'$	$\tau M e_0^2$	[5]	
(79, 17, 44)	$\tau\Delta h_1 e_0^2 g$	$\Delta c_0 d_0 e_0^2 + \tau d_0 e_0^2 l$	τ	
(79, 36, 40)	$P^4 h_0^7 Q'$	$P^9 c_0$	$C\tau$	
(80, 16, 45)	$\tau^3 g^4$	$\Delta h_2^2 e_0^3$	τ	
(81, 13, 47)	$h_1^2 g j_1$	$\tau h_2 c_1 g^3$	$C\tau$	crossing
(81, 17, 47)	$\tau^2 h_1 g^4$	$d_0 e_0^4$	τ	
(81, 17, 48)	$\tau h_1 g^4$	$c_0 e_0^2 g^2$	$C\tau$	
(81, 18, 46)	$\Delta h_2^2 e_0^3$	$\tau d_0 e_0^4$	τ	
(83, 11, 44)	$\tau\Delta j_1 + \tau^2 g C'$	$?M\Delta h_1 d_0$	τ	crossing
(84, 6, 43)	$\tau^2 h_1 h_6 g$	$P h_6 d_0$	τ	
(84, 6, 44)	$\tau h_1 h_6 g$	$h_6 c_0 d_0$	$C\tau$	
(85, 14, 45)	$\tau^3 M g^2$	$M\Delta h_0^2 e_0$	τ	
(85, 17, 48)	$\tau\Delta h_1 g^3$	$\Delta c_0 e_0^2 g + M h_1^2 c_0 d_0 e_0$ $+ \tau e_0^3 m$	$C\tau$	
(86, 11, 44)	$h_0^3 h_6 i$	$\tau^2 \Delta^2 c_1 g$	Lemma 7.95	
(86, 16, 47)	$P^2 h_1^7 h_6$	$\tau^2 \Delta h_2^2 e_0 g^2$	$C\tau$	
(86, 19, 48)	$\tau^2 e_0^3 m$	$\Delta h_1 d_0^2 e_0^2$	τ	
(87, 8, 47)	$h_1^2 x_{85,6}$	$x_{88,10}$	$C\tau$	
(87, 10, 45)	$? \tau \Delta h_1 H_1$	$\Delta^2 f_1 + ? \tau^2 \Delta g_2 g$	$C\tau$	
(87, 18, 50)	$\Delta h_2^2 e_0 g^2$	$\tau e_0^4 g$	$C\tau$	
(87, 41, 44)	$P^8 h_0^2 i$	$P^{10} c_0$	$C\tau$	
(88, 11, 48)	$\tau h_2^2 g H_1$	$\Delta h_1 g_2 g + \tau h_2^2 g C'$	$C\tau$	
(89, 13, 47)	$\Delta^2 h_1 f_1$	$\tau \Delta^2 h_2 c_1 g$	Lemma 7.98	crossing

Table 18: Possible hidden η extensions

(s, f, w)	source	target	proof
(64, 5, 34)	h_2D_3	$? \tau k_1$	
(66, 6, 35)	$\tau \Delta_1 h_3^2$	$? \tau^2 \Delta h_2^2 e_0 g$	
(66, 12, 35)	$\Delta^2 h_1^3 h_4$	$? \tau^2 \Delta h_2^2 e_0 g$	
(67, 6, 36)	$h_0 Q_3 + h_2^2 D_3$	$? \tau M h_0 h_2 g$	
(81, 3, 42)	$h_2 h_4 h_6$	$? \tau h_5^2 g$	Lemma 7.86
		$? \tau^2 e_1 g_2$	
		$? \Delta^2 h_2 n$	
(81, 5, 39)	$h_1^3 h_4 h_6$	$? \tau (\Delta e_1 + C_0) g$	Lemma 7.87
(81, 8, 42)	$\tau^2 g D_3$	$? \Delta^2 h_2 n$	
(81, 8, 43)	$\tau g D_3$	$? \tau (\Delta e_1 + C_0) g$	
(86, 6, 45)	$\tau h_1^2 f_2$	$? \tau^2 g Q_3$	
		$? \Delta^2 h_3 d_1$	
(86, 6, 45)	$\tau h_2 h_6 g + ? \tau h_1^2 f_2$	$? \tau^2 g Q_3$	Lemma 7.93
		$? \Delta^2 h_3 d_1$	
(86, 6, 46)	$h_2 h_6 g + ? h_1^2 f_2$	$? h_1^2 x_{85,6}$	
		$? \tau h_2^2 g A'$	
(86, 7, 45)	$\tau h_1 x_{85,6}$	$? \tau^2 g Q_3 ?$	Lemma 7.93
		$? \Delta^2 h_3 d_1 ?$	
(87, 5, 46)	$h_1^2 c_3$	$? \tau h_0 g_2^2$	
(87, 6, 45)	$\tau h_1 h_4 h_6 c_0$	$? \tau^2 h_0 g_2^2$	
(87, 9, 48)	$g Q_3$	$? \tau M h_0 h_2 g^2$	
(88, 8, 48)	g_2^2	$? \tau h_2^2 g C'$	
		$? \Delta h_1 g_2 g$	
(89, 7, 48)	$h_2^2 h_6 g$	$? \tau h_2 g Q_3$	

Table 19: Hidden ν extensions

(s, f, w)	source	target	proof	notes
(20, 6, 11)	$\tau h_0^2 g$	$Ph_1 d_0$	τ	
(20, 6, 12)	$h_0^2 g$	$h_1 c_0 d_0$	$C\tau$	
(22, 4, 13)	$h_2 c_1$	$h_1^2 h_4 c_0$	$C\tau$	
(26, 6, 15)	$\tau h_2^2 g$	$h_1 d_0^2$	τ	
(30, 2, 16)	h_4^2	p	$C\tau$	
(32, 6, 17)	$\Delta h_1 h_3$	$\tau^2 h_1 e_0^2$	tmf	
(39, 9, 21)	$\Delta h_1 d_0$	τd_0^3	tmf	
(40, 10, 23)	$\tau h_0^2 g^2$	$h_1 d_0^3$	τ	
(40, 10, 24)	$h_0^2 g^2$	$h_1 c_0 e_0^2$	$C\tau$	
(42, 8, 25)	$h_2 c_1 g$	$h_1^6 h_5 c_0$	$C\tau$	
(45, 3, 24)	$h_3^2 h_5$	Mh_2	$C\tau$	crossing
(45, 4, 24)	$h_0 h_3^2 h_5$	$Mh_0 h_2$	$C\tau$	crossing
(45, 9, 24)	$\tau \Delta h_1 g$	$\tau^2 d_0 e_0^2$	mmf	
(46, 10, 27)	$\tau h_2^2 g^2$	$h_1 d_0 e_0^2$	τ	
(48, 6, 26)	$h_2 h_5 d_0$	τgn	[45]	crossing
(51, 8, 27)	τMh_2^2	MPh_1	τ	
(51, 8, 28)	Mh_2^2	$Mh_1 c_0$	[16, Table 31]	
(52, 10, 29)	$\Delta h_1 h_3 g$	$\tau^2 h_1 e_0^2 g$	τ	
(52, 11, 28)	$\tau^2 e_0 m$	$\Delta h_0^2 d_0 e_0$	mmf	
(53, 7, 30)	i_1	gt	$C\tau$	
(54, 11, 32)	$h_1^6 h_5 e_0$	$h_2 e_0^2 g$	τ	
(57, 10, 30)	$h_0 h_2 h_5 i$	$\tau^2 d_0^2 l$	tmf	
(59, 12, 33)	$Ph_1^3 h_5 e_0$	$\tau d_0^2 e_0^2$	τ	
(59, 13, 32)	$\tau \Delta h_1 d_0 g$	$\tau^2 d_0^2 e_0^2$	mmf	
(60, 14, 35)	$\tau h_0^2 g^3$	$h_1 d_0^2 e_0^2$	τ	
(62, 8, 33)	$\Delta e_1 + C_0$	$\tau Mh_0 g$	Lemma 7.106	
(62, 12, 37)	$h_2 c_1 g^2$	$h_1^8 D_4$	$C\tau$	
(63, 6, 33)	$\tau h_1 H_1$	$\tau^2 Mh_1 g$	Lemma 7.108	crossing
(65, 3, 34)	$h_2 h_5^2$	$\tau h_1 Q_3$	$C\tau$	
(65, 9, 36)	$h_1^2 X_2$	$Mh_2 g$	τ	
(65, 13, 36)	$\tau \Delta h_1 g^2 + Ph_1 h_5 c_0 e_0$	$\tau^2 e_0^4$	mmf	
(66, 6, 36)	$\Delta_1 h_3^2$	$h_2^2 C'$	$C\tau$	
(66, 14, 39)	$\tau h_2^2 g^3$	$h_1 e_0^4$	τ	
(67, 8, 36)	$h_2^2 A'$	$h_1 h_3 (\Delta e_1 + C_0)$	Lemma 7.112	
(68, 13, 36)	$Ph_2 h_5 j$	$\Delta^2 h_0^2 h_2 g$	$C\tau$	
(69, 9, 38)	$h_2^2 C'$	$\tau^2 d_1 g^2$	Lemma 7.115	
(70, 9, 37)	$\tau h_1 D_3'$	$\tau M d_0^2$	Lemma 7.116	indet
(70, 12, 37)	$\Delta^2 h_2 c_1$	$\Delta^2 h_1^2 h_4 c_0$	τ	
(70, 14, 37)	$\tau \Delta^2 h_1^2 g + \tau^3 m^2$	$\tau^2 \Delta h_1 d_0 e_0^2$	mmf	
(71, 8, 39)	$h_3^3 H_1$	$h_3 C''$	$C\tau$	
(71, 12, 39)	$\tau Mh_2^2 g$	$Mh_1 d_0^2$	τ	
(71, 14, 38)	$\Delta^2 h_0 h_2 g$	$\tau^4 e_0^2 g^2$	mmf	
(72, 14, 41)	$\Delta h_1 h_3 g^2$	$\tau^2 h_1 e_0^2 g^2$	τ	
(72, 15, 40)	$\tau^2 e_0 g m$	$\Delta h_2^2 d_0^2 e_0$	mmf	
(73, 11, 41)	$h_2^2 C''$	$\tau g^2 t$	Lemma 7.118	

Table 19: Hidden ν extensions

(s, f, w)	source	target	proof	notes
(74, 14, 39)	$\tau\Delta^2h_2^2g$	$\tau^5e_0g^3$	τ	
(77, 3, 40)	$h_3^2h_6$	τh_1x_1	Lemma 7.123	indet
(77, 7, 41)	$h_1x_{76,6}$	c_1A'	$C\tau$	
(77, 15, 42)	$\Delta^2h_2^3g$	$\tau^2d_0e_0^2l$	τ	
(77, 16, 41)	$\tau^5e_0g^3$	$\tau^3d_0e_0^2l$	mmf	
(78, 8, 40)	$h_0^6h_4h_6$	Δ^2p	$C\tau$	
(78, 9, 40)	$h_0^7h_4h_6$	$\tau\Delta^2h_1d_1$	Lemma 7.124	
(79, 17, 45)	$\Delta h_1e_0^2g + Mh_1^3d_0e_0$	$\tau d_0e_0^4$	mmf	
(80, 18, 47)	$\tau h_0^2g^4$	$h_1d_0e_0^4$	τ	
(80, 18, 48)	$h_0^2g^4$	$h_1c_0e_0^2g^2$	$C\tau$	
(82, 6, 44)	h_5^2g	$h_0h_2h_4Q_3$	$C\tau$	
(82, 8, 44)	$h_2^2x_{76,6}$	$Ph_1x_{76,6}$	$C\tau$	
(82, 10, 42)	$P^2h_2h_6$	Δ^2h_0x	$C\tau$	
(82, 12, 44)	$\tau(\Delta e_1 + C_0)g$	$?M\Delta h_1^2d_0$	τ	
(82, 12, 45)	$(\Delta e_1 + C_0)g$	τMh_0g^2	Lemma 7.130	
(82, 16, 49)	$h_2c_1g^3$	$h_1^{14}h_6c_0$	$C\tau$	
(83, 7, 43)	$\tau h_0^2h_6g$	$Ph_1h_6d_0$	τ	
(83, 7, 44)	$h_0^2h_6g$	$h_1h_6c_0d_0$	$C\tau$	
(83, 11, 44)	τ^2gC'	$M\Delta h_0^2e_0$	Lemma 7.132	crossing
(83, 11, 45)	$\Delta j_1 + \tau gC'$	$\tau^2Mh_1g^2$	Lemma 7.132	
(84, 9, 46)	h_2gD_3	B_6d_1	$C\tau$	
(85, 5, 44)	τh_1f_2	$h_1x_{87,7} + \tau^2g_2^2$	Lemma 7.135	
(85, 5, 45)	$h_2h_6c_1$	$h_1^2h_4h_6c_0$	$C\tau$	
(85, 7, 46)	$h_2h_4Q_3$	$h_0g_2^2$	$C\tau$	
(85, 17, 48)	$\tau\Delta h_1g^3$	$\tau^2e_0^4g$	mmf	
(86, 18, 51)	$\tau h_2^2g^4$	$h_1e_0^4g + Mh_1^7e_0g$	τ	
(87, 12, 48)	h_2^2gA'	$\Delta h_1^2g_2g$	Lemma 7.138	

Table 20: Possible hidden ν extensions

(s, f, w)	source	target
(70, 5, 36)	$h_0^3h_3h_6$	$?h_0h_4D_2$
(75, 6, 40)	$h_0h_3d_2$	$?M\Delta h_1^2h_3$
(81, 4, 42)	$h_0h_2h_4h_6$	$?M\Delta h_1d_0$
(86, 11, 44)	$h_0^3h_6i$	$? \tau\Delta^2h_1f_1$
(87, 5, 46)	$h_1^2c_3$	$? \tau M\Delta h_1g$
(87, 7, 45)	$x_{87,7}$	$? \tau^2M\Delta h_1g$
(87, 9, 46)	τ^2gQ_3	$? \tau M\Delta h_1g$
(87, 10, 45)	$\tau\Delta h_1H_1$	$? \tau^2M\Delta h_1g$

Table 21: Miscellaneous hidden extensions

(s, f, w)	type	source	target	proof
(16, 2, 9)	σ	h_1h_4	h_4c_0	$C\tau$
(20, 4, 11)	ϵ	τg	d_0^2	[16 , Table 33]
(30, 2, 16)	σ	h_4^2	x	$C\tau$
(30, 2, 16)	η_4	h_4^2	$h_1h_5d_0$	[16 , Table 33]
(32, 6, 17)	ϵ	Δh_1h_3	$\Delta h_1^2d_0$	tmf
(32, 6, 17)	κ	Δh_1h_3	$\tau d_0l + \Delta c_0d_0$	tmf
(44, 4, 24)	θ_4	g_2	$x_{74,8}$	$C\tau$
(45, 3, 23)	ϵ	$\tau h_3^2h_5$	MP	Lemma 7.140
(45, 3, 24)	ϵ	$h_3^2h_5$	Mc_0	Lemma 7.140
(45, 3, 24)	κ	$h_3^2h_5$	Md_0	Lemma 7.142
(45, 3, 24)	$\bar{\kappa}$	$h_3^2h_5$	τMg	Lemma 7.144
(45, 3, 24)	$\{\Delta h_1h_3\}$	$h_3^2h_5$	$M\Delta h_1h_3$	Lemma 7.145
(45, 3, 24)	$\theta_{4.5}$	$h_3^2h_5$	M^2	Lemma 7.146
(62, 2, 32)	σ	h_5^2	p'	$C\tau$
(62, 2, 32)	κ	h_5^2	h_0h_4A	Lemma 7.149
(62, 2, 32)	ρ_{15}	h_5^2	$?h_0x_{77,7}$ $? \tau^2 m_1$	Lemma 7.150
(62, 2, 32)	θ_4	h_5^2	$h_0^2g_3$	$C\tau$
(63, 7, 33)	ϵ	$\tau X_2 + \tau C'$	d_0Q_2	$C\tau$
(63, 7, 33)	κ	$\tau X_2 + \tau C'$	$M\Delta h_1h_3$	$C\tau$
(63, 7, 33)	η_4	$\tau X_2 + \tau C'$	$h_1x_{78,9}$	$C\tau$
(64, 2, 33)	ρ_{15}	h_1h_6	Ph_6c_0	Lemma 7.151
(64, 2, 33)	ρ_{23}	h_1h_6	$P^2h_6c_0$	Lemma 7.151
(65, 10, 35)	ϵ	τMg	Md_0^2	Lemma 7.152
(69, 4, 36)	σ	p'	h_0h_4A	$C\tau$
(77, 12, 41)	ϵ	$M\Delta h_1h_3$	$?M\Delta h_1^2d_0$	Lemma 7.153
(79, 3, 41)	σ	$h_1h_4h_6$	$h_4h_6c_0$	$C\tau$

Bibliography

- [1] J. F. Adams, *On the non-existence of elements of Hopf invariant one*, Ann. of Math. (2) **72** (1960), 20–104. MR0141119 (25 #4530)
- [2] M. G. Barratt, J. D. S. Jones, and M. E. Mahowald, *Relations amongst Toda brackets and the Kervaire invariant in dimension 62*, J. London Math. Soc. (2) **30** (1984), no. 3, 533–550, DOI 10.1112/jlms/s2-30.3.533. MR810962 (87g:55025)
- [3] M. G. Barratt, M. E. Mahowald, and M. C. Tangora, *Some differentials in the Adams spectral sequence. II*, Topology **9** (1970), 309–316. MR0266215 (42 #1122)
- [4] Mark Behrens, Michael Hill, Michael J. Hopkins, and Mark Mahowald, *Detecting exotic spheres in low dimensions using coker J* , J. London Math. Soc., to appear.
- [5] Robert Burklund, Daniel C. Isaksen, and Zhouli Xu, *Synthetic stable stems*, in preparation.
- [6] Robert Bruner, *A new differential in the Adams spectral sequence*, Topology **23** (1984), no. 3, 271–276, DOI 10.1016/0040-9383(84)90010-7. MR770563 (86c:55016)
- [7] Robert R. Bruner, *The cohomology of the mod 2 Steenrod algebra: A computer calculation*, Wayne State University Research Report **37** (1997).
- [8] ———, *Squaring operations in $\text{Ext}_{\mathcal{A}}(F_2, F_2)$* (2004), preprint.
- [9] Bogdan Gheorghe, *The motivic cofiber of τ* , Doc. Math. **23** (2018), 1077–1127. MR3874951
- [10] Bogdan Gheorghe, Daniel C. Isaksen, Achim Krause, and Nicolas Ricka, *\mathbb{C} -motivic modular forms* (2018), submitted, available at [arXiv:1810.11050](https://arxiv.org/abs/1810.11050).
- [11] Bogdan Gheorghe, Guozhen Wang, and Zhouli Xu, *The special fiber of the motivic deformation of the stable homotopy category is algebraic* (2018), submitted, available at [arXiv:1809.09290](https://arxiv.org/abs/1809.09290).
- [12] Bertrand J. Guillou and Daniel C. Isaksen, *The η -local motivic sphere*, J. Pure Appl. Algebra **219** (2015), no. 10, 4728–4756, DOI 10.1016/j.jpaa.2015.03.004. MR3346515
- [13] M. A. Hill, M. J. Hopkins, and D. C. Ravenel, *On the nonexistence of elements of Kervaire invariant one*, Ann. of Math. (2) **184** (2016), no. 1, 1–262, DOI 10.4007/annals.2016.184.1.1. MR3505179
- [14] Daniel C. Isaksen, *The cohomology of motivic $A(2)$* , Homology Homotopy Appl. **11** (2009), no. 2, 251–274. MR2591921 (2011c:55034)
- [15] ———, *When is a fourfold Massey product defined?*, Proc. Amer. Math. Soc. **143** (2015), no. 5, 2235–2239, DOI 10.1090/S0002-9939-2014-12387-2. MR3314129
- [16] ———, *Stable stems*, Mem. Amer. Math. Soc. **262** (2019), no. 1269, viii+159, DOI 10.1090/memo/1269. MR4046815
- [17] ———, *The homotopy of \mathbb{C} -motivic modular forms*, preprint, available at [arXiv:1811.07937](https://arxiv.org/abs/1811.07937).
- [18] ———, *The Mahowald operator in the cohomology of the Steenrod algebra* (2020), preprint.
- [19] Daniel C. Isaksen, Guozhen Wang, and Zhouli Xu, *Classical and motivic Adams charts* (2020), preprint, available at s.wayne.edu/isaksen/adams-charts.
- [20] ———, *Classical and motivic Adams-Novikov charts* (2020), preprint, available at s.wayne.edu/isaksen/adams-charts.
- [21] ———, *Classical algebraic Novikov charts and \mathbb{C} -motivic Adams charts for the cofiber of τ* (2020), preprint, available at s.wayne.edu/isaksen/adams-charts.
- [22] Daniel C. Isaksen and Zhouli Xu, *Motivic stable homotopy and the stable 51 and 52 stems*, Topology Appl. **190** (2015), 31–34, DOI 10.1016/j.topol.2015.04.008. MR3349503
- [23] Michel A. Kervaire and John W. Milnor, *Groups of homotopy spheres. I*, Ann. of Math. (2) **77** (1963), 504–537, DOI 10.2307/1970128. MR148075
- [24] Stanley O. Kochman, *A chain functor for bordism*, Trans. Amer. Math. Soc. **239** (1978), 167–196, DOI 10.2307/1997852. MR488031

- [25] ———, *Stable homotopy groups of spheres*, Lecture Notes in Mathematics, vol. 1423, Springer-Verlag, Berlin, 1990. A computer-assisted approach. MR1052407 (91j:55016)
- [26] Stanley O. Kochman and Mark E. Mahowald, *On the computation of stable stems*, The Čech centennial (Boston, MA, 1993), Contemp. Math., vol. 181, Amer. Math. Soc., Providence, RI, 1995, pp. 299–316, DOI 10.1090/conm/181/02039. MR1320997 (96j:55018)
- [27] Achim Krause, *Periodicity in motivic homotopy theory and over BP_*BP* , Ph.D. thesis, Universität Bonn, 2018.
- [28] Wen-Hsiung Lin, *A proof of the strong Kervaire invariant in dimension 62*, First International Congress of Chinese Mathematicians (Beijing, 1998), AMS/IP Stud. Adv. Math., vol. 20, Amer. Math. Soc., Providence, RI, 2001, pp. 351–358. MR1830191
- [29] Mark Mahowald and Martin Tangora, *Some differentials in the Adams spectral sequence*, Topology **6** (1967), 349–369. MR0214072 (35 #4924)
- [30] J. Peter May, *Matric Massey products*, J. Algebra **12** (1969), 533–568. MR0238929 (39 #289)
- [31] Haynes Robert Miller, *SOME ALGEBRAIC ASPECTS OF THE ADAMS-NOVIKOV SPECTRAL SEQUENCE*, ProQuest LLC, Ann Arbor, MI, 1975. Thesis (Ph.D.)—Princeton University. MR2625232
- [32] John Milnor, *Differential topology forty-six years later*, Notices Amer. Math. Soc. **58** (2011), no. 6, 804–809. MR2839925
- [33] R. Michael F. Moss, *Secondary compositions and the Adams spectral sequence*, Math. Z. **115** (1970), 283–310. MR0266216 (42 #1123)
- [34] S. P. Novikov, *Methods of algebraic topology from the point of view of cobordism theory*, Izv. Akad. Nauk SSSR Ser. Mat. **31** (1967), 855–951 (Russian). MR0221509
- [35] Piotr Pstragowski, *Synthetic spectra and the cellular motivic category* (2018), preprint, available at [arXiv:1803.01804](https://arxiv.org/abs/1803.01804).
- [36] Douglas C. Ravenel, *Complex cobordism and stable homotopy groups of spheres*, Pure and Applied Mathematics, vol. 121, Academic Press, Inc., Orlando, FL, 1986. MR860042 (87j:55003)
- [37] Martin C. Tangora, *On the cohomology of the Steenrod algebra*, Math. Z. **116** (1970), 18–64. MR0266205 (42 #1112)
- [38] Hirosi Toda, *Composition methods in homotopy groups of spheres*, Annals of Mathematics Studies, No. 49, Princeton University Press, Princeton, N.J., 1962. MR0143217 (26 #777)
- [39] *Topological modular forms*, Mathematical Surveys and Monographs, vol. 201, American Mathematical Society, Providence, RI, 2014. Edited by Christopher L. Douglas, John Francis, André G. Henriques and Michael A. Hill. MR3223024
- [40] Vladimir Voevodsky, *Motivic cohomology with $\mathbf{Z}/2$ -coefficients*, Publ. Math. Inst. Hautes Études Sci. **98** (2003), 59–104, DOI 10.1007/s10240-003-0010-6. MR2031199 (2005b:14038b)
- [41] ———, *Motivic Eilenberg-MacLane spaces*, Publ. Math. Inst. Hautes Études Sci. **112** (2010), 1–99, DOI 10.1007/s10240-010-0024-9. MR2737977 (2012f:14041)
- [42] Guozhen Wang, github.com/pouiyter/morestablestems.
- [43] ———, *Computations of the Adams-Novikov E_2 -term* (2020), preprint, available at github.com/pouiyter/MinimalResolution.
- [44] Guozhen Wang and Zhouli Xu, *The triviality of the 61-stem in the stable homotopy groups of spheres*, Ann. of Math. (2) **186** (2017), no. 2, 501–580, DOI 10.4007/annals.2017.186.2.3. MR3702672
- [45] ———, *Some extensions in the Adams spectral sequence and the 51-stem*, Algebr. Geom. Topol. **18** (2018), no. 7, 3887–3906, DOI 10.2140/agt.2018.18.3887. MR3892234
- [46] Zhouli Xu, *The strong Kervaire invariant problem in dimension 62*, Geom. Topol. **20** (2016), no. 3, 1611–1624, DOI 10.2140/gt.2016.20.1611. MR3523064